

Dust coagulation with porosity evolution; effects on planetesimal formation and opacity evolution

Akimasa Kataoka (ITA, Heidelberg University)

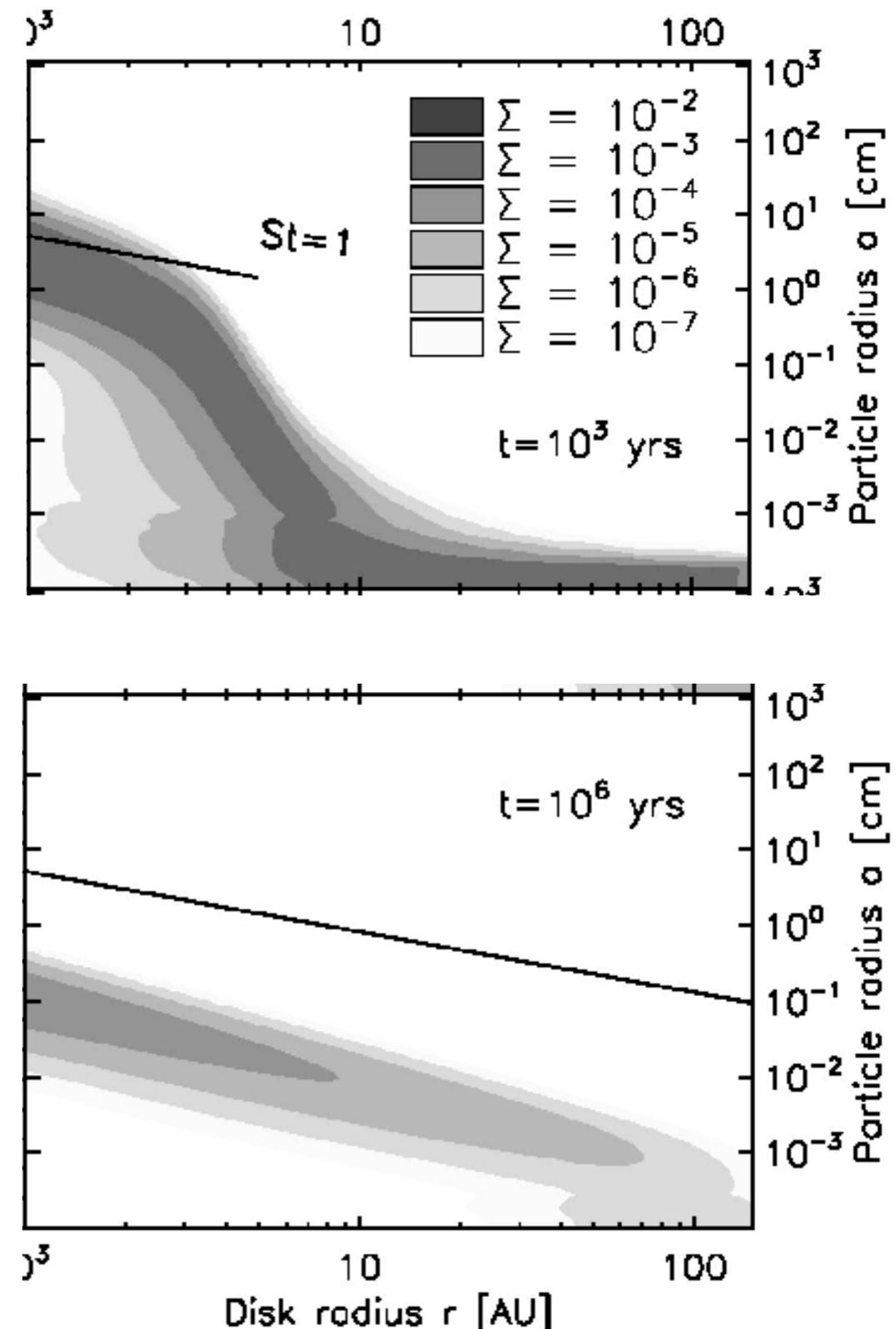
H.Tanaka (Hokkaido Univ.), S.Okuzumi (Tokyo Tech), C.P. Dullemond (ITA)

time = 0.00e+000

Planet formation in size evolution

Barriers in dust coagulation

- Radial drift barrier for planetesimal formation (e.g., Adachi et al. 1976)
 - Dust grains have to "jump" the barrier
- Radial drift barrier for millimeter-wave observations (e.g., Beckwith & Sargent 1991)
 - Dust grains has to "jump" the barrier and we should "keep" them
- Fragmentation barrier (e.g., Blum & Münch 1993)
 - Dust should overcome high-speed collisions such as ~ 50 m/s
- Bouncing barrier (e.g., Zsom et al. 2010)



Brauer et al. 2008

Possible solution

Porous dust aggregates

- Radial drift barrier for planetesimal formation
→ Rapid coagulation by large cross section (Okuzumi et al. 2012)
- Radial drift barrier for millimeter-wave observations
→ Do they account for millimeter-wave emission? Opacity?
- Fragmentation barrier
→ Ice is more sticky than silicate (Wada et al. 2009)
- Bouncing barrier
→ Highly porous dust aggregates do not bounce (Wada et al. 2011)



Wada et al. 2009

Dust aggregate

radius : a , mass: m
internal density: ρ

Possible solution

Porous dust aggregates

- Radial drift barrier for planetesimal formation
→ Rapid coagulation by large cross section (Okuzumi et al. 2012)
- Radial drift barrier for millimeter-wave observations
→ Do they account for millimeter-wave emission? Opacity?
- Fragmentation barrier
→ Ice is more sticky than silicate (Wada et al. 2009)
- Bouncing barrier
→ Highly porous dust aggregates do not bounce (Wada et al. 2011)



Numerical simulations

$$u_{\text{col,crit}} = \begin{cases} 80 (r/0.1\mu\text{m})^{-5/6} & [\text{m/s}] \text{ for ice} \\ 8 (r/0.1\mu\text{m})^{-5/6} & [\text{m/s}] \text{ for silicate} \end{cases}$$

(Wada et al. 2009, 2013)

Laboratory experiments

Quarz (Colwell et al. 2003), **SiO₂**, **MgSiO₃** (Blum and Wurm 2000), **Graphite**, **Al₂O₃** (Reisshaus et al. 2006), **Ice** (Gundlach and Blum 2015), **CO₂** (Musiolik et al. 2016)

Dust aggregate

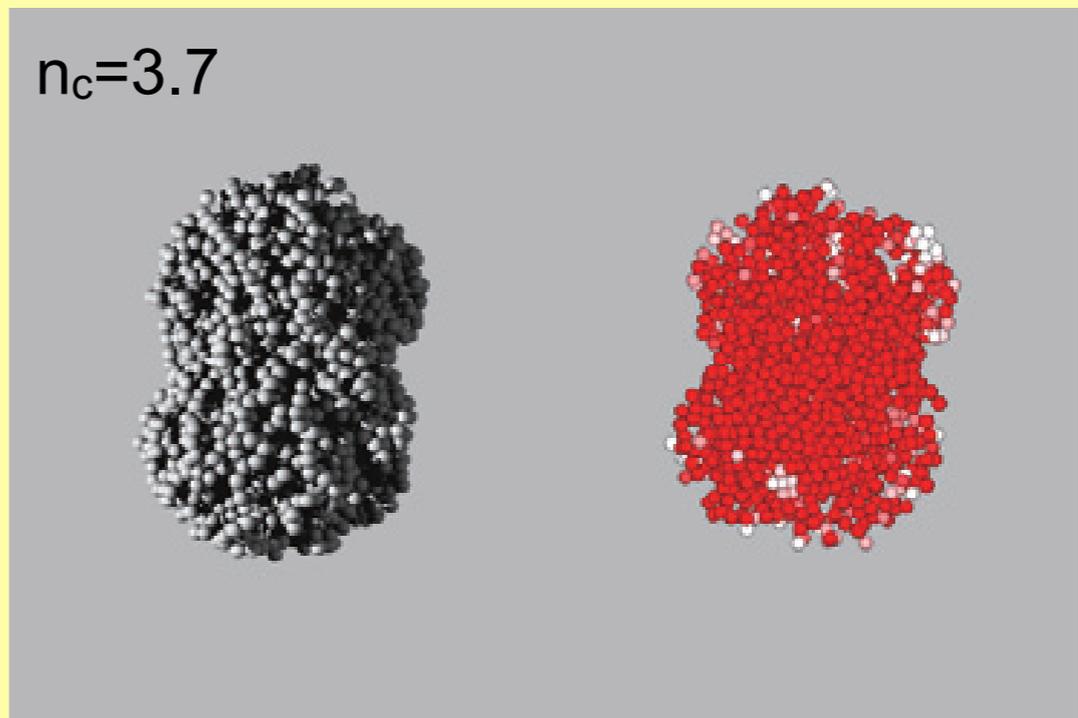
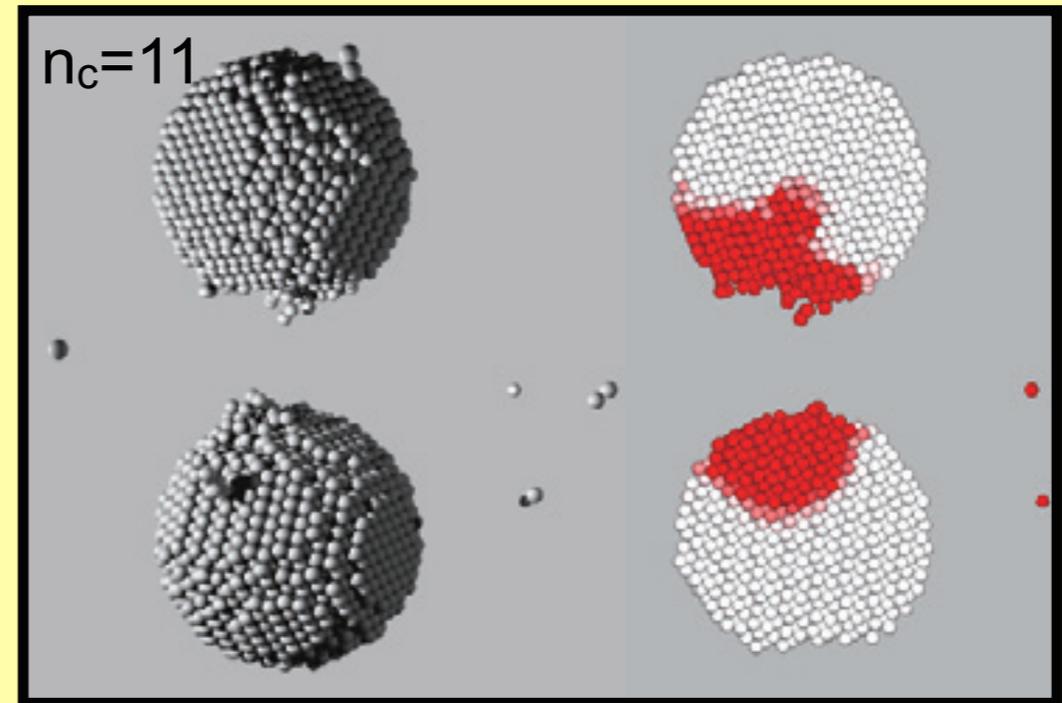
radius : a , mass: m
internal density: ρ

Possible solution

Porous dust aggregates

- Radial drift barrier for planetesimal formation
→ Rapid coagulation by large cross section (Okuzumi et al. 2012)
- Radial drift barrier for millimeter-wave observations
→ Do they account for millimeter-wave emission? Opacity?
- Fragmentation barrier
→ Ice is more sticky than silicate (Wada et al. 2009)
- Bouncing barrier
→ Highly porous dust aggregates do not bounce (Wada et al. 2011)

