

超新星放出ガス中のダスト形成 と衝撃波によるダストの破壊

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Outline

1. Introduction

- Properties of interstellar dust in the Galaxy
- Observations of dust at high redshift

2. Formation and evolution of dust in Pop III

Type II-P and pair-instability SNe

3. Formation of dust in various types of SNe

4. Implication and Summary

1. Introduction

1-1-1. Cosmic dust

○ **Cosmic dust** : solid particles with radii of a few nm to ~0.1 mm in astronomical environments
interplanetary dust, interstellar dust, intergalactic dust ...

Milky Way (optical)



Milky Way (infrared)



Dust grains absorb UV/optical lights and reemit it by their thermal radiation at IR wavelengths!

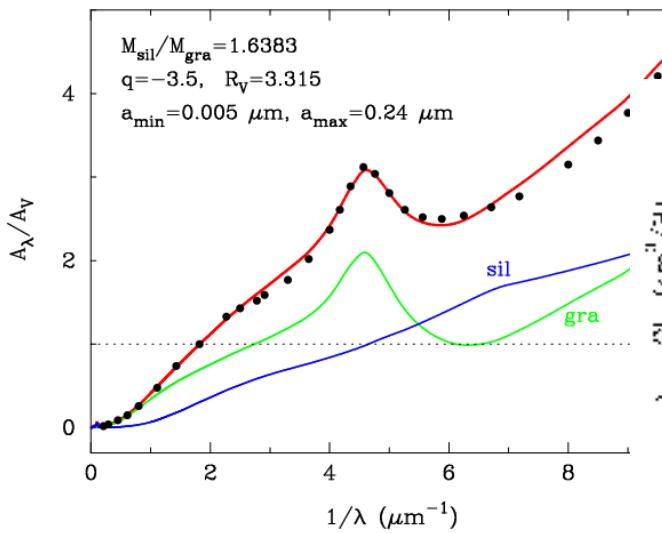
1-1-2. Interstellar dust in our Galaxy

○ Dust in our Galaxy → when and where is dust formed?

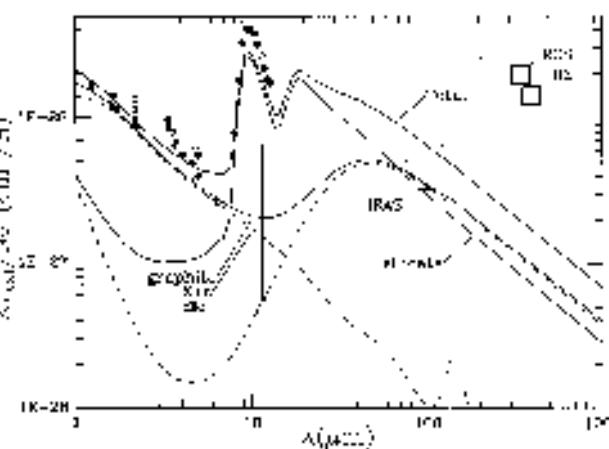
- composition : **graphite** (carbonaceous) grains
silicate (SiO_2 , MgSiO_3 , MgFeSiO_4 , ...) grains
 - size : $n(a) = f(a)da = a^{-3.5} da$ (**$0.005\text{--}0.25 \mu\text{m}$**)
 - amount : $M_{\text{dust}} / M_{\text{gas}} \sim 1 / 140$ ($\sim 10^9 \text{ M}_{\odot}$)

→ MRN dust model (e.g., Mathis+’77; Draine & Lee’84)

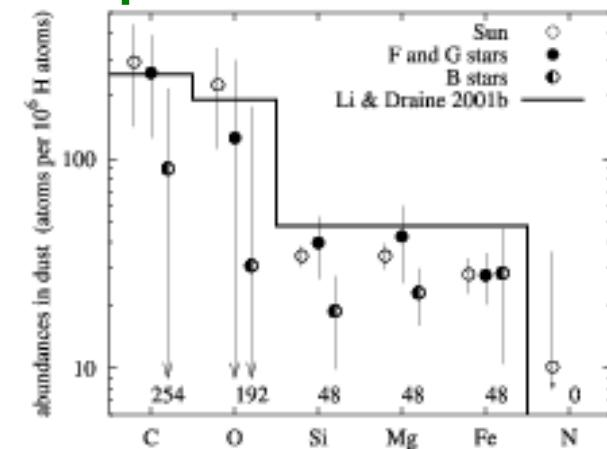
extinction curve

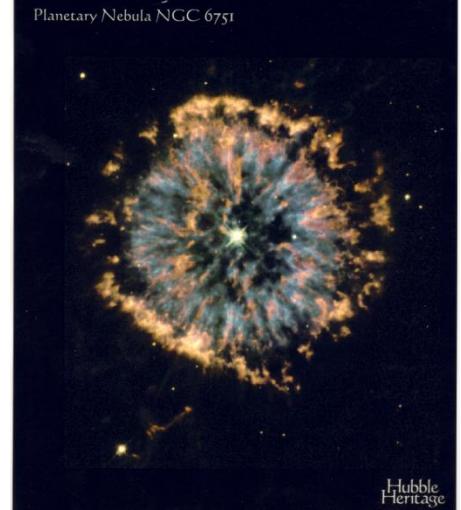


IR spectral feature



depletion of elements

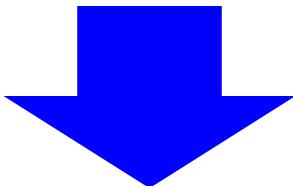


Hubble
Heritage

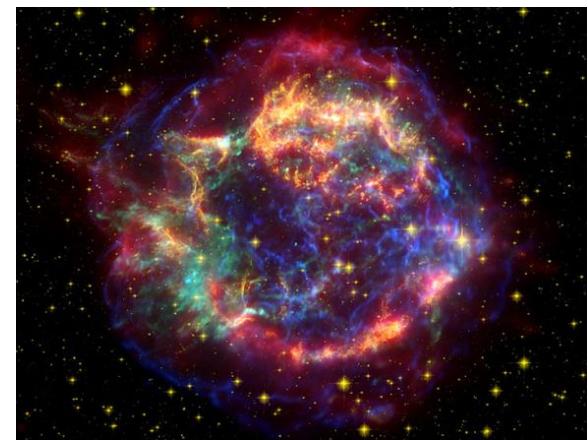
1-1-3. Formation site of dust

O Formation sites of dust

- abundant metal (metal : N > 5)
- low gas temperature ($T < \sim 2000$ K)
- high gas density ($n > \sim 10^6$ cm $^{-3}$)



- mass-loss winds of AGB stars
- expanding ejecta of supernovae
- **molecular clouds (grain growth only)**
- **red giant, W-R stars, novae, protoplanetary disk ...**



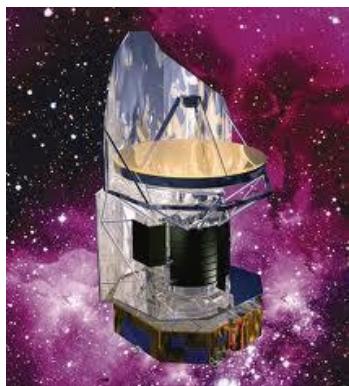
→ relative contribution of these sources is unclear!

1-2-1. Discovery of massive dust at $z > 5$

- The submm observations have confirmed the presence of dust in excess of 10^8 M_{\odot} in 30% of $z > 5$ quasars

