Magnetic Reconnection Associated With Kink Modes: Three-Dimensional Full Kinetic Simulations

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In order to explain the explosive energy conversion by magnetic reconnection, the rate of reconnection has to be sufficiently large

$$E_{\text{rec}} \sim 0.1 \, V_A B_0,$$

and the fast reconnection has to persist for the time scale of the phenomena in space plasmas

a few minutes $\sim$ a few hours.
**Diffusion Region**

Magnetic dissipation takes place in the diffusion region where the MHD constraints break down. Magnetic field satisfies the diffusion equation

\[
\frac{\partial B}{\partial t} = \frac{\eta}{\mu_0} \nabla^2 B
\]

\(\eta\): Electric resistivity

It is still an open question how the resistivity is produced in collisionless plasmas.
Magnetic Dissipation in 2D System

It has been suggested that magnetic dissipation in 2D system is provided by the electron inertia resistivity due to the meandering/Speiser motion of the electrons.

\[ E_{rec} = \eta \frac{j}{n_e} \frac{1}{e} (\nabla \cdot P_e) + \frac{m_e dV_e}{e} \frac{d}{dt} \]

In order for the inertia resistivity to be effective, the electron diffusion region has to be compact,

\[ \frac{E_{rec}}{V_{A0} B_0} \sim \frac{\sigma}{l} \sim 0.1 \]

\[ \begin{align*}
\frac{\sigma}{\lambda e} \\
1 < \lambda_i
\end{align*} \]

which is believed to be guaranteed by the Hall effects [Birn et al., 2001].
It has been suggested that the elongated diffusion region is unstable in the current density direction, and anomalous resistivity could enhance the magnetic dissipation.
Large-Scale Particle-In-Cell Simulations in 3D System

Simulation code: 3D-EM-PIC + AMR (Adaptive Mesh Refinement)

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\[
\begin{align*}
B_x &= B_0 \tanh(z/\lambda) \\
J_y &= J_0 \operatorname{sech}^2(z/\lambda)
\end{align*}
\]

\[
\begin{align*}
L_x \sqcup L_y \sqcup L_z &= 31 \lambda_i \sqcup 7.7 \lambda_i \sqcup 31 \lambda_i \\
N_x \sqcup N_y \sqcup N_z &= 1024 \sqcup 256 \sqcup 1024 \\
m_i/m_e &= 25 \\
The \text{number of particles} &< 1.5 \times 10^9
\end{align*}
\]
In 3D system a kink mode is driven by the ions drifting in the cross-field direction. The kink mode deforms the current sheet structure drastically.

Nevertheless, the time profile of the reconnection electric field is almost identical to the run without the kink mode, and the fast reconnection is achieved.

$$E_{\text{rec}} \sim 0.1 V_{A0}B_0$$
Generalized Ohm’s Law

How is the fast reconnection achieved in the system with kink mode?

\[
\frac{1}{\langle n_e \rangle} \left( \langle n_e \bar{V}_e \rangle \times \langle \bar{B} \rangle \right)_{y'} + \frac{1}{e \langle n_e \rangle} \langle (\nabla' \cdot \bar{P}_e) y' \rangle + \frac{m_c}{e \langle n_e \rangle} \left( \langle dV_{ey'} \rangle \right) + \frac{1}{\langle n_e \rangle} \langle \delta n_e \delta E_{y'} \rangle
\]
2D Simulations in the YZ Plane

In order to clarify the dissipation mechanism associated with the kink mode, we perform 2D simulations in the system where the inertia effects are excluded. It is confirmed that magnetic dissipation is provided, with an amplitude independent of the mass ratio.

[Ozaki et al., 1996; Horiuchi and Sato, 1999; Shinohara et al., 2001]
Plasma Heating due to Wave-Particle Interaction

It is found that the electrons are strongly heated in the thin current sheet. We also find a small-scale kink structure embedded in the large-scale kinked current sheet.

Hybrid-scale kink mode: \( \lambda \sim (\lambda_i \lambda_e)^{1/2} \)

[Shinohara et al., 2001; Daughton, 2003]

It is the hybrid-scale kink mode that is responsible for the electron heating.
Magnetic Dissipation Associated with Kink Modes

Under a steady-state condition, the energy gain should be balanced with the energy loss.

\[ \vec{j} \cdot \vec{E} \sim \frac{1}{2} n_e m_e V_{ec}^2 \nu_e + \frac{1}{2} n_i m_i V_{ic}^2 \nu_i \]

\( \nu_i, \nu_e : \) Effective collision frequency

If the effective collision converts the bulk energy to the thermal energy,

\[ \vec{j} \cdot \vec{E} \sim \frac{\partial (nT)}{\partial t} \]

should be satisfied.
Magnetic Dissipation in 3D System

\[ \left\langle -E'_{y'} \right\rangle_{xline} = \eta^{eff}_{xline} \left\langle j_{y'} \right\rangle_{xline} = \frac{1}{\langle n_e \rangle} \left( \langle n_e \bar{V}_e \rangle \times \langle B \rangle \right)_{y'} + \frac{1}{e \langle n_e \rangle} \left\langle (\nabla' \cdot \bar{P}_e)_{y'} \right\rangle + \frac{m_e}{e \langle n_e \rangle} \left\langle \frac{dV_{ey'}}{dt} \right\rangle + \frac{1}{\langle n_e \rangle} \left\langle \delta n_e \delta E'_{y'} \right\rangle + \eta^{an}_{y'} \left\langle j_{y'} \right\rangle_{xline} \]

Anomalous resistivity

\[ \eta^{an} = \left\langle -E'^{(2D)}_{y'} \right\rangle / \left\langle j_{y'}^{(2D)} \right\rangle \]

Electron inertial effects

Due to anomalous resistivity

It is found that the anomalous resistivity enhances the magnetic dissipation, compensating the reduction in the inertia resistivity, and makes a significant contribution for realizing the fast reconnection.
Summary and Conclusions

We have performed a 3D PIC simulation in a large system in order to investigate how kink modes affect magnetic reconnection.

We found that the kink mode broadens the width of the current sheet and decreases the inertia resistivity. Instead, the anomalous resistivity associated with the kink mode compensates the depletion so as to keep the high reconnection rate.

The present result suggests that the electron dynamics in the electron diffusion region is automatically adjusted so as to produce sufficient dissipation for the fast magnetic reconnection.