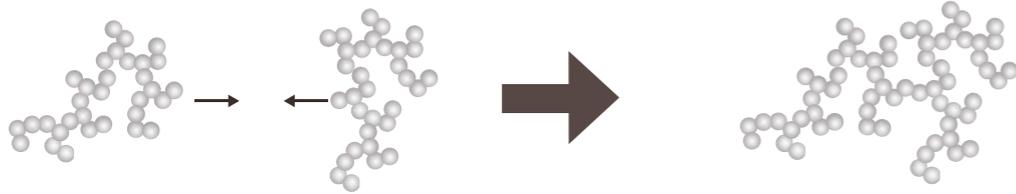
Coordination number : n_c



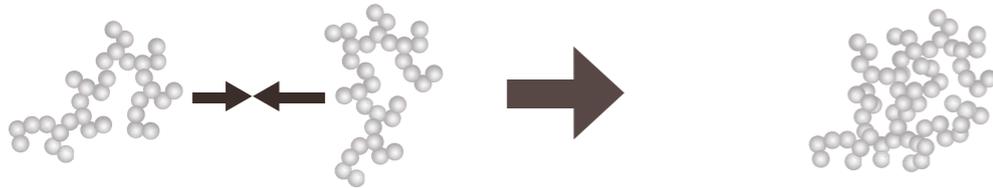
Wada et al. 2011

Structure evolution of dust aggregates

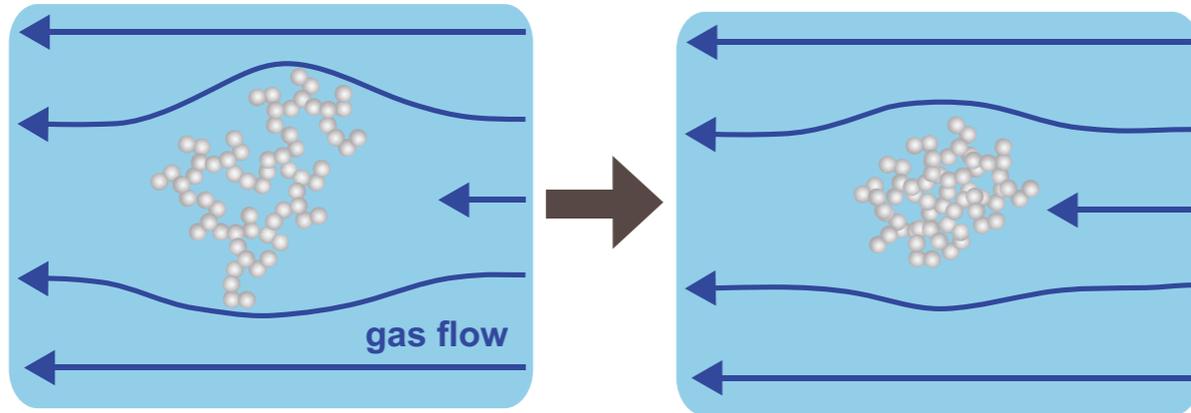
(a) Hit-and-stick



(b) Collisional compression



(c) Gas compression



(d) Self-gravitational compression



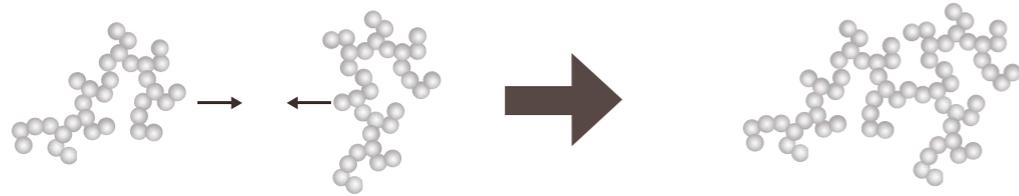
Kataoka et al. 2013b



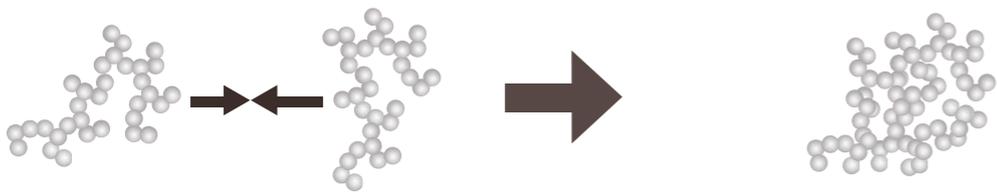
→ Growth with fractal dimension of ~ 2 due to low-velocity collisions

Structure evolution of dust aggregates

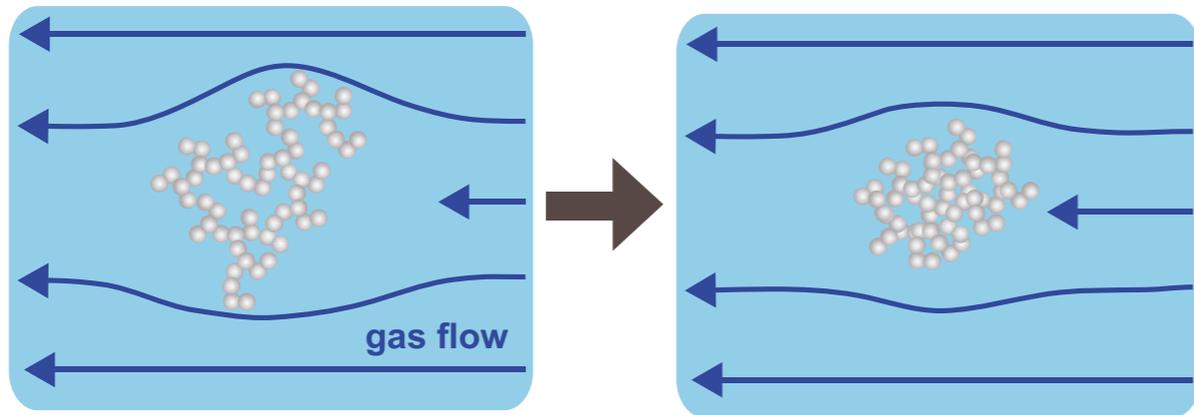
(a) Hit-and-stick



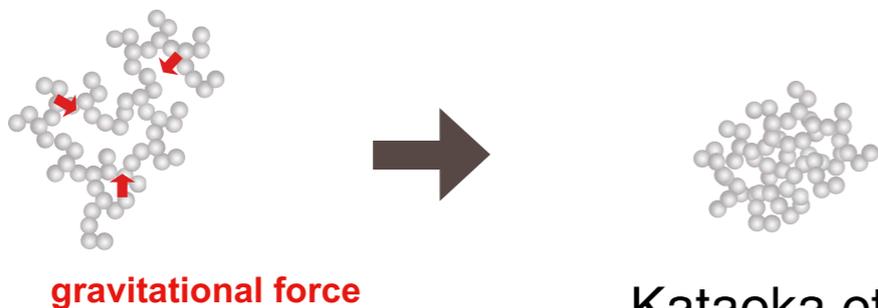
(b) Collisional compression



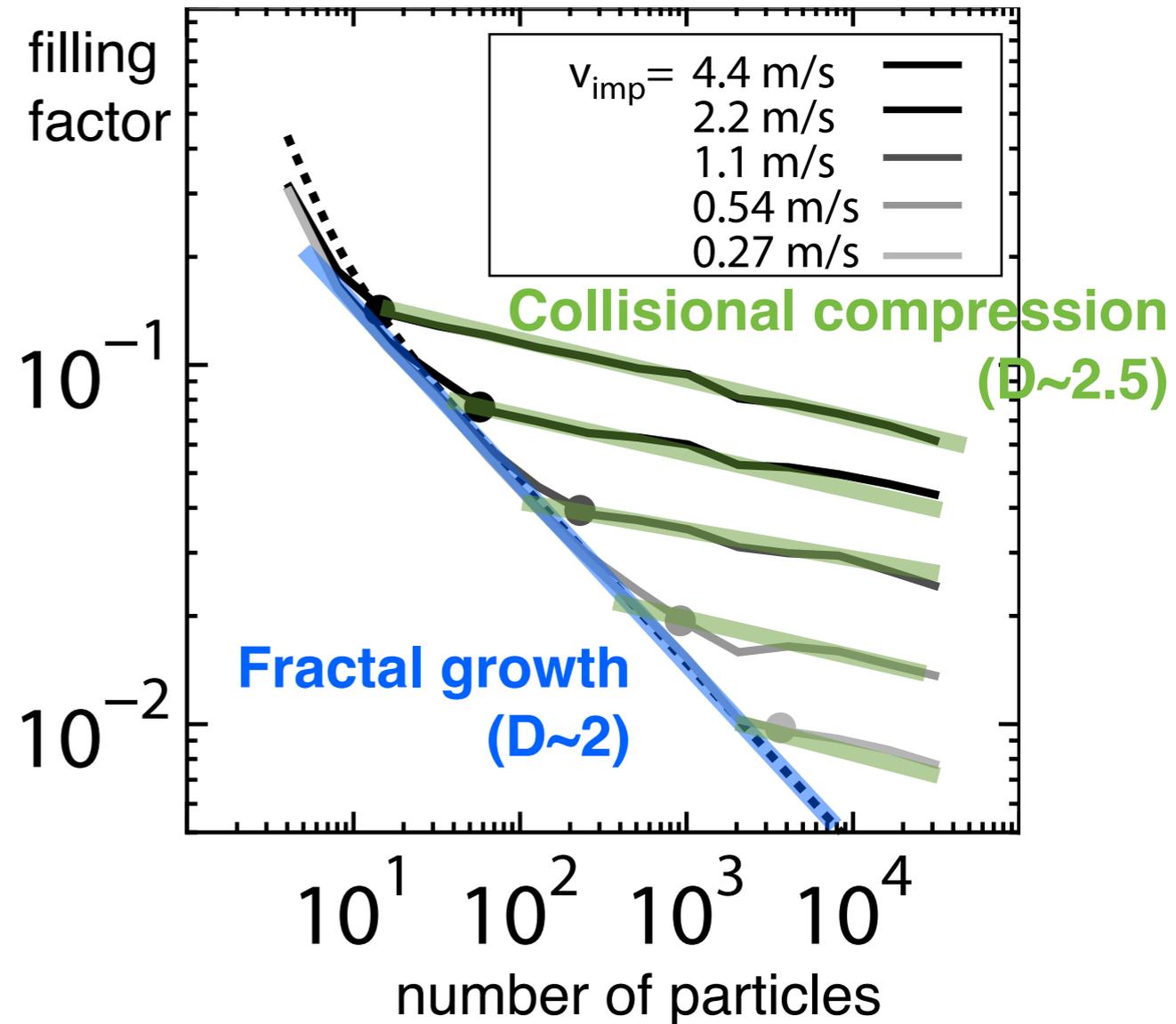
(c) Gas compression



(d) Self-gravitational compression



Kataoka et al. 2013b

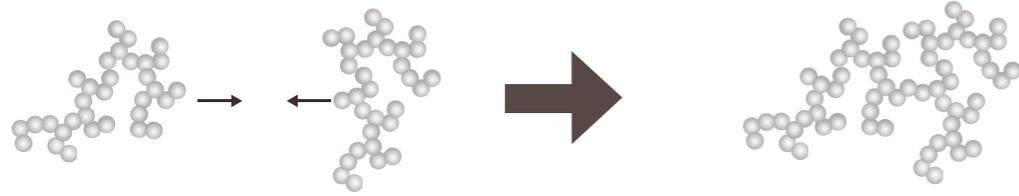


Suyama et al. 2008

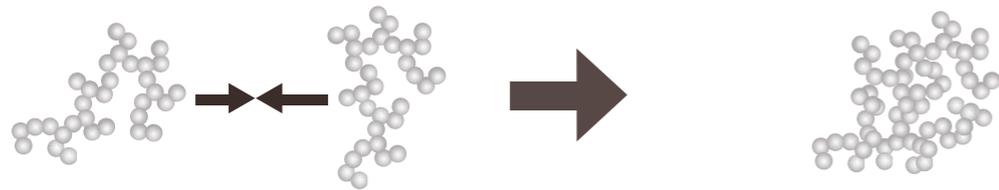
Collisional compression : ineffective to make aggregates compact