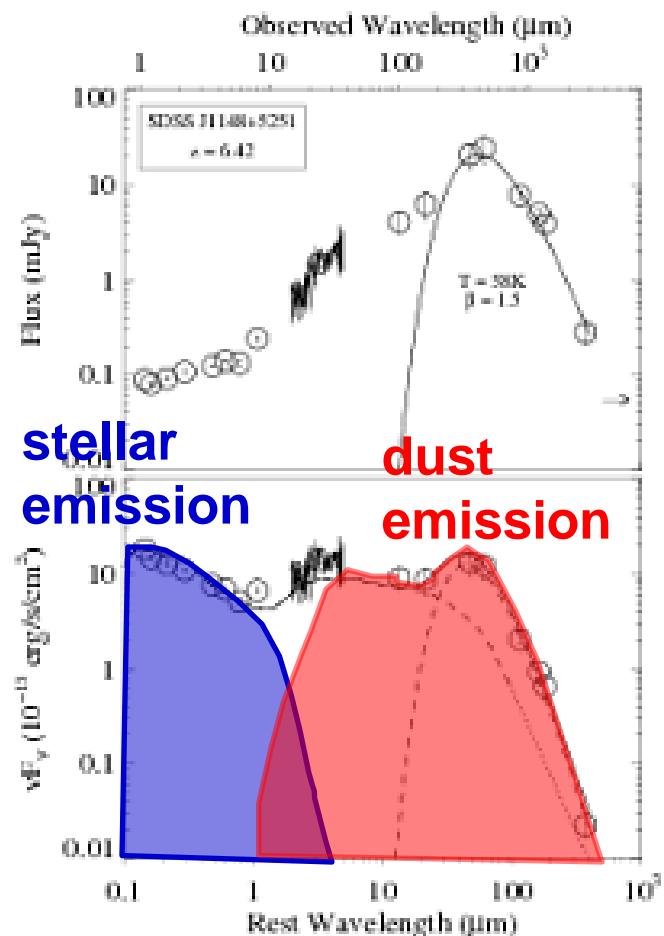
SDSS J1148+5251 at $z=6.4$

- age : 890 Myr
- IR luminosity : $\sim(1-3)\times10^{13} \text{ L}_{\odot}$
- dust mass : $(2-7)\times10^8 \text{ M}_{\odot}$
- SFR : $\sim3000 \text{ M}_{\odot}/\text{yr}$ (Salpeter IMF)
- gas mass : $\sim3\times10^{10} \text{ M}_{\odot}$ (Walter+'04)
- metallicity : \sim solar



***Herschel (3.5m)**

- PACS :
70, 100, 160 μm
- SPIRE
250, 350, 500 μm



1-2-2. What are dust sources in high-z quasar?

- Supernovae (Type II SNe)

→ ~0.1 Msun per SN is sufficient

(Morgan & Edmunds'03; Maiolino+'06; Li+'08)

→ > 1 Msun per SN (Dwek+'07)

- AGB stars + SNe

(Valiante+'09; Gall+'11; Dwek & Cherchneff'11)

→ 0.01-0.05 Msun per AGB (Zhukovska & Gail '08)

→ 0.01-1 Msun per SN

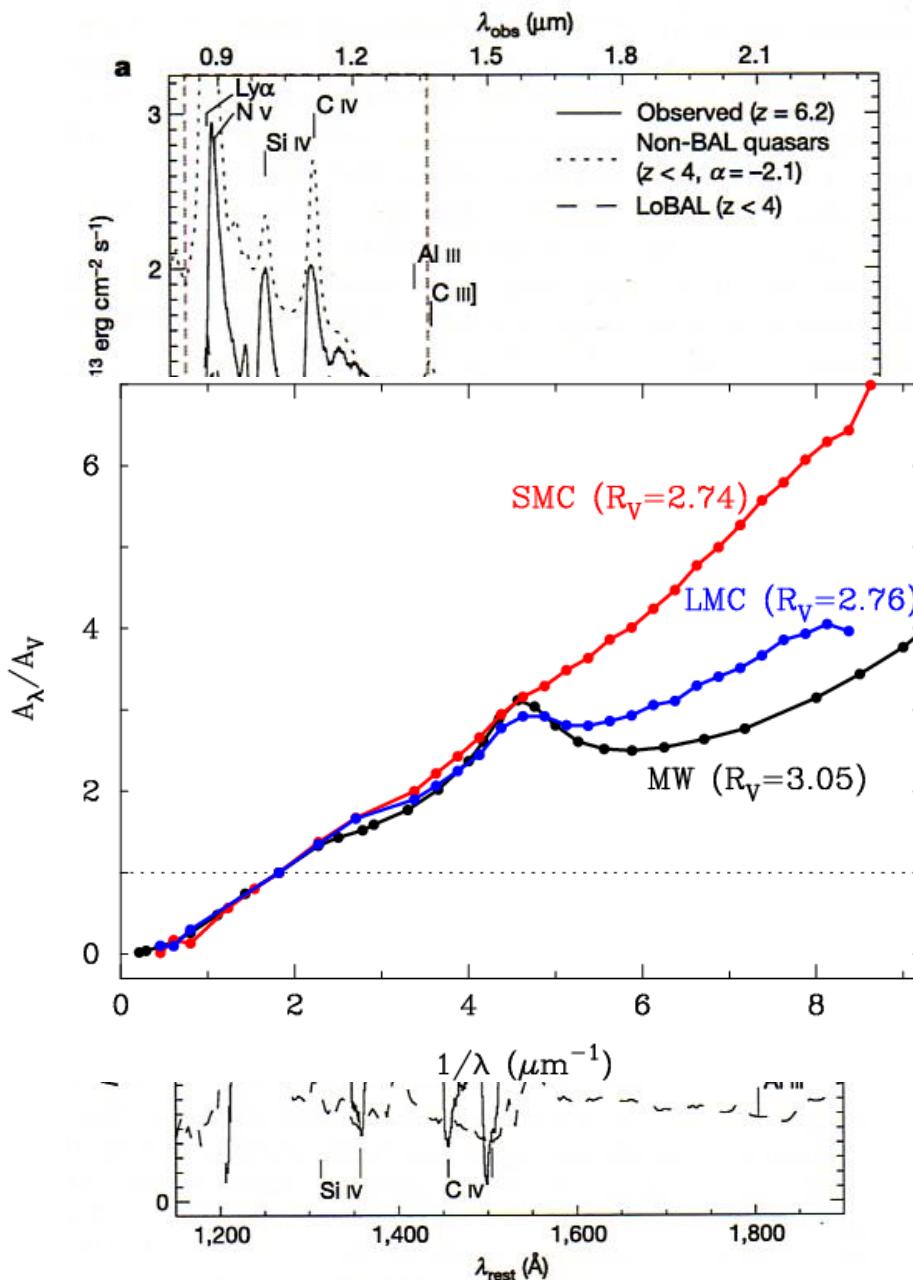
- Grain growth in dense clouds + AGB stars + SNe

(Draine'09; Michalowski+'10; Pipino+'11; Mattsson'11)

→ Tgrowth ~ $10^7 (Z / Z_{\text{sun}})$ yr

- Quasar outflows (Elvis+'02)

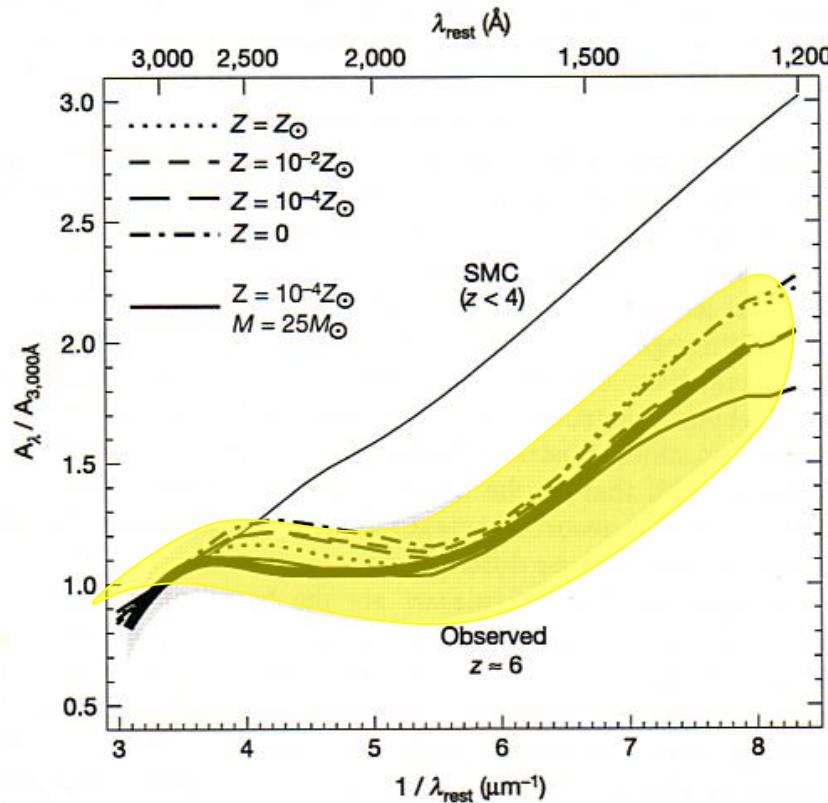
1-3-1. Extinction curves at high-z quasars