Structure evolution of dust aggregates

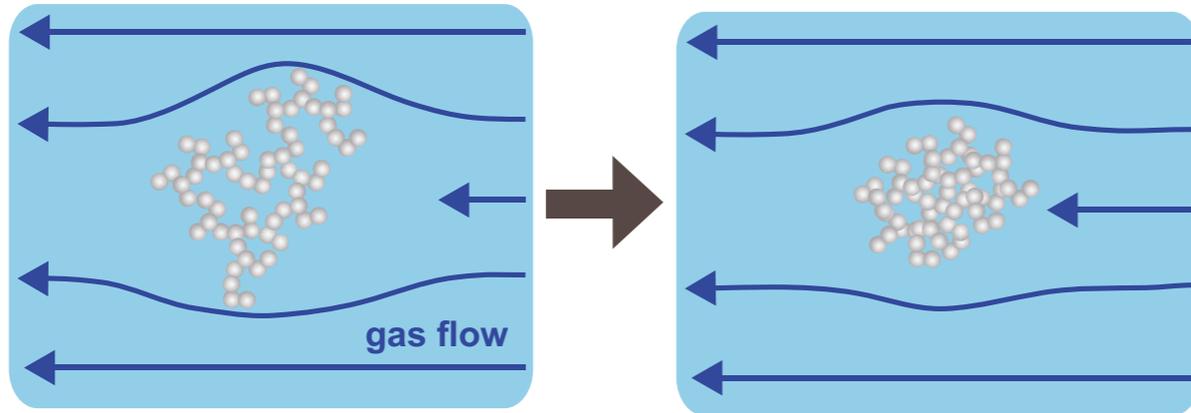
(a) Hit-and-stick



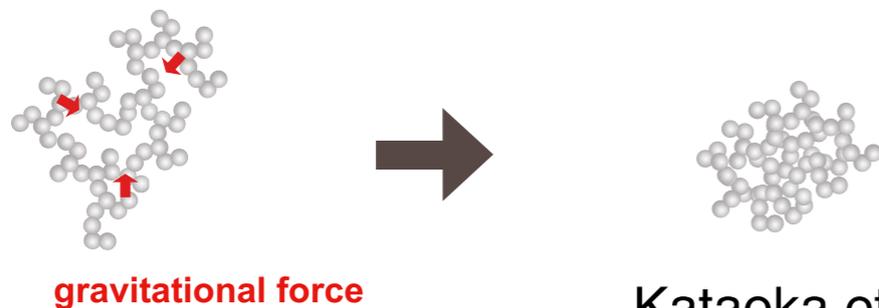
(b) Collisional compression



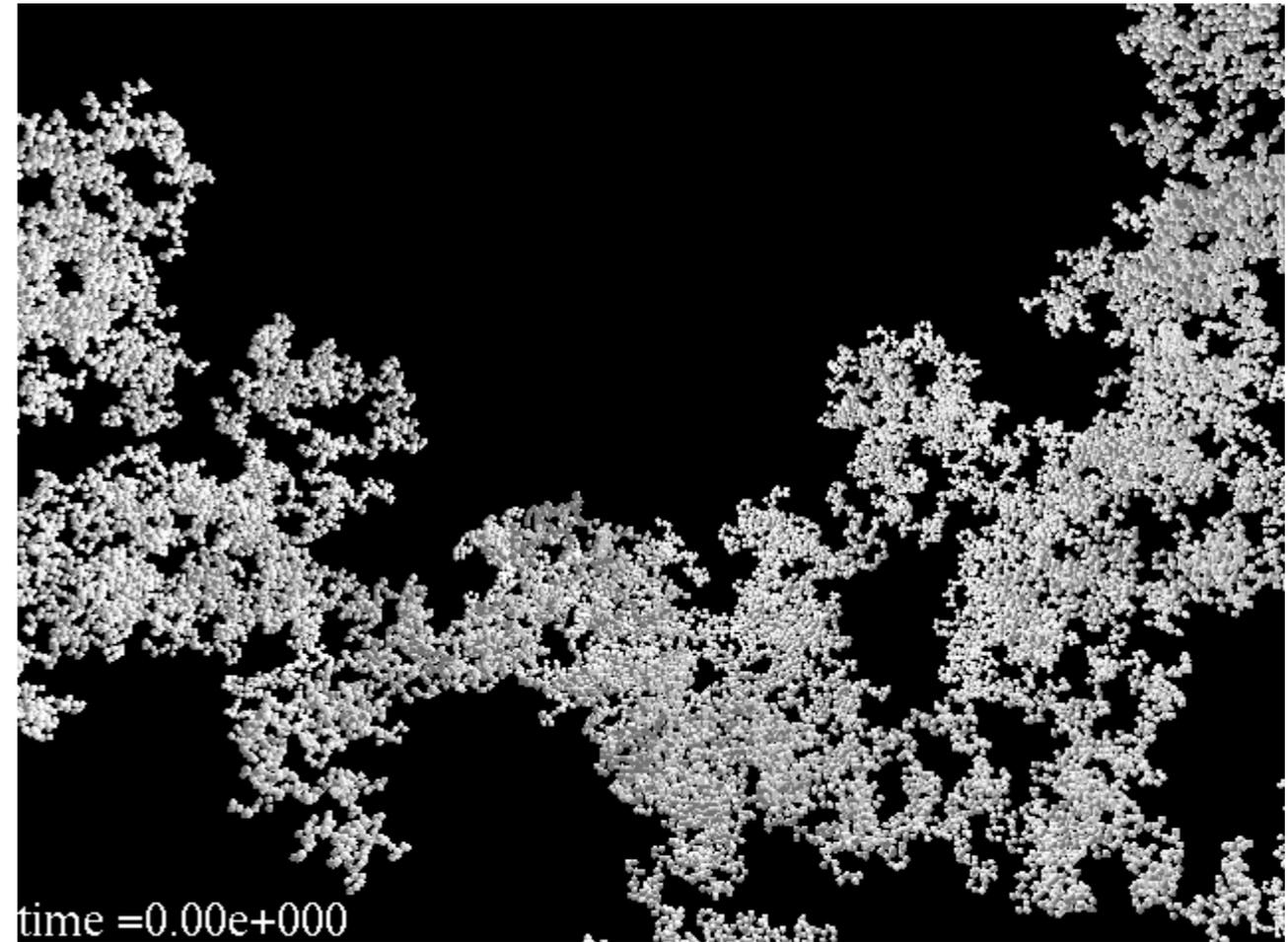
(c) Gas compression



(d) Self-gravitational compression



Kataoka et al. 2013b

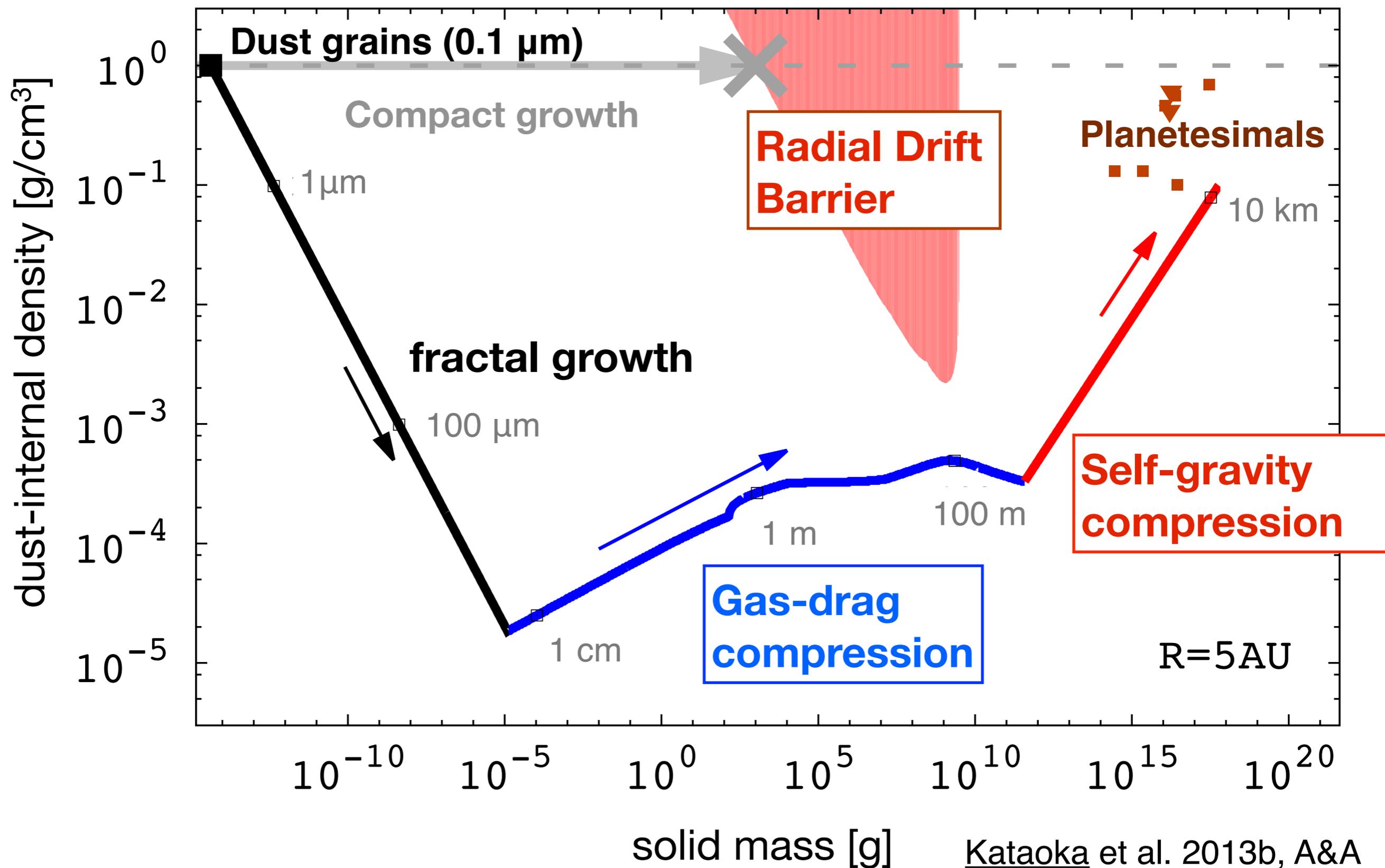


- compressive strength

$$P = \frac{E_{\text{roll}}}{r_0^3} \phi^3 \quad (\text{Kataoka et al. 2013a, A\&A})$$

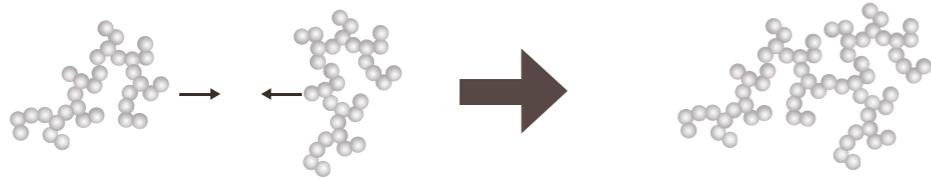
- external pressure
 - ram pressure of the gas
 - self-gravity

Planetesimal formation via fluffy aggregates

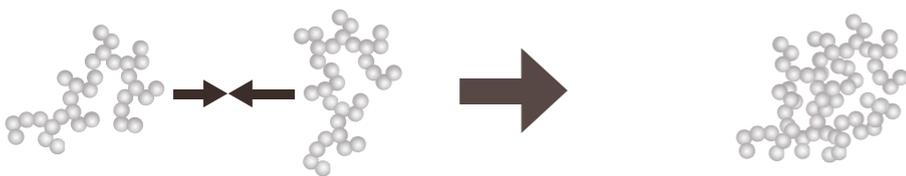


Dependence on the orbital radius

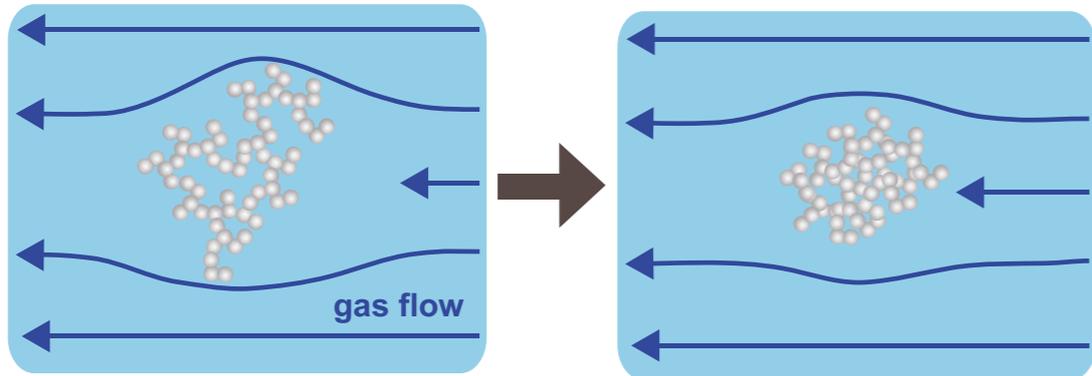
(a) Hit-and-stick



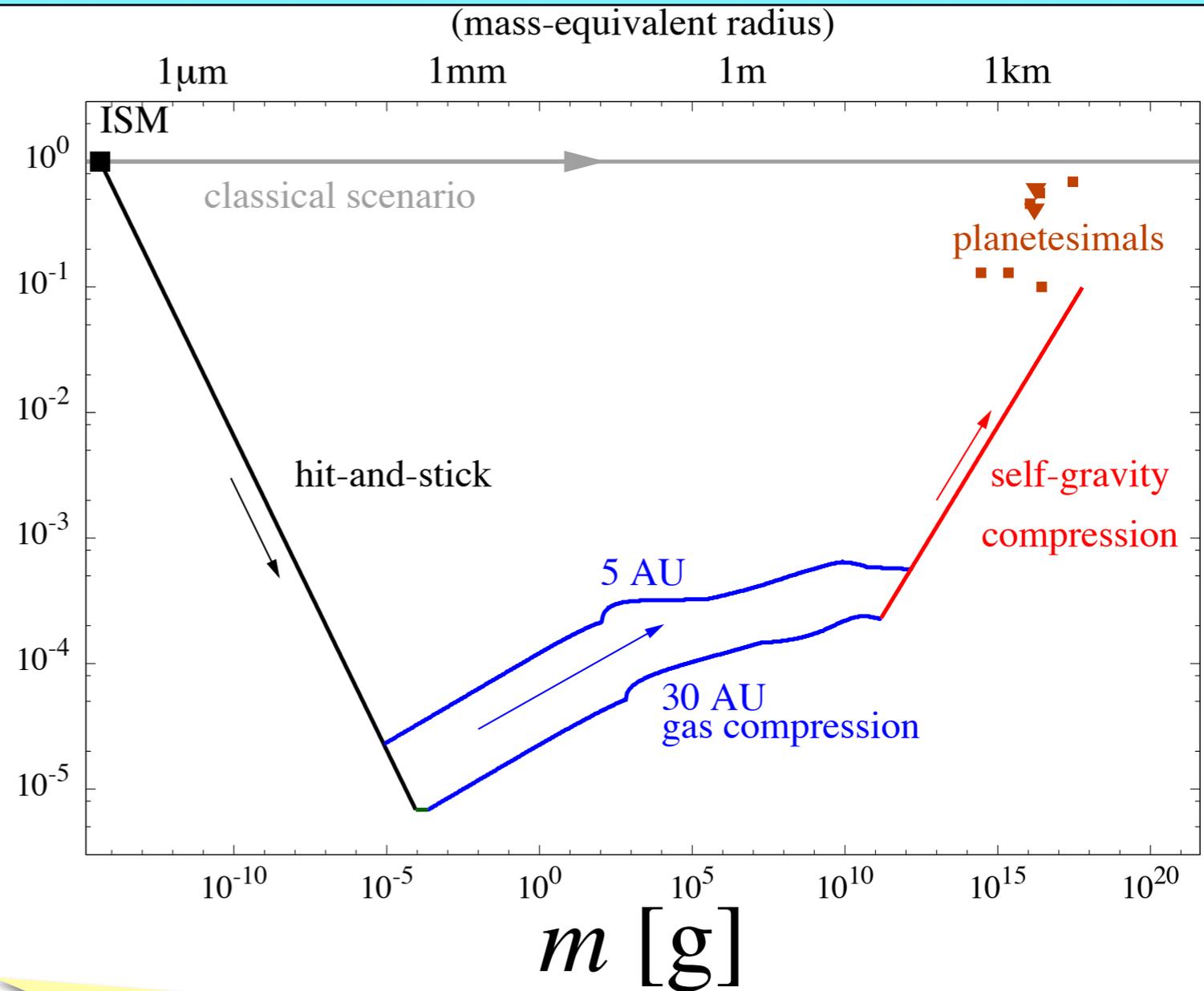
(b) Collisional compression



(c) Gas compression

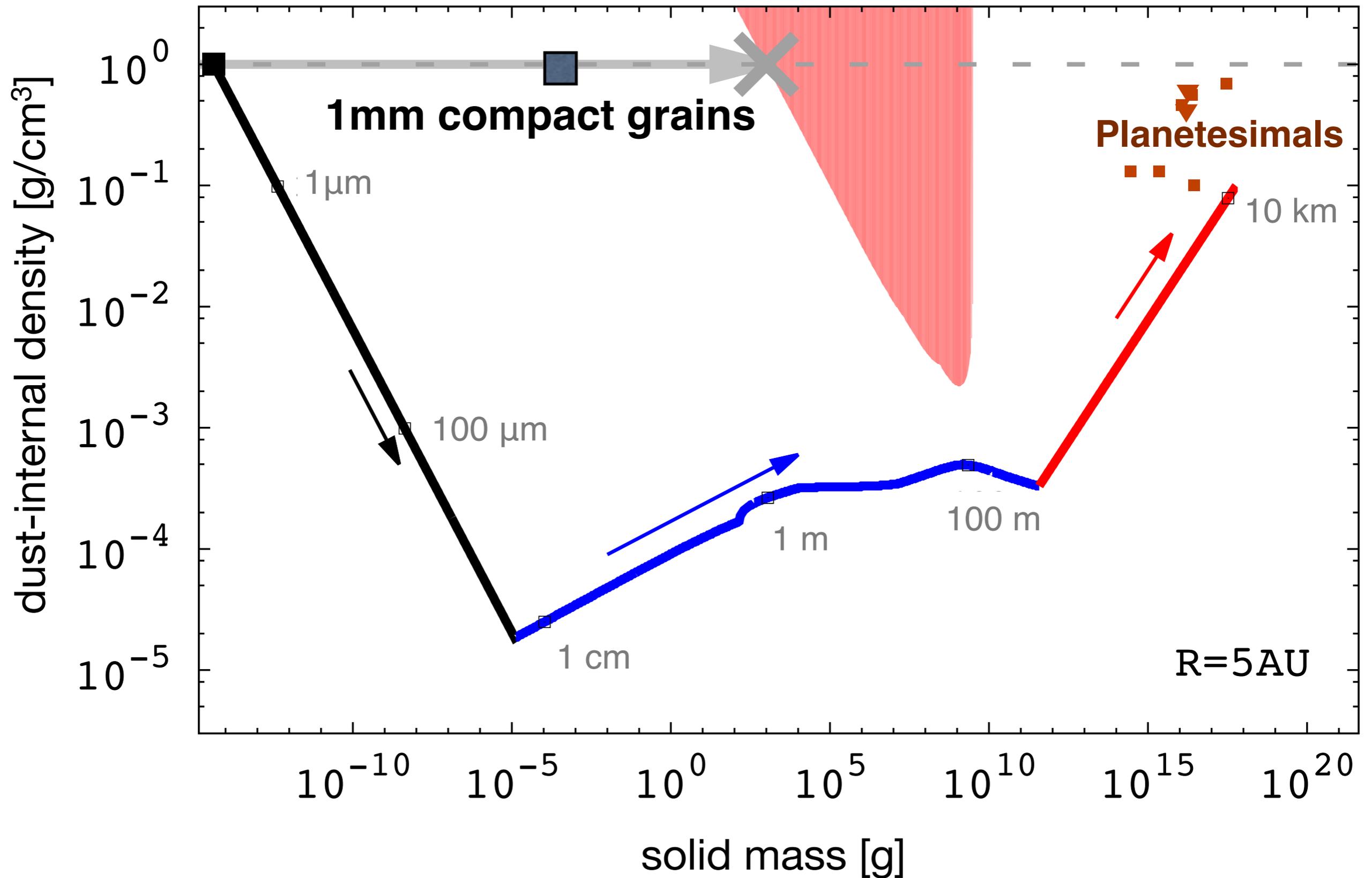


(d) Self-gravitational compression

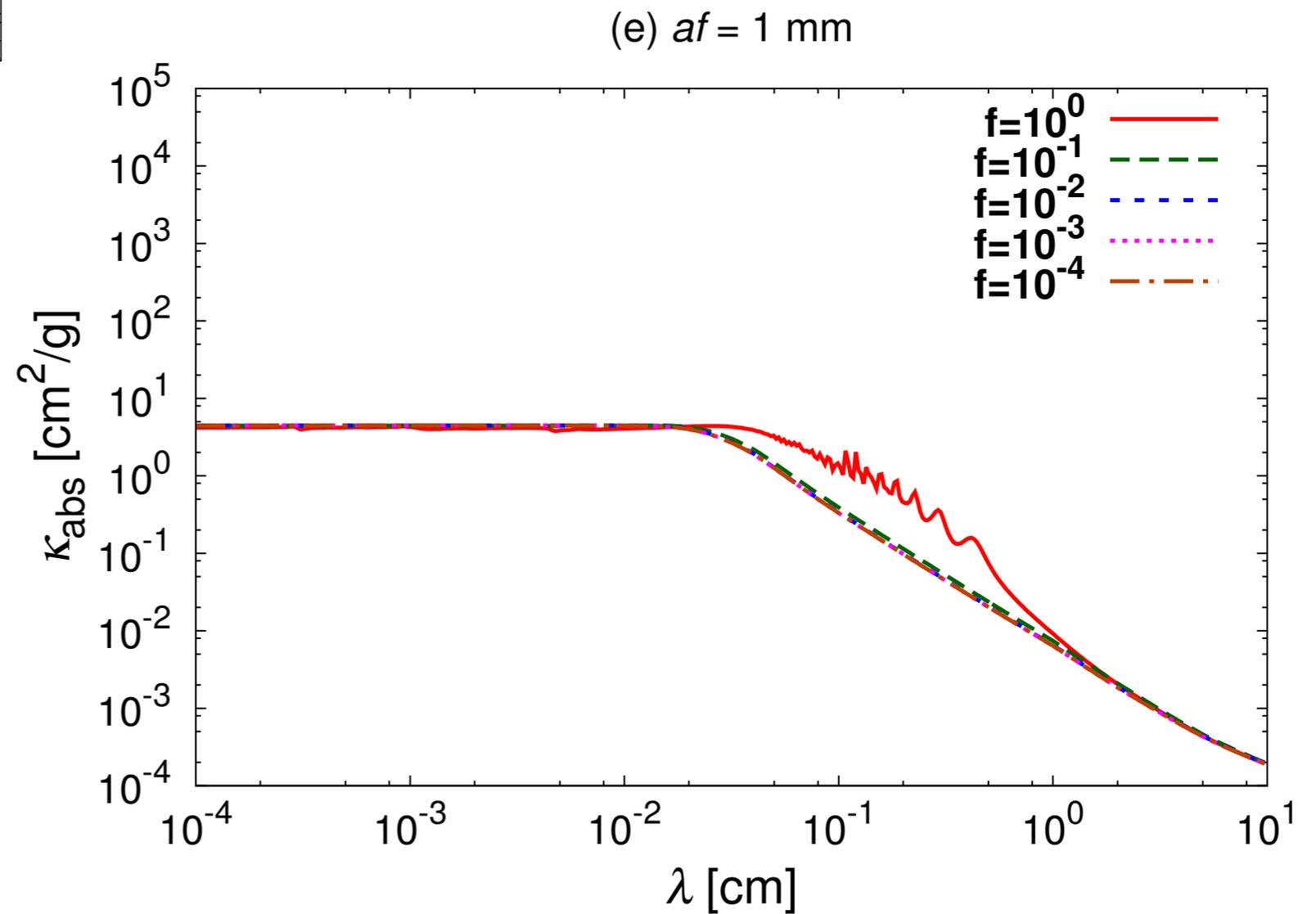
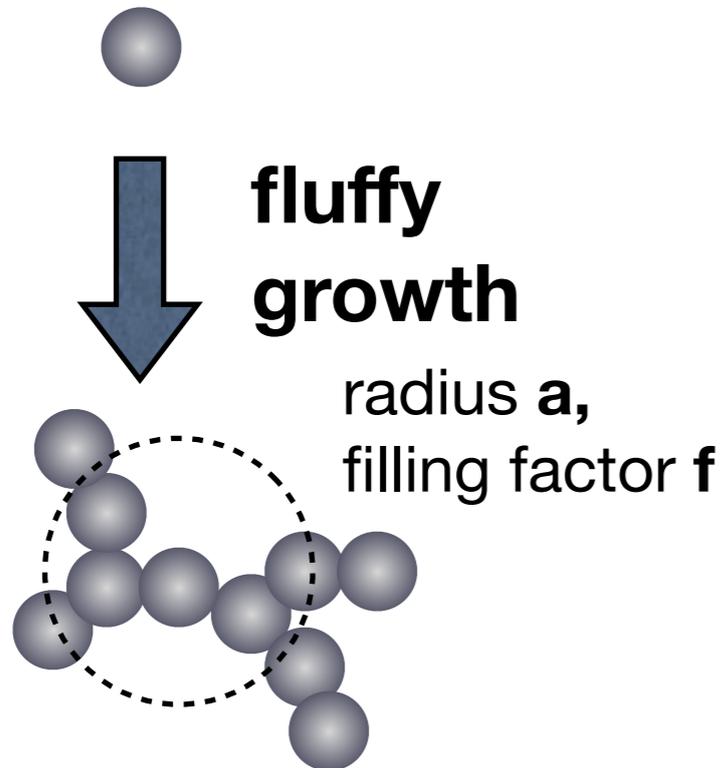
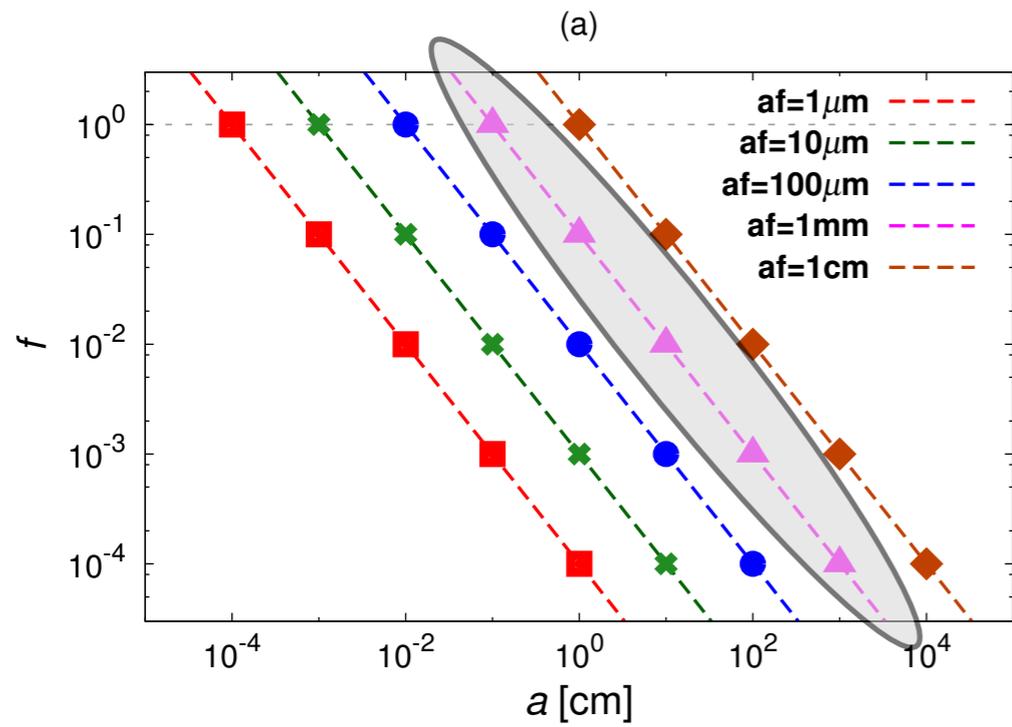


$$\rho_{\text{eq}} = \left(\frac{r_0^3 P_{\text{gas}}}{E_{\text{roll}}} \right)^{1/3} \rho_0.$$

What is the observed "mm-sized grains"?