Maiolino+’04, Nature, 431, 533

SDSS J1048+4637 at $z=6.2$

Broad absorption line (BAL) quasars



different dust properties
from those at low redshift

1-3-2. Extinction curves at $3.9 < z < 6.4$

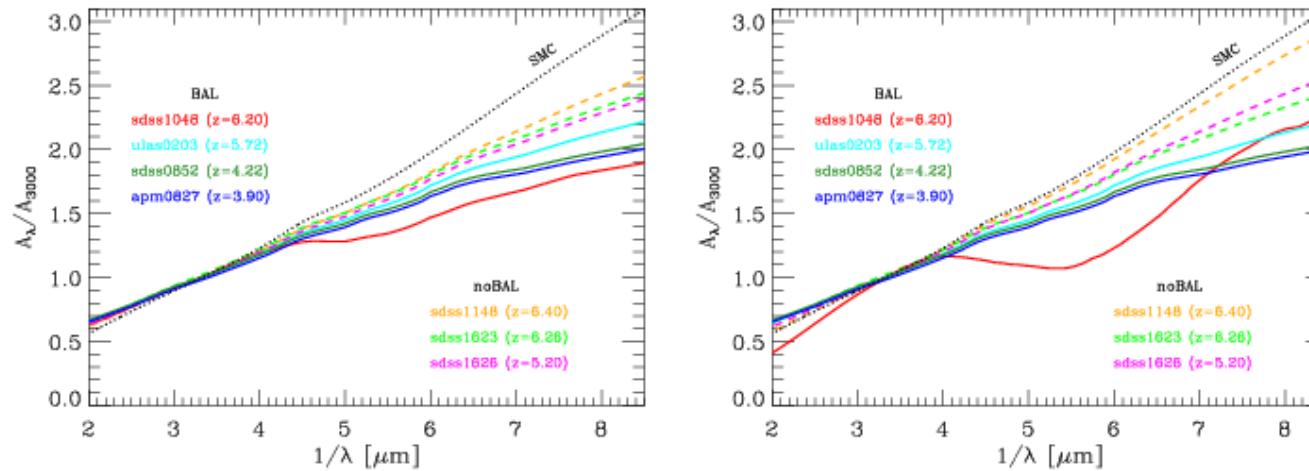


Fig. 4. Best fit extinction curves of reddened quasars. The solid lines are for BAL quasars, while dashed lines are for non-BAL quasars. For comparison the SMC extinction curve is also shown and labeled in the Figure (dotted black line). The panel on the left shows the results assuming a minimum intrinsic slope $\alpha_{\lambda,\min} = -2.9$, while the panel on the right is obtained with $\alpha_{\lambda,\min} = -2.6$.

Gallerani+’10, A&A,
523, 85

7 of 33 requires
substantial dust
extinction,
which deviates
from the SMC

The mean
Extinction
curves for
BAL quasars
deviates from
the SMC with
level > 95 %

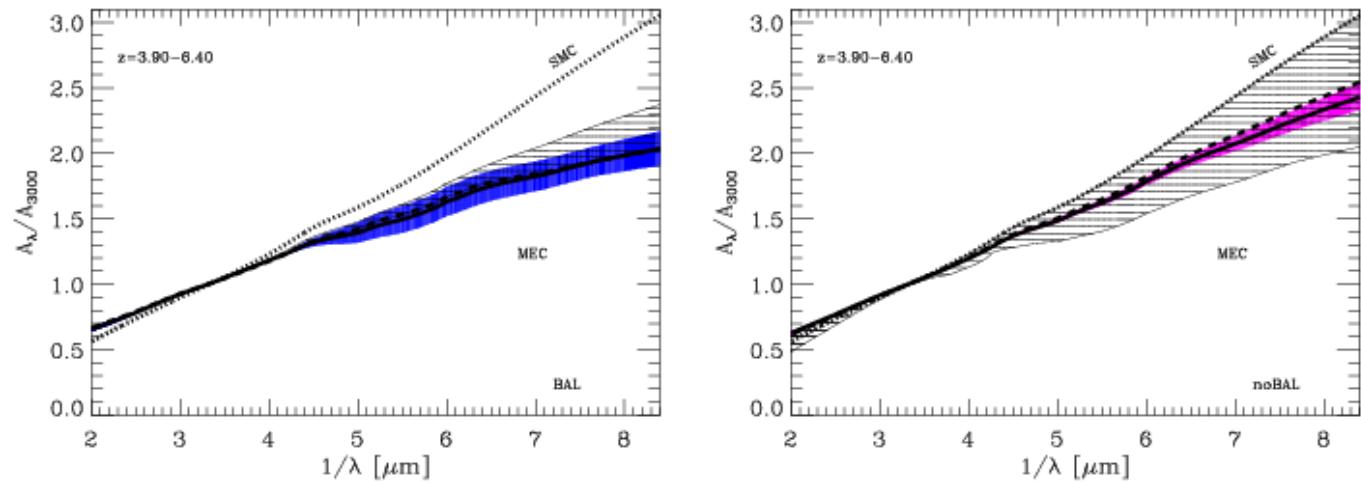
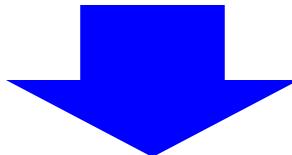


Fig. 6. Mean (MEC) and simultaneous/global (GEC) extinction curve of reddened quasars divided into BAL (left) and non-BAL (right). The coding is the same as in Fig. 5.

1-4. Summary of Introduction

- There is clear evidence for huge amounts of dust at $z > 4$, but the dust sources remain unexplained
 - SNe? AGB stars? grain growth in the dense clouds? quasar outflow? any other sources?
- Properties (composition & size) of dust at high z are likely to be different from those at low z
 - high- z quasars and GRBs are good targets to probe the extinction curves in their host galaxies



At $z > 4$, short-lived SNe II ($M = 8\text{-}40 \text{ M}_{\odot}$) dominate the dust production over AGB stars ($M < 8 \text{ M}_{\odot}$) ??

1-5. Aim of our study

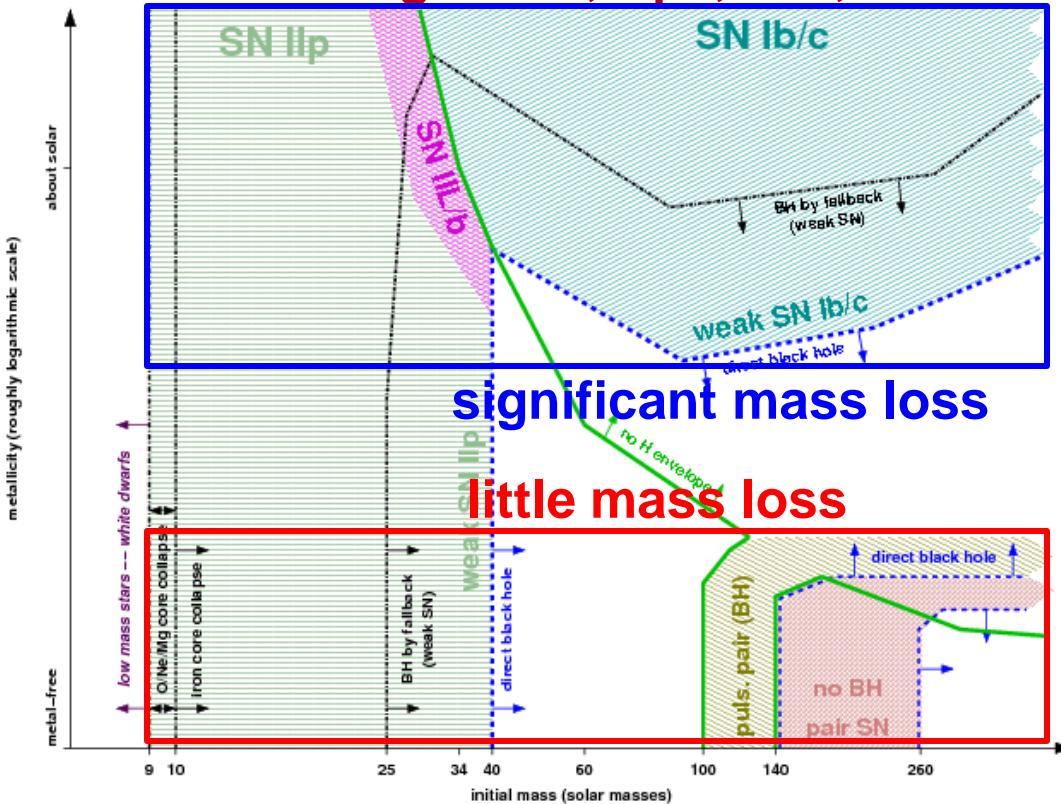
- Dust absorbs stellar light and emits it by thermal radiation
 - plays a crucial role in interpreting the evolution history of the universe from high-z observations
- Dust has great impacts on formation processes of stars
 - forming molecules (mainly H₂) on the surface
(e.g., Cozaux & Spaans'04)
 - providing additional cooling pathways of gas through thermal emission
(e.g., Omukai+'05; Schneider+'06)
 - controlling the energy balance in interstellar space

We aim at revealing the evolution of composition, size, and amount of dust, considering formation and destruction processes of dust consistently!!