Opacity evolution



How can we understand porosity?

Stokes number

$$St = \frac{3\Omega}{4\rho_g v_{th}} \frac{m}{\pi a^2} \quad (\text{Epstein regime})$$

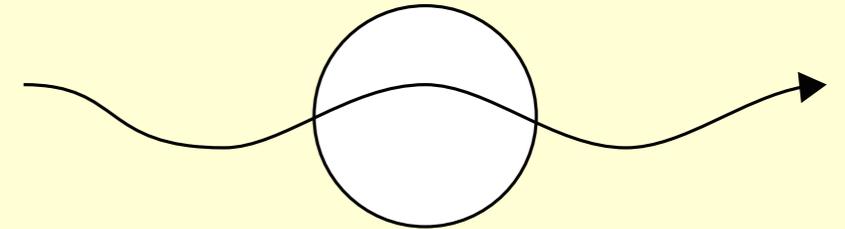
Opacity

$$\kappa_{abs} = \frac{\pi a^2}{m} Q_{abs} = f \left(\frac{m}{\pi a^2} \right)$$

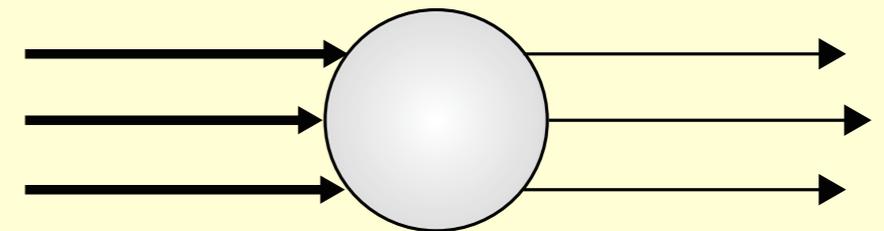
$$\frac{(\text{mass})}{(\text{area})} \sim \frac{a^3 f}{a^2} \sim a f$$

$af = \text{const} \Leftrightarrow \text{mass-to-area ratio} = \text{const}$

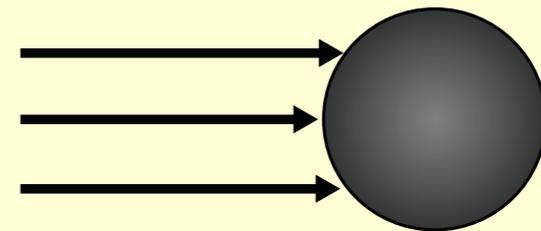
(a) $x < 1$



(b) $x > 1$, optically thin

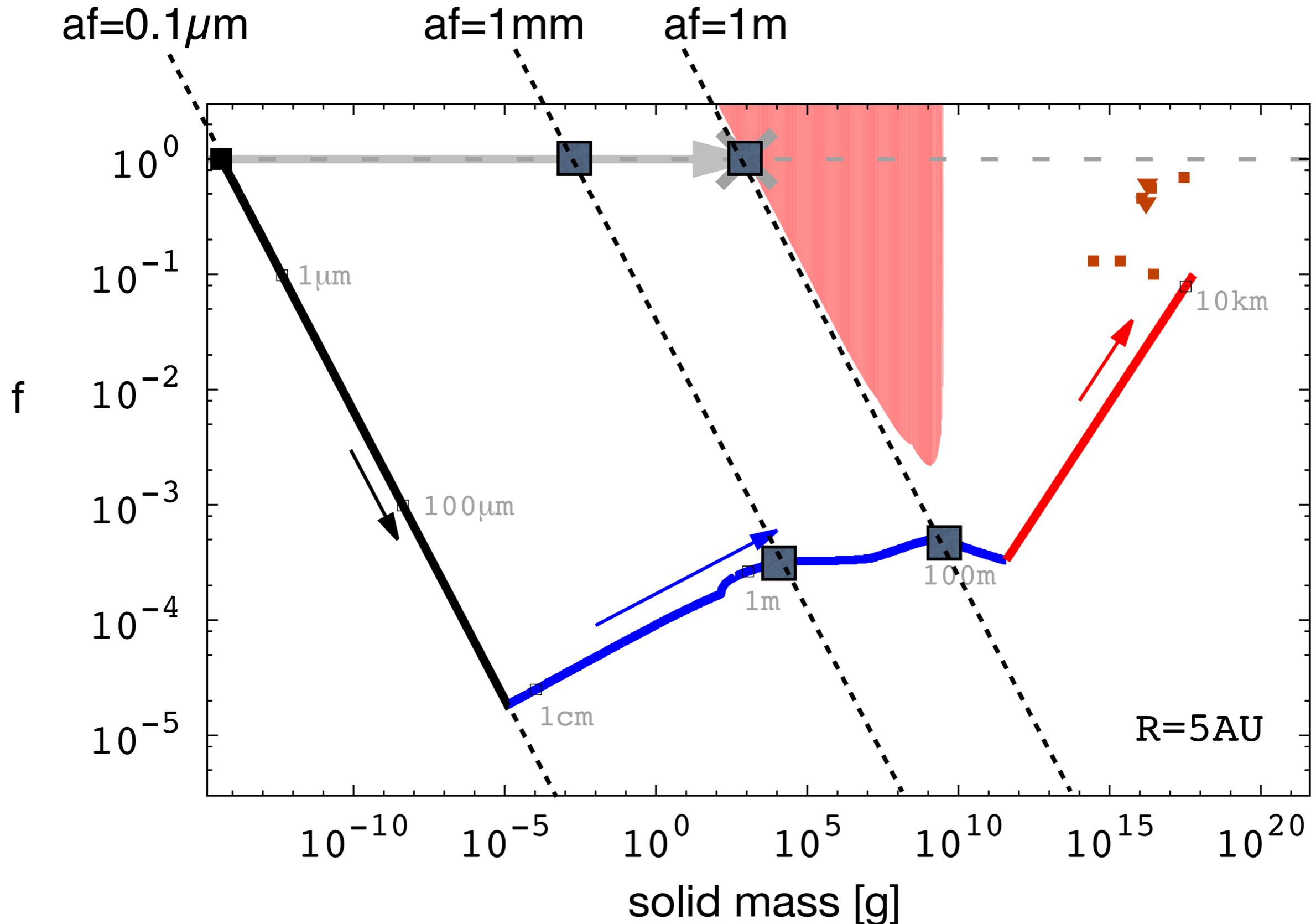


(c) $x > 1$, optically thick

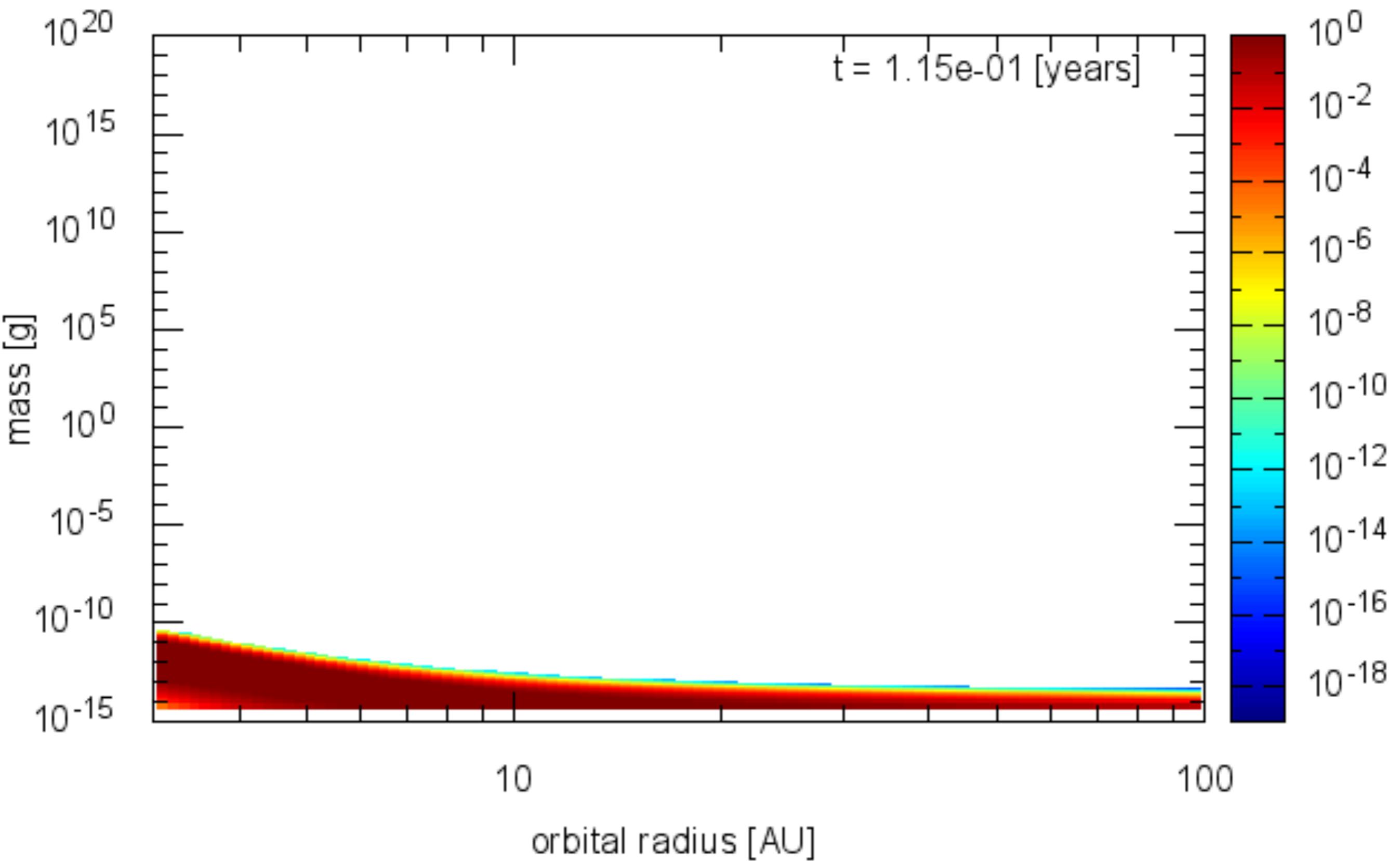


(Kataoka et al. 2014)

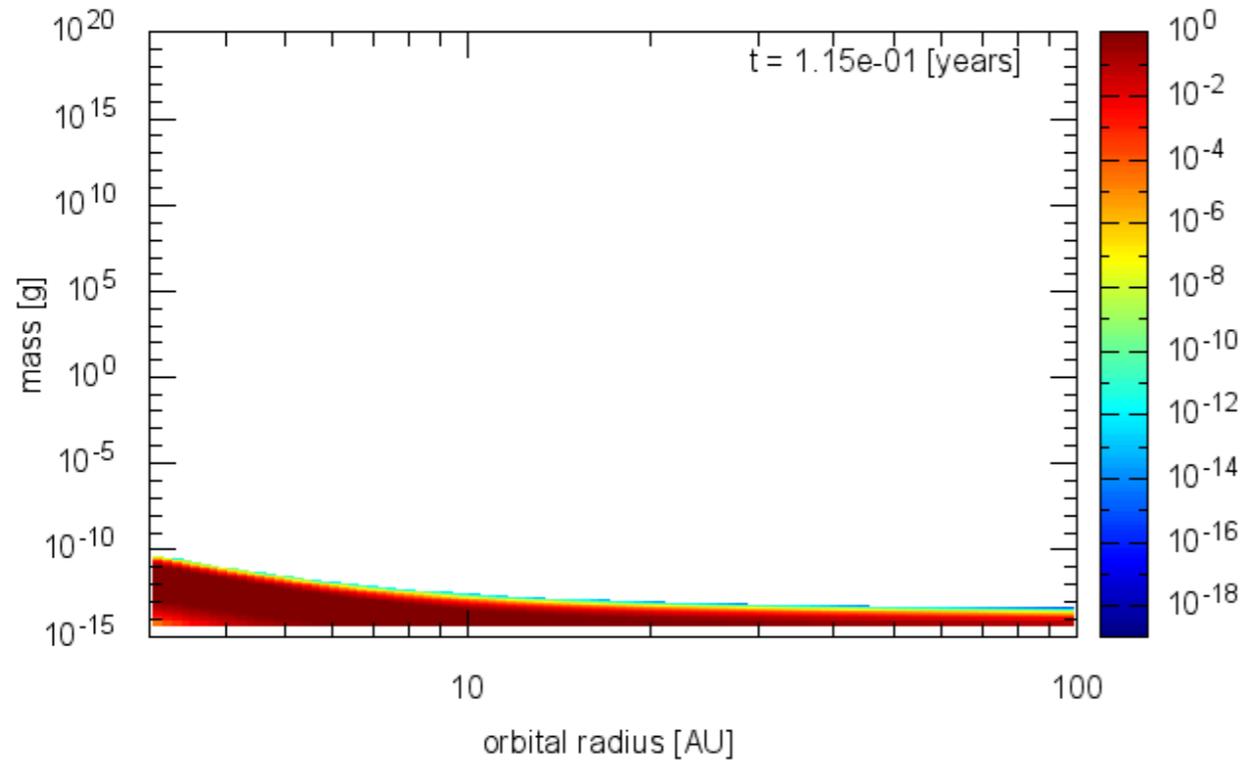
What are “mm-sized grains”?



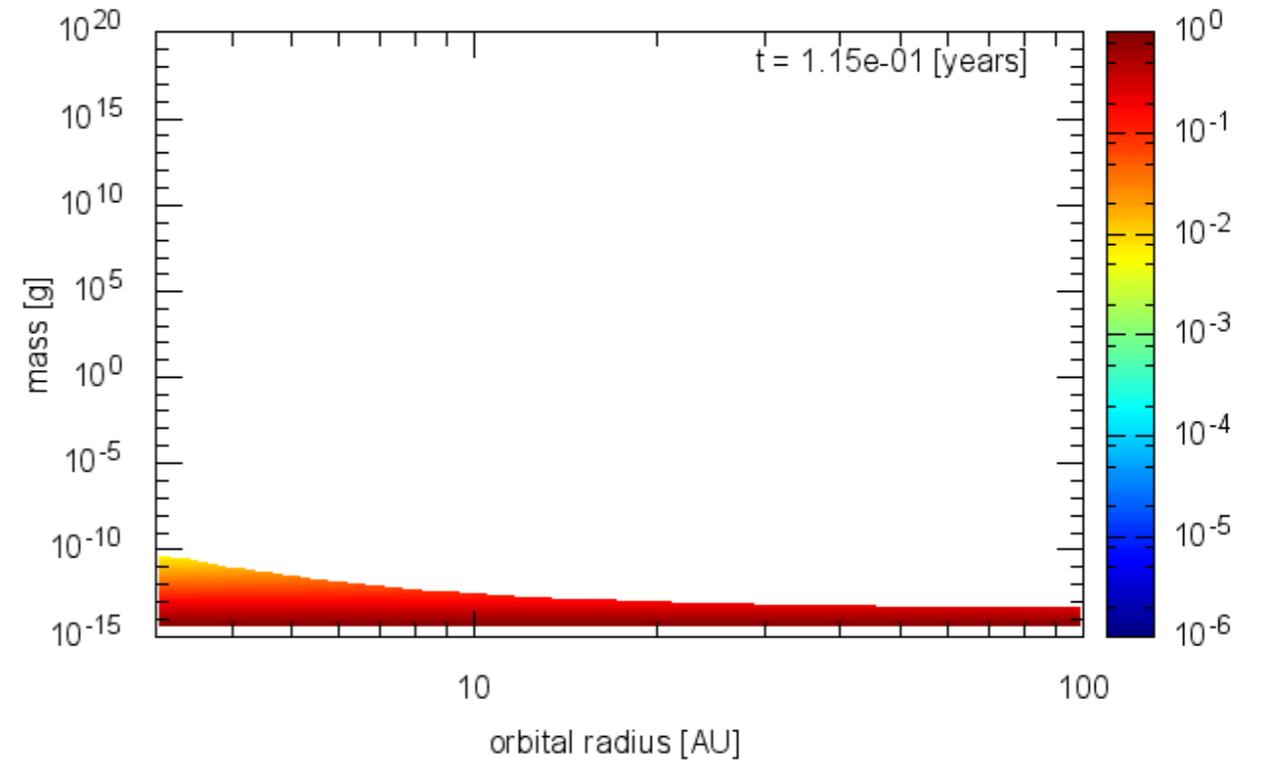
dust spatial density



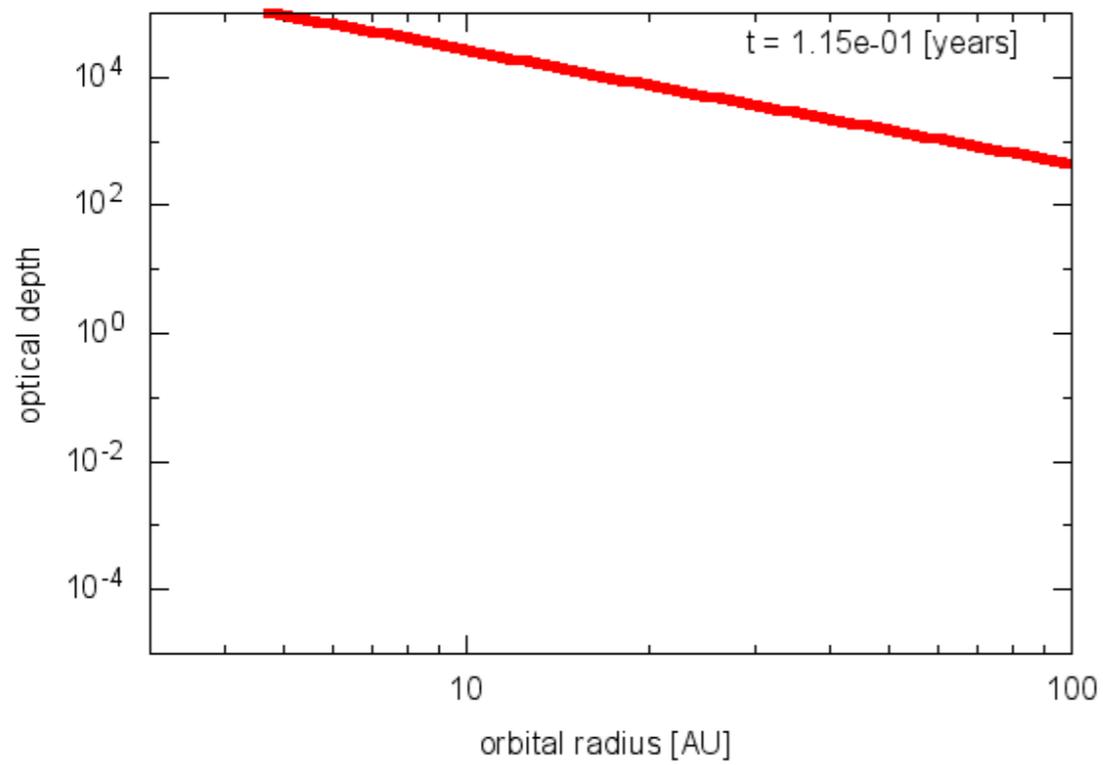
Dust density



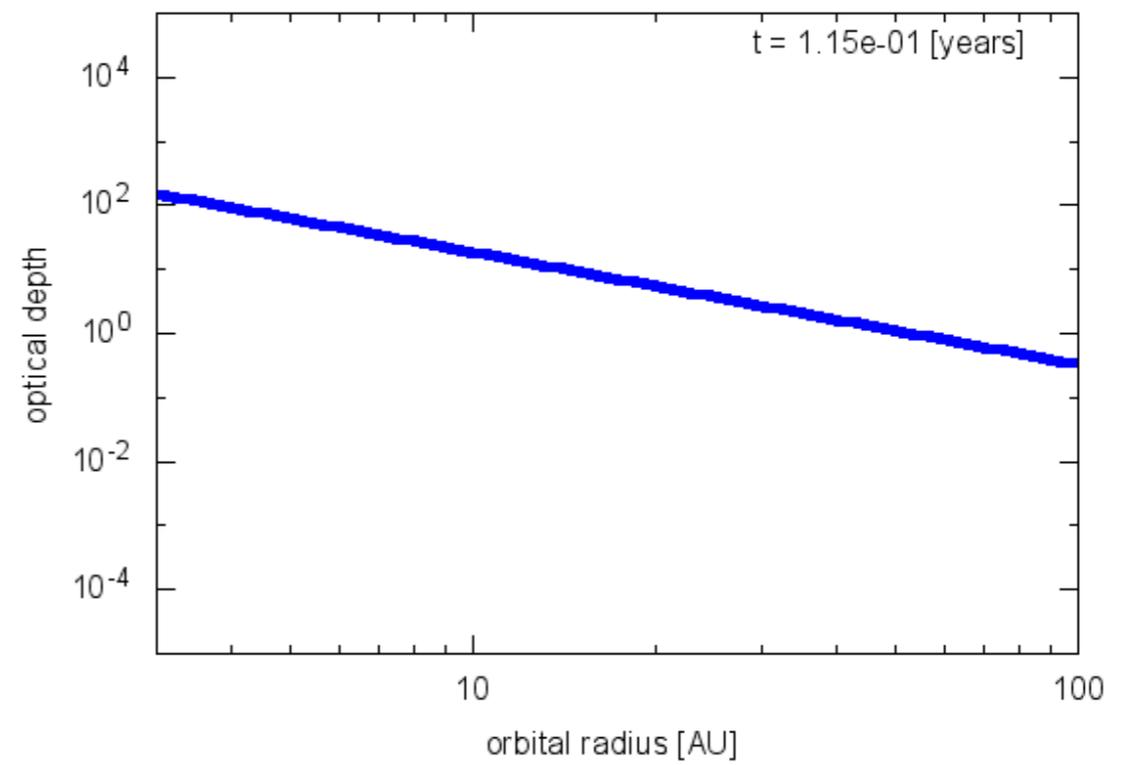
filling factor

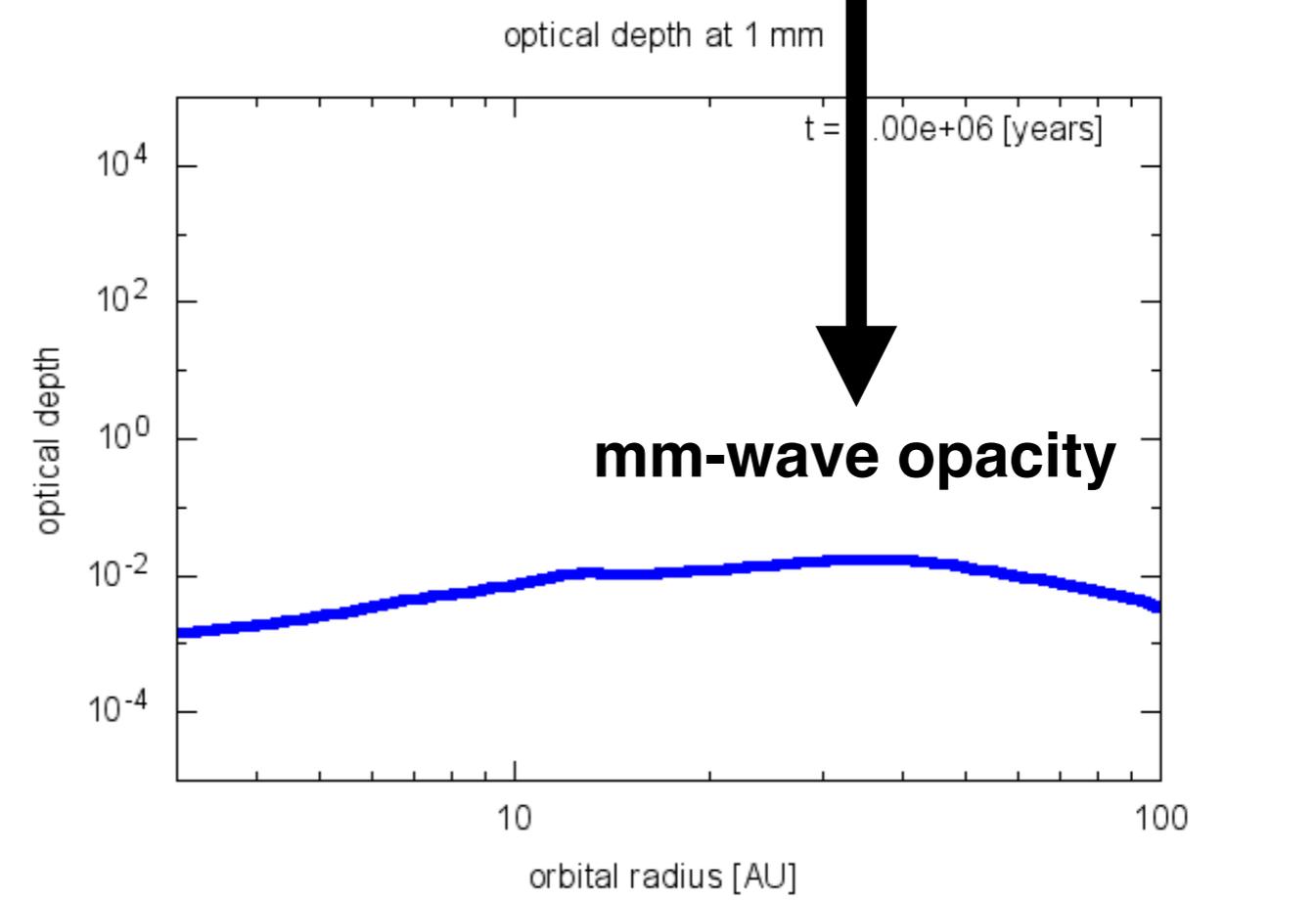
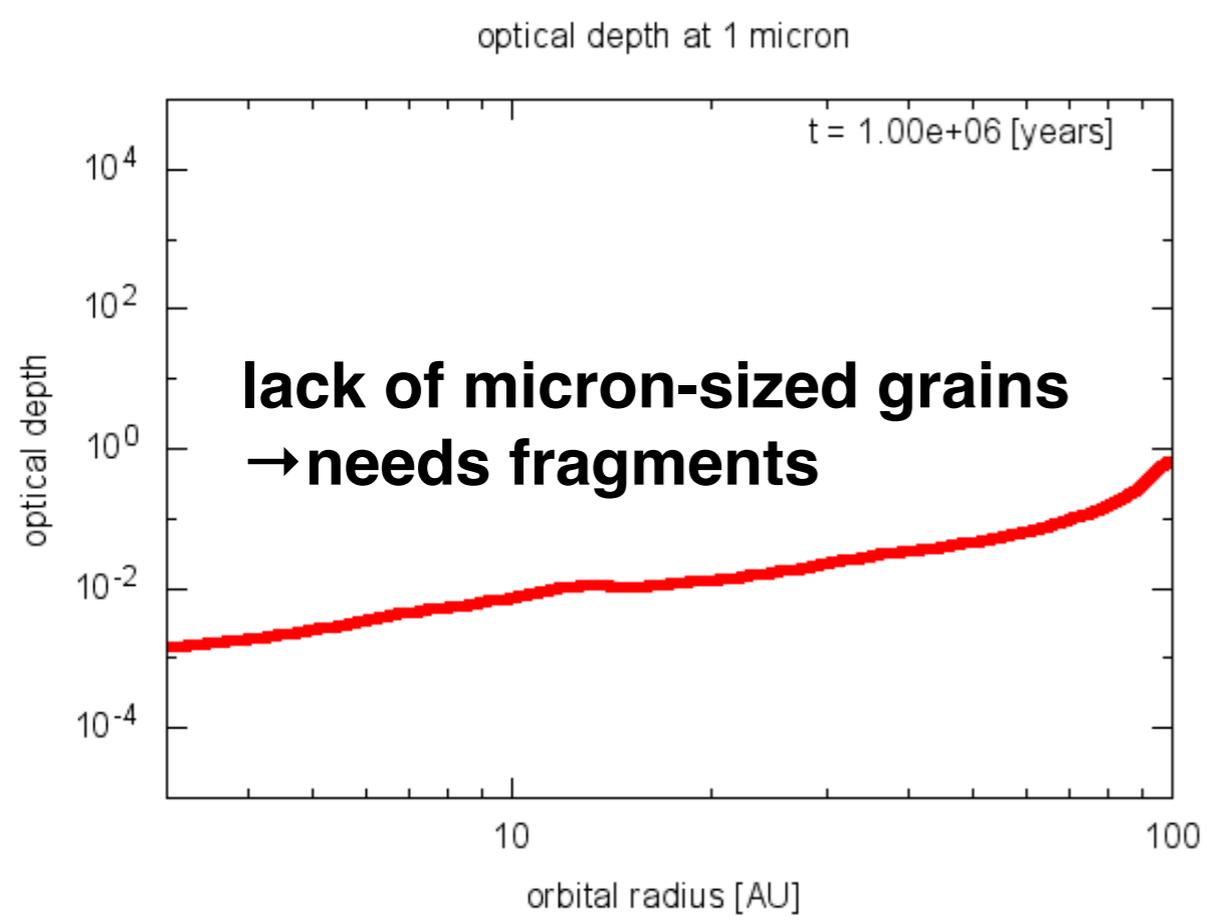
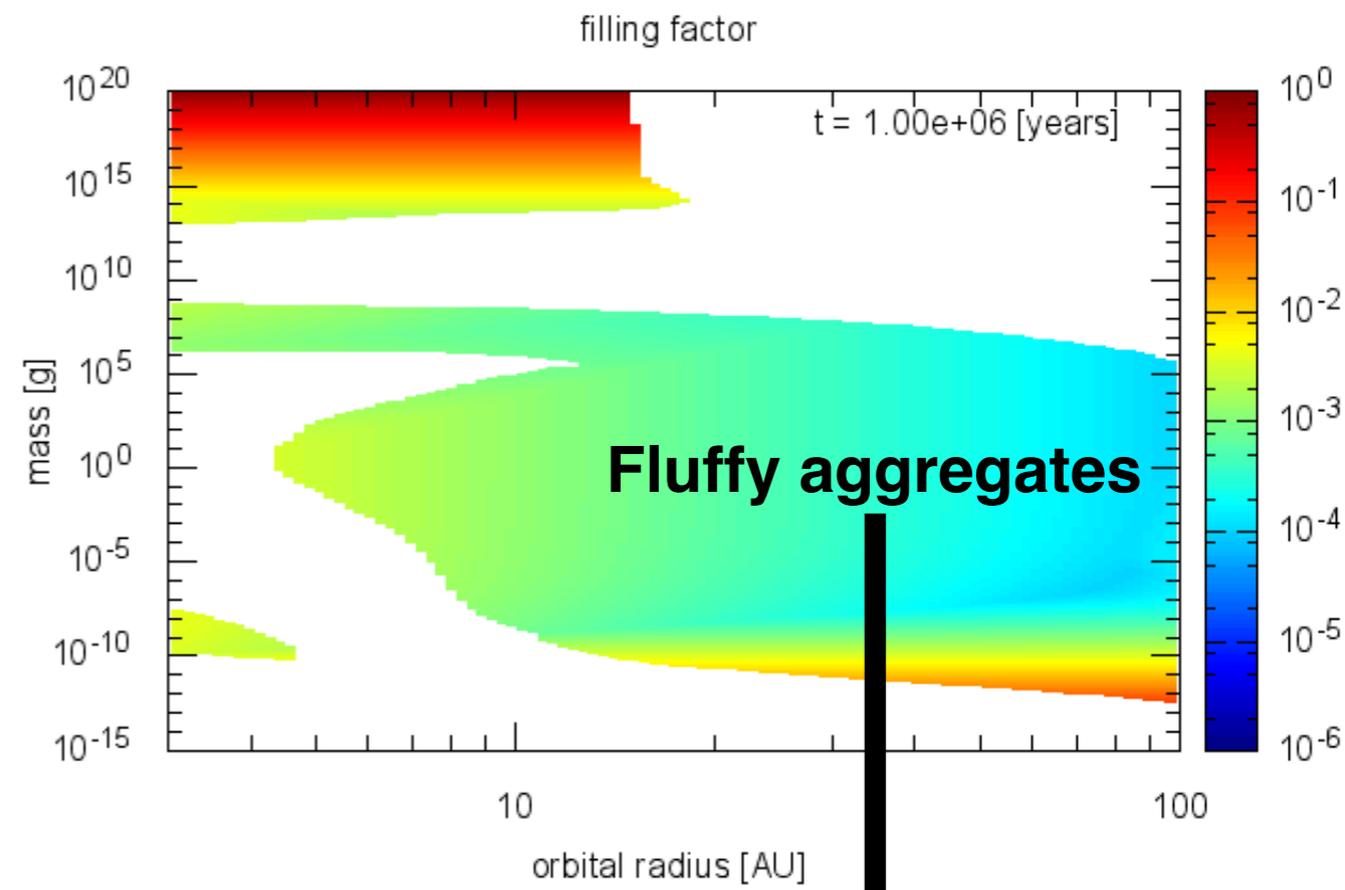
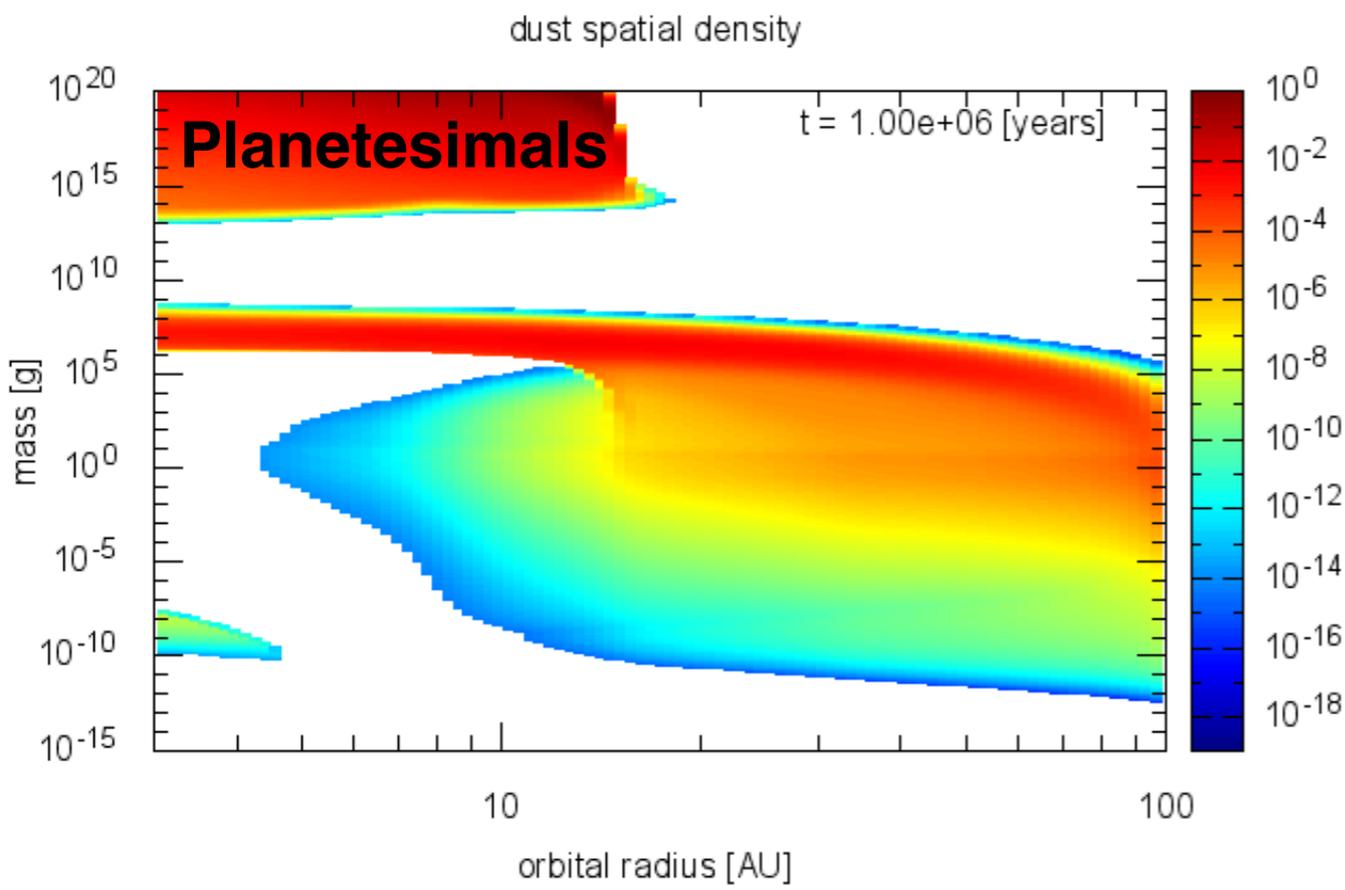


Optical depth at 1 micron



Optical depth at 1 mm





Conclusions

- We investigate the the static compressive strength of highly porous aggregates ($f < 0.1$)
- we reveal the overall porosity evolution
- the path avoids the three barriers: radial drift, fragmentation, and bouncing barriers

([Kataoka et al. 2013a](#), A&A, 554, A4, [Kataoka et al. 2013b](#), A&A, 557, L4)