2. Formation and evolution of dust in Pop III SNe II-P and PISNe

2-0. Death of single massive stars

Heger+’03, ApJ, 591, 288



At low metallicity ($Z < 10^{-4} Z_{\text{sun}}$)

▪ Type II-P SNe:

$M_{\text{ZAMS}} = 8-40 \text{ Msun}$

▪ pair-instability SNe:

$M_{\text{ZAMS}} = 140-260 \text{ Msun}$

At high metallicity

▪ Type II-P SNe:

$M_{\text{ZAMS}} = 8-25 \text{ Msun?}$
massive H envelope

▪ Type IIb SNe:

$M_{\text{ZAMS}} = 25-35 \text{ Msun?}$
very thin H-envelope

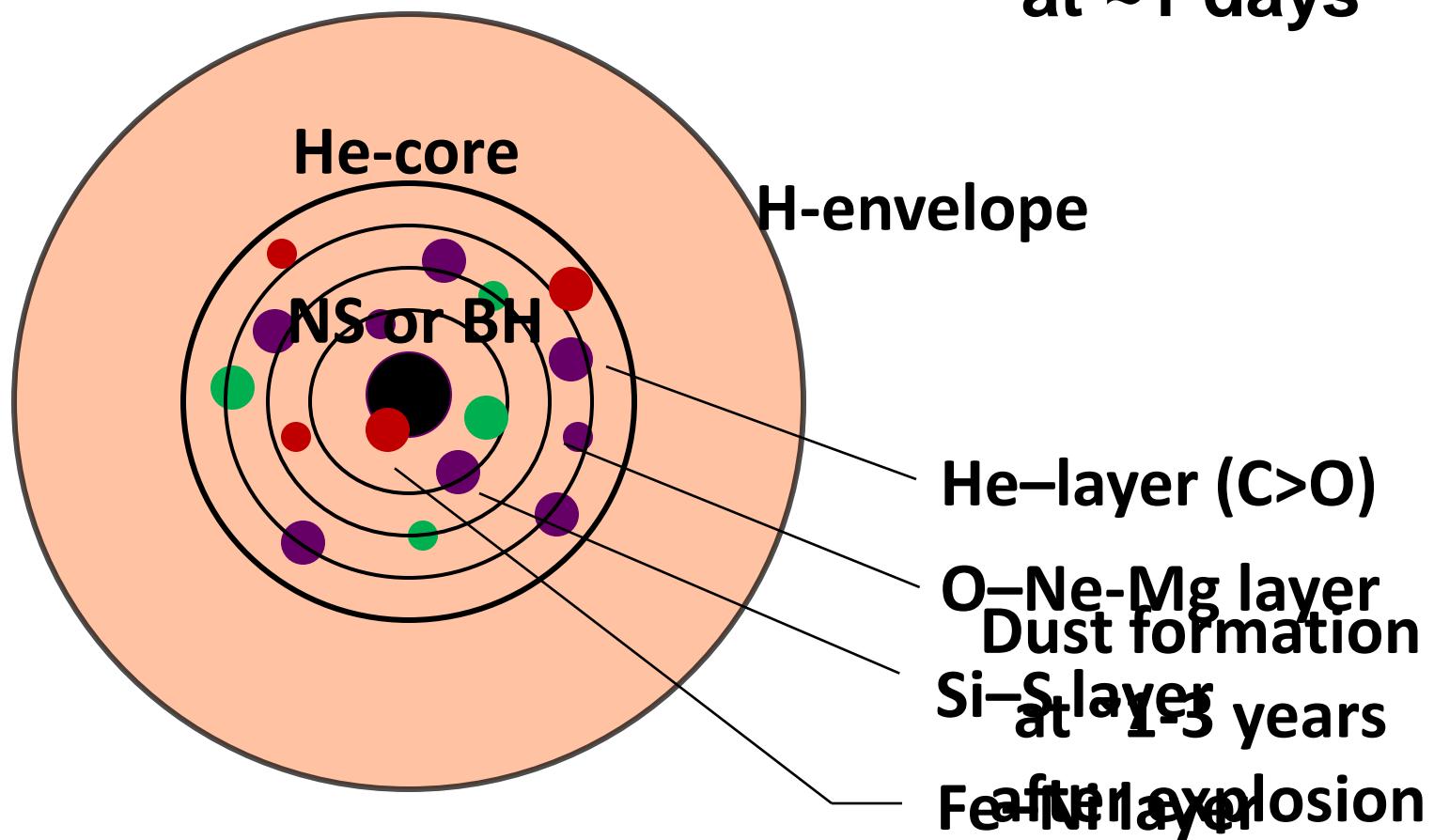
▪ Type Ib/Ic SNe :

$M_{\text{ZAMS}} > 35 \text{ Msun?}$
no H / He envelope

▪ Type Ia SNe :

thermonuclear explosion
of C+O white dwarfs
 $M_{\text{pre-explosion}} \sim 1.4 \text{ Msun}$

2-1. Dust Formation in Pop III SNe

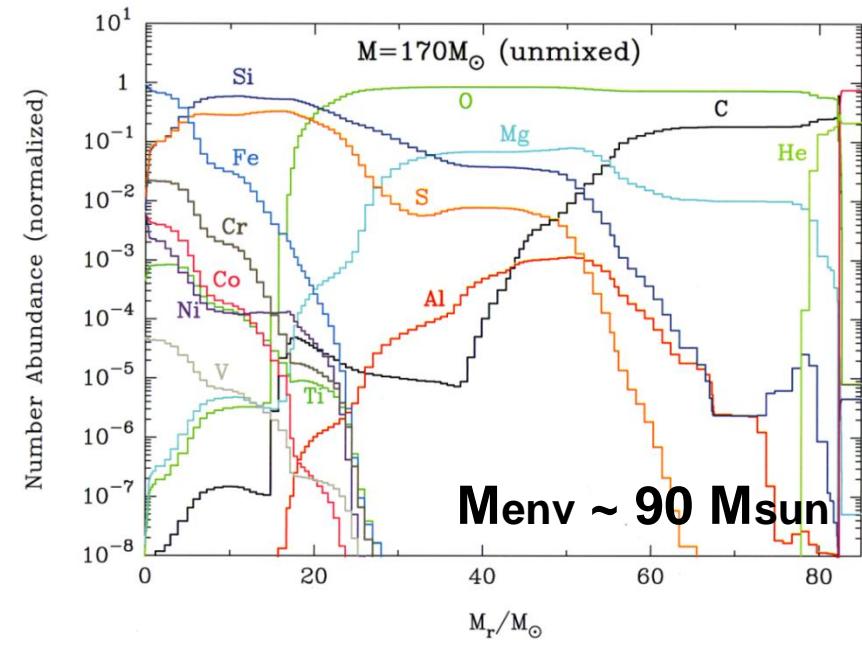
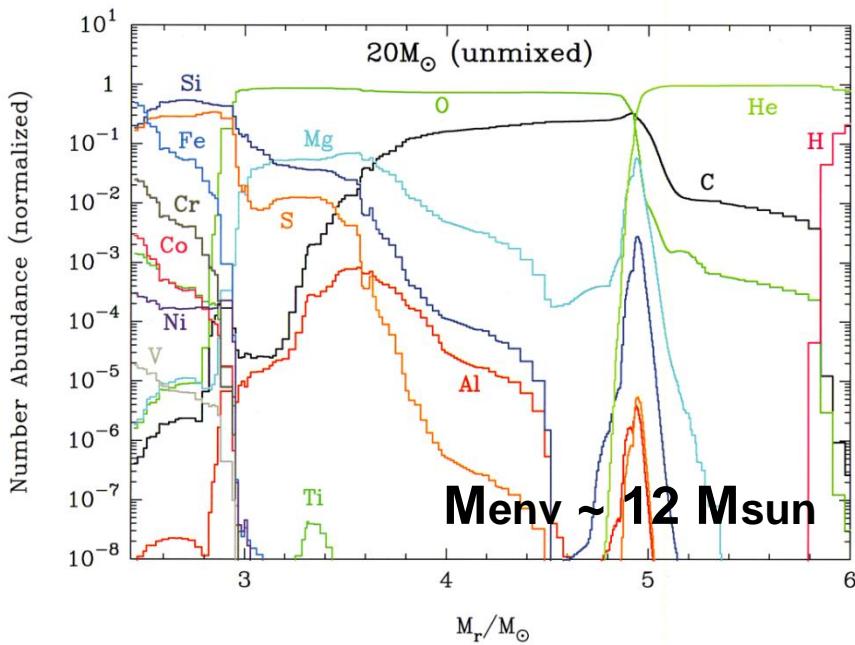


2-1-1. Dust formation in primordial SNe

Nozawa+’03, ApJ, 598, 785

○ Population III SNe model (Umeda & Nomoto’02)

- SNe II-P : $M_{\text{ZAMS}} = 13, 20, 25, 30 \text{ Msun}$ ($E_{51}=1$)
- PISNe : $M_{\text{ZAMS}} = 170 \text{ Msun}$ ($E_{51}=20$), 200 Msun ($E_{51}=28$)



- nucleation and grain growth theory (Kozasa & Hasegawa’88)
- no mixing of elements within the He-core
- complete formation of CO and SiO

2-1-1.5. Calculations of dust formation

- nucleation and grain growth theory taking account of chemical reaction at condensation

(Kozasa & Hasegawa'87)

- key species:
gas species with the least collision frequency among reactants



key species controls the kinetics of the nucleation and grain growth

Dust species	Chemical reactions
$\text{Fe}_{(\text{s})}$	$\text{Fe}_{(\text{g})} \rightarrow \text{Fe}_{(\text{s})}$
$\text{FeS}_{(\text{s})}$	$\text{Fe}_{(\text{g})} + \text{S}_{(\text{g})} \rightarrow \text{FeS}_{(\text{s})}$
$\text{Si}_{(\text{s})}$	$\text{Si}_{(\text{g})} \rightarrow \text{Si}_{(\text{s})}$
$\text{Ti}_{(\text{s})}$	$\text{Ti}_{(\text{g})} \rightarrow \text{Ti}_{(\text{s})}$
$\text{V}_{(\text{s})}$	$\text{V}_{(\text{g})} \rightarrow \text{V}_{(\text{s})}$
$\text{Cr}_{(\text{s})}$	$\text{Cr}_{(\text{g})} \rightarrow \text{Cr}_{(\text{s})}$
$\text{Co}_{(\text{s})}$	$\text{Co}_{(\text{g})} \rightarrow \text{Co}_{(\text{s})}$
$\text{Ni}_{(\text{s})}$	$\text{Ni}_{(\text{g})} \rightarrow \text{Ni}_{(\text{s})}$
$\text{Cu}_{(\text{s})}$	$\text{Cu}_{(\text{g})} \rightarrow \text{Cu}_{(\text{s})}$
$\text{C}_{(\text{s})}$	$\text{C}_{(\text{g})} \rightarrow \text{C}_{(\text{s})}$
$\text{SiC}_{(\text{s})}$	$\text{Si}_{(\text{g})} + \text{C}_{(\text{g})} \rightarrow \text{SiC}_{(\text{s})}$
$\text{TiC}_{(\text{s})}$	$\text{Ti}_{(\text{g})} + \text{C}_{(\text{g})} \rightarrow \text{TiC}_{(\text{s})}$
$\text{Al}_2\text{O}_3_{(\text{s})}$	$2\text{Al}_{(\text{g})} + 3\text{O}_{(\text{g})} \rightarrow \text{Al}_2\text{O}_3_{(\text{s})}$
$\text{MgSiO}_3_{(\text{s})}$	$\text{Mg}_{(\text{g})} + \text{SiO}_{(\text{g})} + 2\text{O}_{(\text{g})} \rightarrow \text{MgSiO}_3_{(\text{s})}$
$\text{Mg}_2\text{SiO}_4_{(\text{s})}$	$2\text{Mg}_{(\text{g})} + \text{SiO}_{(\text{g})} + 3\text{O}_{(\text{g})} \rightarrow \text{Mg}_2\text{SiO}_4_{(\text{s})}$
$\text{SiO}_2_{(\text{s})}$	$\text{SiO}_{(\text{g})} + \text{O}_{(\text{g})} \rightarrow \text{SiO}_2_{(\text{s})}$
$\text{MgO}_{(\text{s})}$	$\text{Mg}_{(\text{g})} + \text{O}_{(\text{g})} \rightarrow \text{MgO}_{(\text{s})}$
$\text{Fe}_3\text{O}_4_{(\text{s})}$	$3\text{Fe}_{(\text{g})} + 4\text{O}_{(\text{g})} \rightarrow \text{Fe}_3\text{O}_4_{(\text{s})}$
$\text{FeO}_{(\text{s})}$	$\text{Fe}_{(\text{g})} + \text{O}_{(\text{g})} \rightarrow \text{FeO}_{(\text{s})}$

2-1-2. Nucleation rate of dust

Steady-state classical nucleation rate

$$J_s(t) = \alpha_s \Omega \left(\frac{2\sigma}{\pi m_1} \right)^{\frac{1}{2}} \Pi c_1^2(t) \exp \left[-\frac{4}{27} \frac{\mu^3}{(\ln S)^2} \right]$$

$$J_s(t) = 4\pi a_0^2 \alpha_s \left(\frac{kT}{2\pi m_1} \right)^{\frac{1}{2}} c_1^2(t) \frac{1}{3} \left(\frac{\mu}{\pi} \right)^{\frac{1}{2}} \exp \left[-\frac{4}{27} \frac{\mu^3}{(\ln S)^2} \right]$$

$$P_{i,\text{eq}} \qquad kT \qquad \underbrace{\qquad}_{i}$$

α_s : sticking probability of key species ($\alpha_s = 1$, in the calculations)

Ω : volume of the condensate per key species ($\Omega = 4\pi a_0^3/3$)

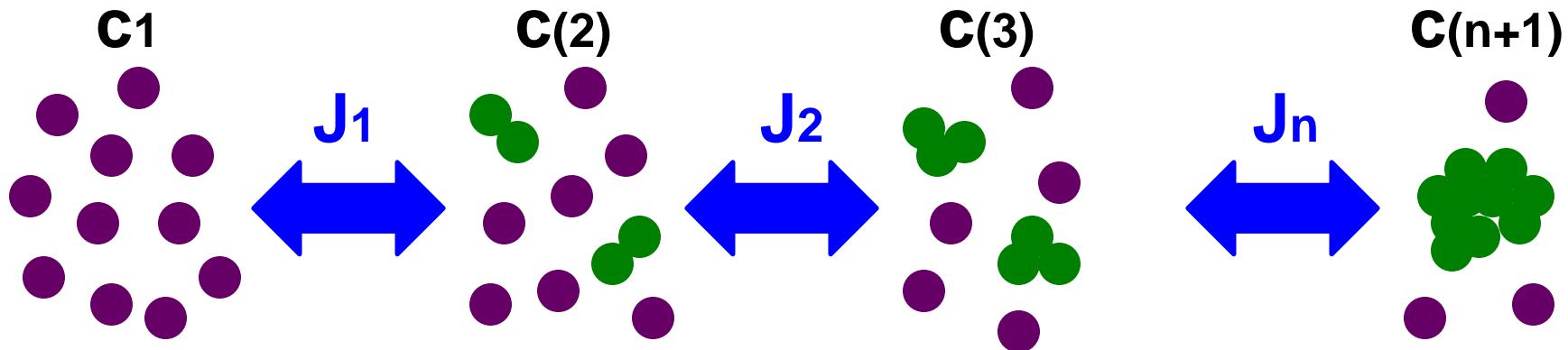
σ : surface energy of the condensate

m_1 : mass of key species

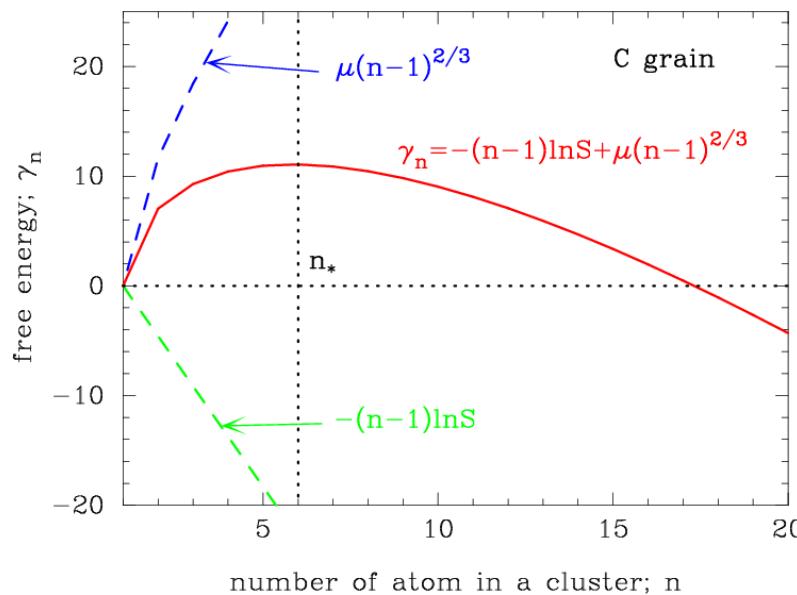
$c_1(t)$: number density of key species

μ : $\mu \equiv 4\pi a_0^2 \sigma / kT$; energy barrier for nucleation

2-1-2.5. Steady-state nucleation rate



$$J_s = J_1 = J_2 = \dots = J_n = \left\{ \frac{1}{(\alpha_1 c_1)} + \sum \frac{1}{(\alpha_n c_1^2)} \exp[g(n)] \right\}^{-1}$$