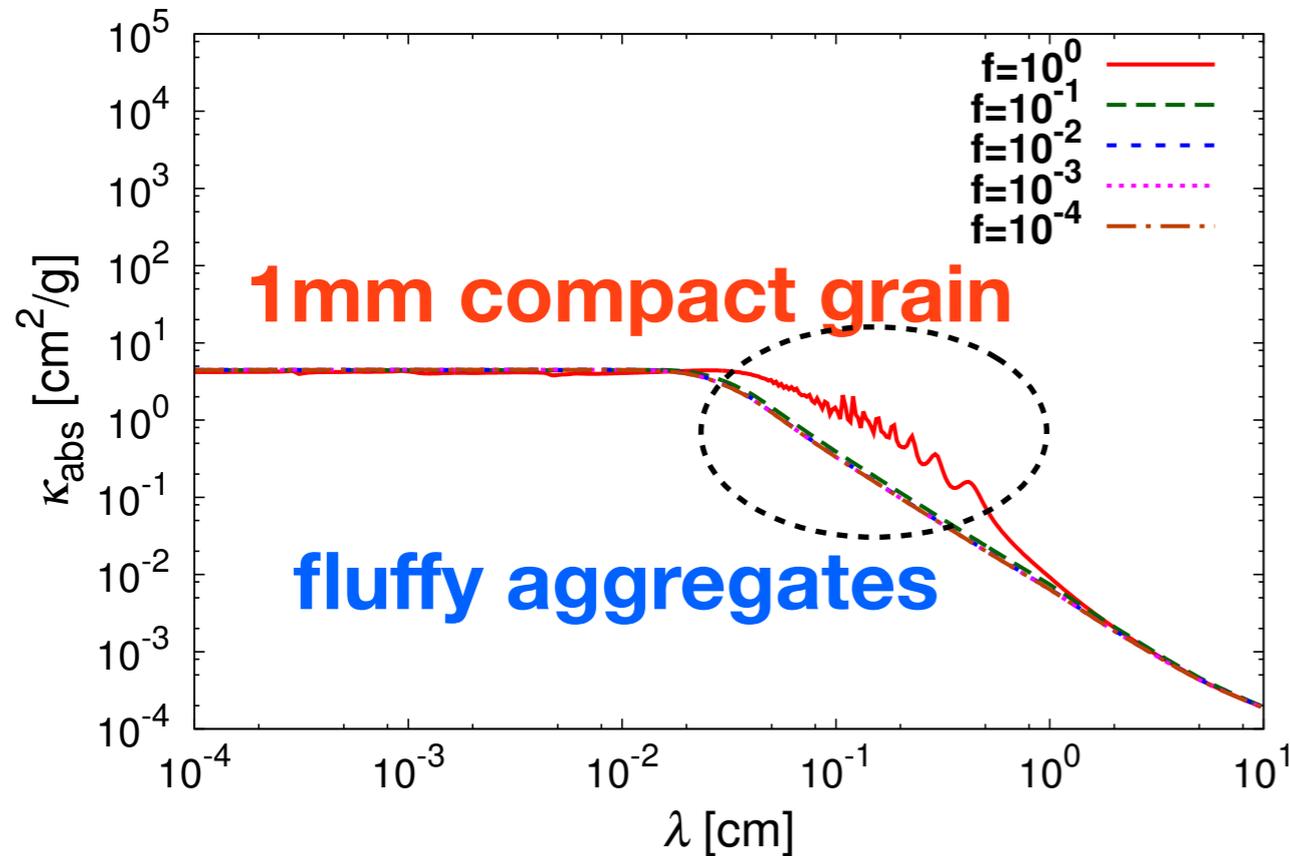
- Opacities are characterized by af

([Kataoka et al., 2014](#), A&A, 568, A42)

- Planetesimals are formed inside ~ 10 AU
- Fluffiness does not help for the radial drift barrier in outer disk.

How to observe fluffy aggregates?

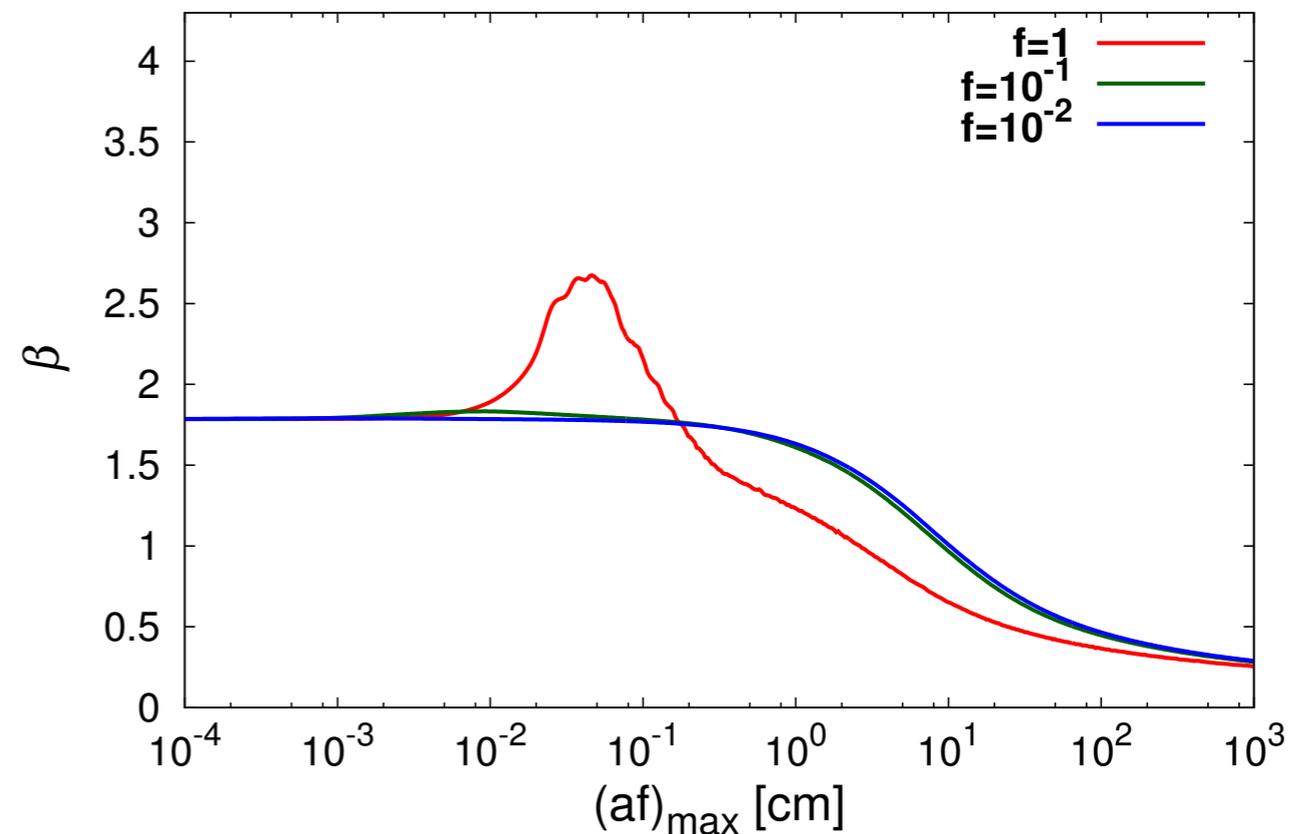
$a \times f = 1 \text{ mm}$



a: radius

f: filling factor

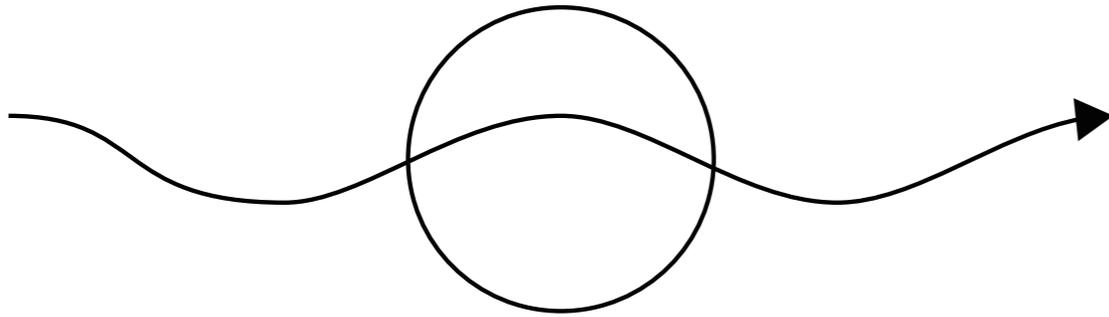
Opacity index β



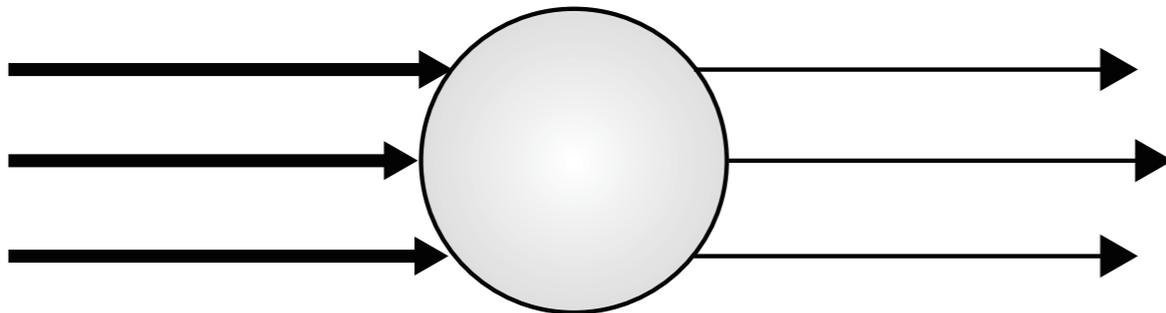
$\beta > 1.7$ only when grains are compact

Understanding opacity

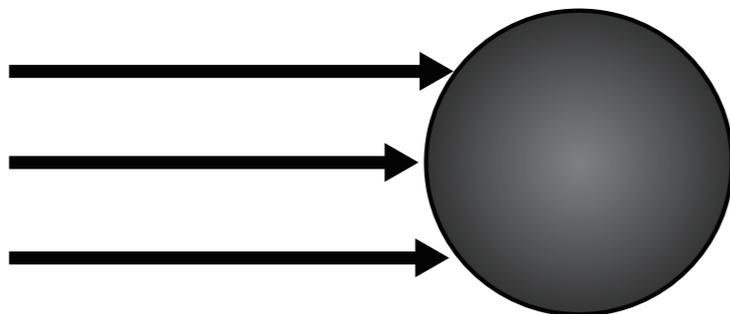
(a) $x < 1$



(b) $x > 1$, optically thin



(c) $x > 1$, optically thick



absorption/scattering opacity

$$\kappa_{\text{abs}} = \frac{\pi a^2}{m} Q_{\text{abs}},$$

$$\kappa_{\text{sca}} = \frac{\pi a^2}{m} Q_{\text{sca}}.$$

(m : grain mass)

→ $Q_{\text{abs}}, Q_{\text{sca}}$ is calculated with Mie theory

Useful parameters

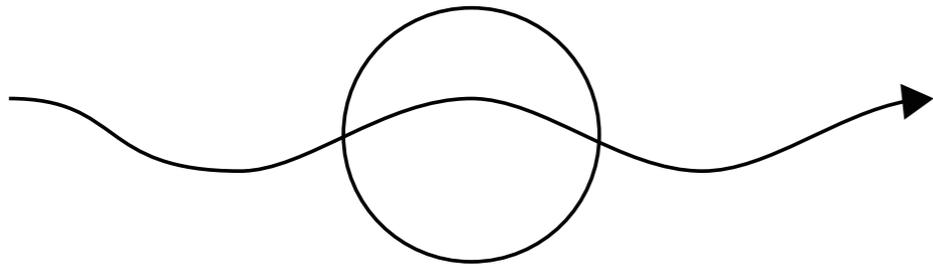
• size parameter $x \equiv \frac{2\pi a}{\lambda}$

• optical depth kx

→ derive the approximated equations

Understanding opacity : absorption

(a) $x < 1$



(a) $x \ll 1 \rightarrow$ Rayleigh regime

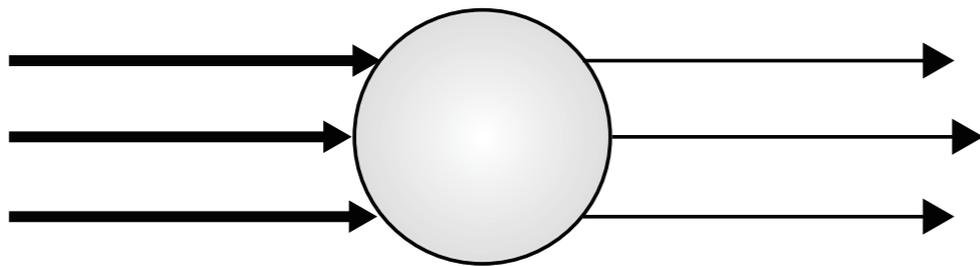
if $n-1 \ll 1, k \ll 1$

$$Q_{\text{abs}} \simeq Q_{\text{abs},1} \equiv \frac{24nkx}{(n^2 + 2)^2}$$

if $f \ll 1, n-1 \propto f, k \propto f$, thus

$$Q_{\text{abs},1} \propto af \text{ for } f \ll 1$$

(b) $x > 1$, optically thin

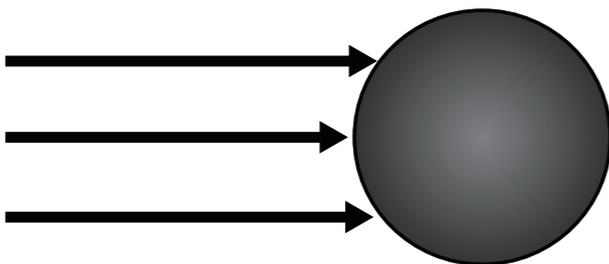


(b) $x \gg 1$ & $kx \ll 1$

$$Q_{\text{abs}} \simeq Q_{\text{abs},2} \equiv \frac{8kx}{3n} (n^3 - (n^2 - 1)^{3/2})$$

$$Q_{\text{abs},2} \propto af \text{ for } f \ll 1$$

(c) $x > 1$, optically thick



(c) $x \gg 1$ & $kx \gg 1 \rightarrow$ geometric regime

$$Q_{\text{abs}} \simeq Q_{\text{abs},3} \equiv \int_0^{\pi/2} (1 - R(\theta_i)) \sin 2\theta_i d\theta_i,$$

$$\approx 1 \text{ for } f \ll 1$$

1. K_{abs} is proportional to af
2. we derive the piecewise formula of abs. opacity