$$g(n) = -(n-1) \ln S + \mu(n-1)^{\frac{2}{3}}$$

$$(n_* - 1)^{\frac{1}{3}} = \frac{2\mu}{3 \ln S}$$

2-1-3. Basic equations of dust formation

Equation of conservation for key species

$$1 - \frac{c_1(t)}{\tilde{c}_1(t)} = \int_{t_0}^t \frac{J(t')}{\tilde{c}_1(t')} \frac{4\pi}{3\Omega} r^3(t, t') dt'$$

$$V(t)\tilde{c}_1(t) - V(t)c_1(t) = \int_{t_0}^t V(t')J(t')n[r(t, t')]dt'$$

$$\frac{\partial t}{\partial t} = \alpha_s \frac{1}{3} \left(\frac{1}{2\pi m_1} \right)^{1/2} c_1(t) = \frac{1}{3} a_0 \tau_{\text{coll}}$$

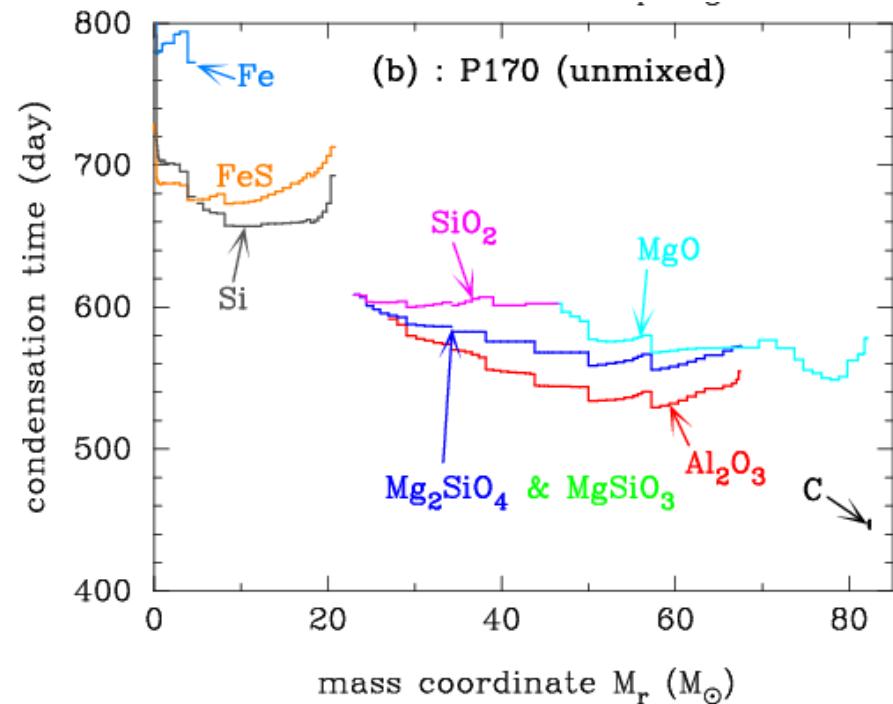
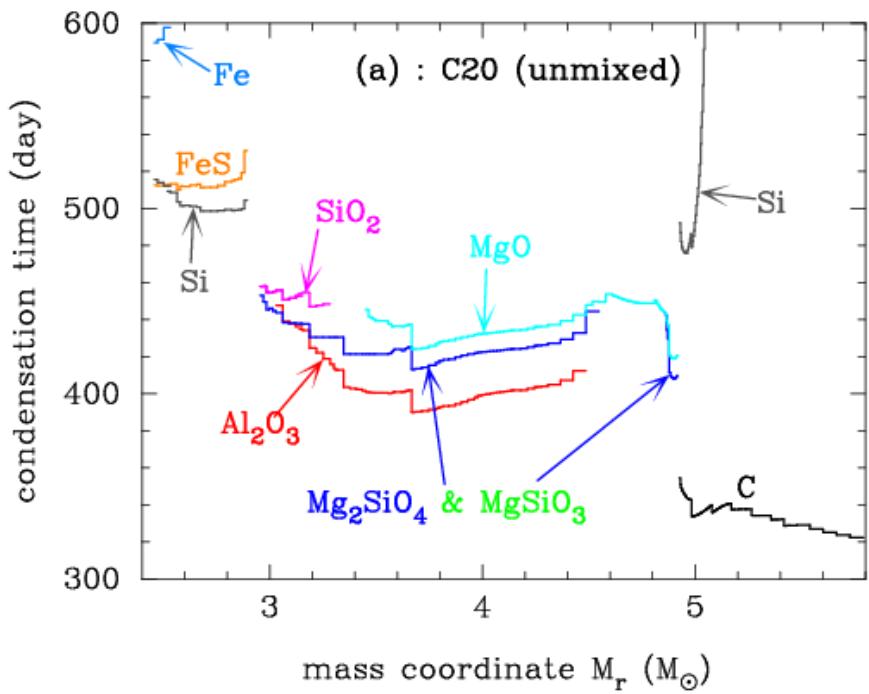
$$\frac{\partial V_d}{\partial t} = 4\pi r^2 \frac{\partial r}{\partial t} = \alpha_s \Omega 4\pi r^2 \langle v \rangle c_1(t)$$

$$\tau_{\text{coll}}^{-1}(t) = 4\pi a_0^2 \alpha_s \left(\frac{kT}{2\pi m_1} \right)^{1/2} c_1(t) \Big|_{t_0}$$

$$r(t, t_0) = r_* + \int_{t_0}^t \frac{1}{3} a_0 \tau_{\text{coll}}^{-1}(t') dt'$$

2-1-4. Dust formed in primordial SNe

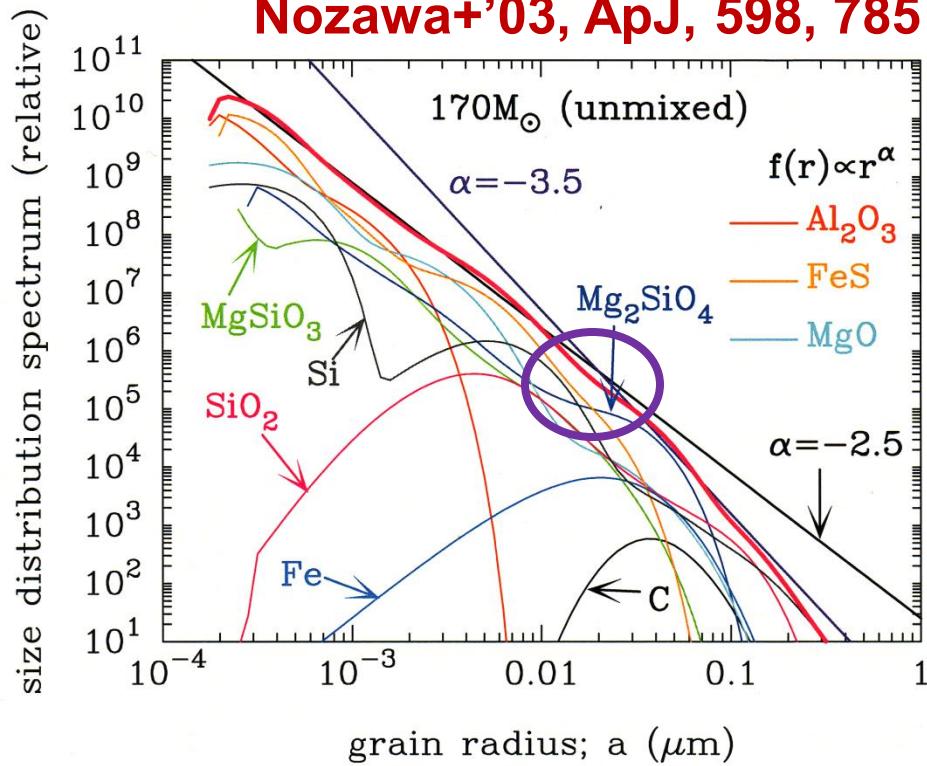
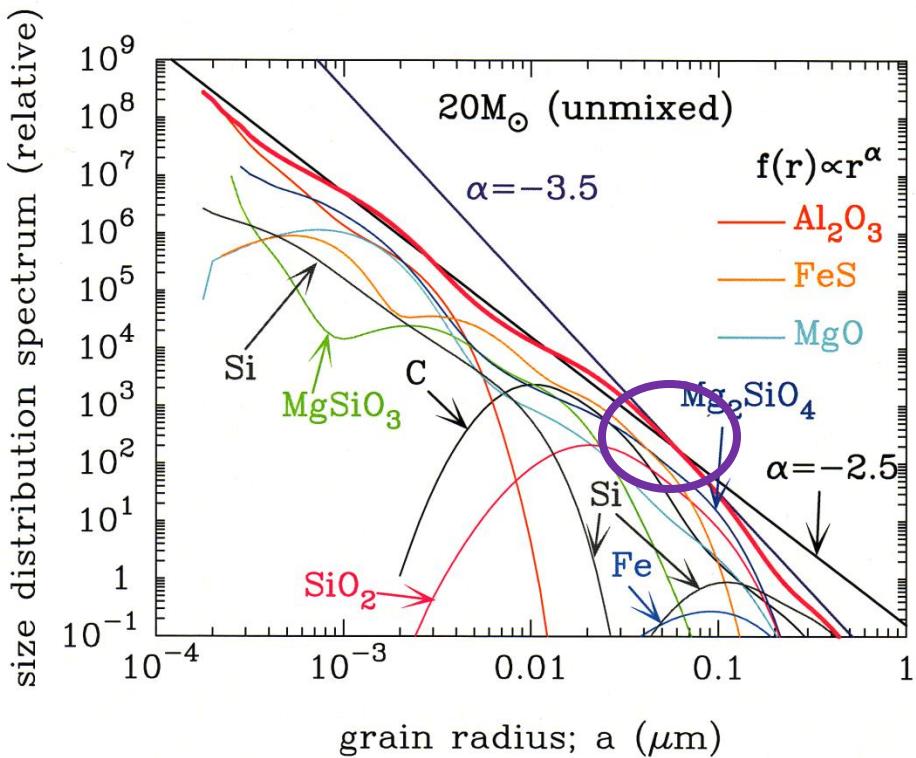
Nozawa+’03, ApJ, 598, 785



- Various dust species (C, MgSiO_3 , Mg_2SiO_4 , SiO_2 , Al_2O_3 , MgO , Si, FeS, Fe) form in the unmixed ejecta, according to the elemental composition of gas in each layer
- The condensation time: **300-600 days** for SNe II-P
400-800 days for PISNe

2-1-5. Size distribution of newly formed dust

Nozawa+’03, ApJ, 598, 785



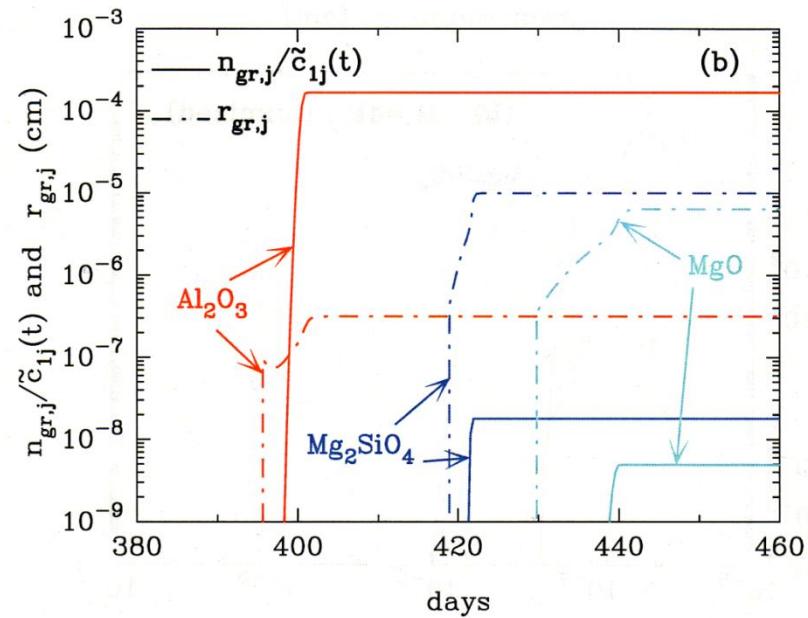
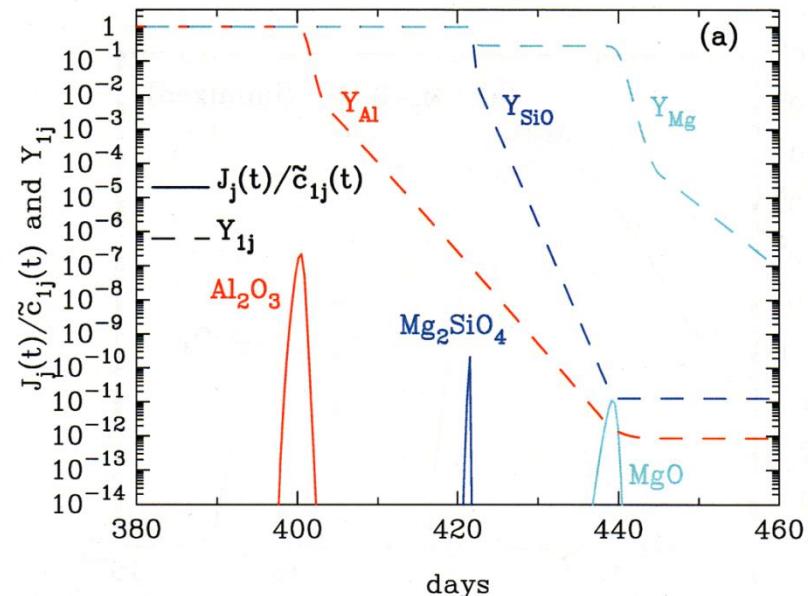
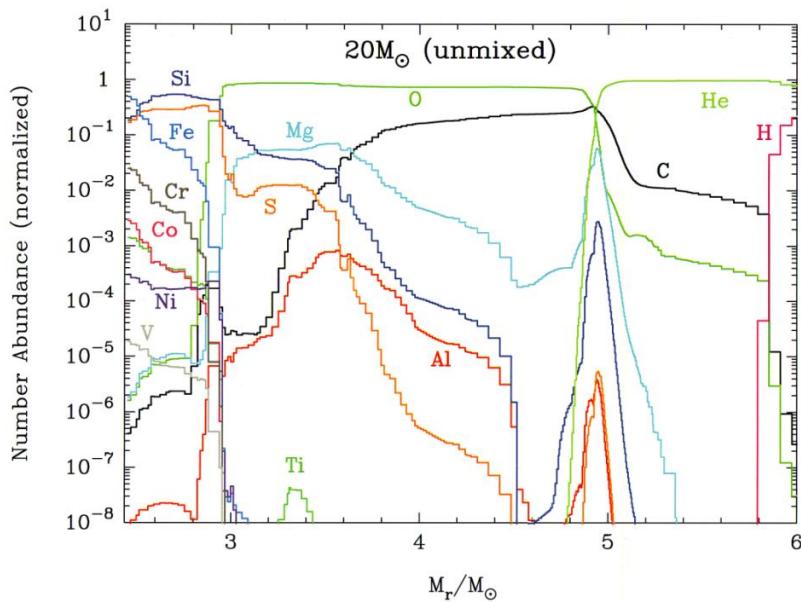
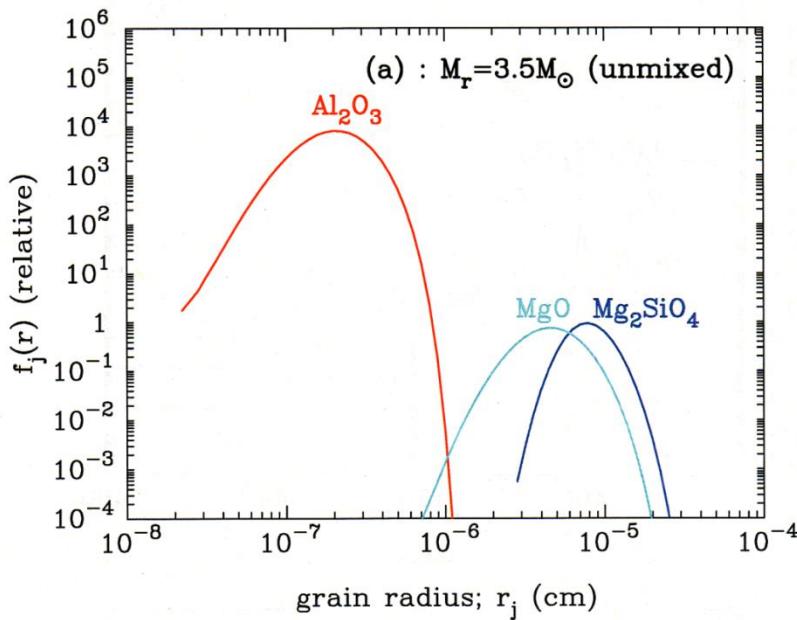
- grain radii range from a few Å up to 1 μm
- average dust radius is smaller for PISNe than SNe II-P

amount of newly formed dust grains

SNe II-P : $M_{\text{dust}} = 0.1\text{-}1 M_{\odot}$, $M_{\text{dust}} / M_{\text{metal}} = 0.2\text{-}0.3$

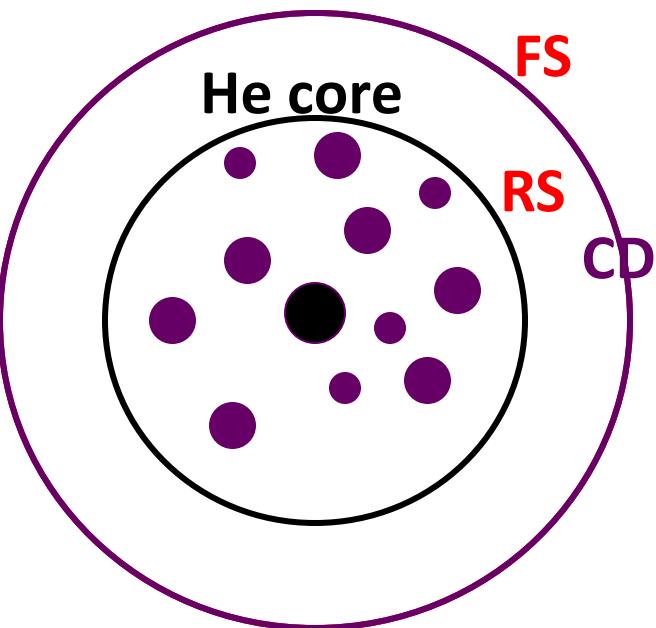
PISNe : $M_{\text{dust}} = 20\text{-}40 M_{\odot}$, $M_{\text{dust}} / M_{\text{metal}} = 0.3\text{-}0.4$

2-1-6. Behavior of dust formation



2-2. Dust Evolution in SNRs

$$T = (1-2) \times 10^4 \text{ K}$$
$$n_{H,0} = 0.1-1 \text{ cm}^{-3}$$



2-2-1. Time evolution of shock wave

- Basic equations (spherical symmetry)

$$\frac{\partial \rho}{\partial t} + \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \rho v) = 0$$

$$\frac{\partial}{\partial t} (\rho v) + \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \rho v^2) = - \frac{\partial P}{\partial r}$$

$$\frac{\partial}{\partial t} \left(\frac{\rho v^2}{2} + \frac{P}{\gamma - 1} \right) + \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \left[\frac{\rho v^2}{2} + \frac{\gamma P}{\gamma - 1} \right] v \right)$$

$$= -(n_e n_H \Lambda_{\text{gas}}(T) + \Lambda_{\text{ic}}(T) + \Lambda_{\text{d}}(n_H, T))$$

$\Lambda_{\text{gas}}(T)$: cooling function of gas by the atomic process

(Sutherland & Dopita 1993; Smith et al. 2001)

$\Lambda_{\text{ic}}(T)$: inverse Compton cooling (Ikeuchi & Ostriker 1986)

$$\Lambda_{\text{ic}}(T) = 5.41 \times 10^{-32} (1+z)^4 n_e (T/10^4 \text{ K}) \quad (\text{we adopt } z = 20)$$

$\Lambda_{\text{d}}(n_H, T)$: cooling of gas through thermal emission of dust

- numerical code : flux splitting method (van Albada et al. 1982)

2-2-2. Initial condition for shock waves

- Hydrodynamical model of SNe (Umeda & Nomoto'02)
 - SNe II : $M_{\text{pr}}=13, 20, 25, 30 \text{ Msun}$ ($E_{51}=1$)
 - PISNe : $M_{\text{pr}}=170$ ($E_{51}=20$), 200 Msun ($E_{51}=28$)
- The ambient medium (homogeneous)
 - gas temperature : $T = 10^4 \text{ K}$
 - gas density : $n_{\text{H},0} = 0.1, 1, \text{ and } 10 \text{ cm}^{-3}$
- Dust Model
 - initial size distribution and spatial distribution of dust
 - results of dust formation calculations
 - treating as a test particle

The calculation is performed from 10 yr up to $\sim 10^6$ yr

2-2-3. Dynamics of dust

- deceleration of dust due to drag force (Baines et al. 1965)

$$\frac{dw_d}{dt} = \frac{F_{\text{drag}}}{m_d} = -\frac{3n_H kT}{2a\rho_d} \sum_i A_i G_i(s_i) \quad (w_d : \text{relative velocity})$$

$$F_{\text{drag}} = m_d \frac{dw_d}{dt} = -\pi a^2 \sum n_i \langle v_i m_i v_i \cos \theta \rangle$$

$$\begin{aligned} \frac{dw_d}{dt} &= -\frac{\pi a^2}{\frac{4}{3}\pi a^3 \rho_d} n_H \sum A_i \langle v_i m_i v_i \cos \theta \rangle \\ &= -\frac{3n_H}{4a\rho_d} kT \sum A_i G_i \end{aligned}$$

$$G_i(s_i) \approx \frac{8s_i}{3\sqrt{\pi}} \left(1 + \frac{9\pi}{64} s_i^2 \right)^{\frac{1}{2}} \quad (\text{Draine \& Salpeter 1979})$$

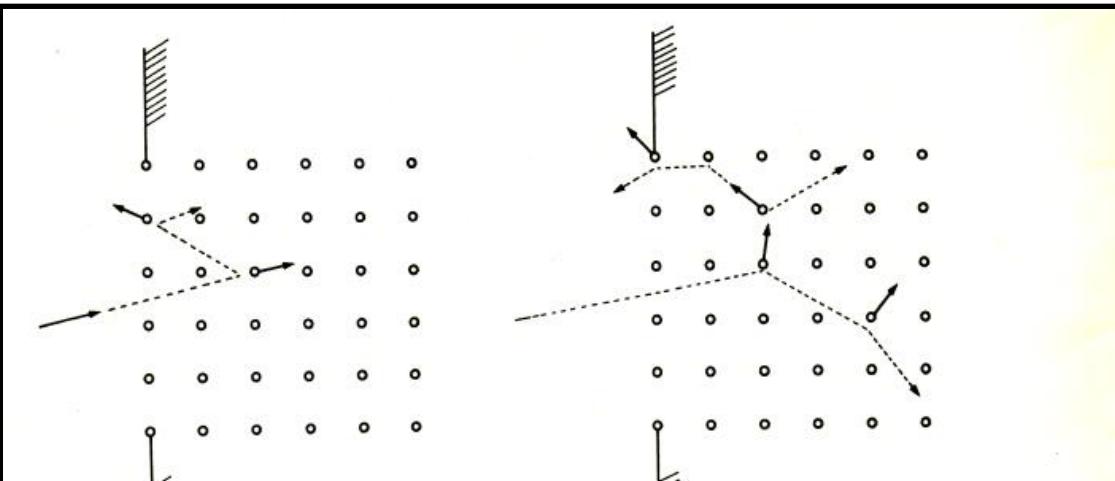
$$\text{where } s_i^2 = m_i w_d^2 / 2kT$$

2-2-4. Erosion rate of dust by sputtering

- dust destruction

996)

$$Y_i(E) = 2$$



$$\frac{dV_d}{dt} = 4\pi a^2 \frac{da}{dt} = -\pi a^2 \frac{4\pi a_0^3}{3} \sum n_i \langle v_i Y_i \rangle$$

• rate
acco

$$\frac{da}{dt} = -\frac{1}{4} \Omega n_{\text{H}} \sum A_i \langle v_i Y_i \rangle$$

is taken
)

$\mathcal{R}($

$$\frac{da}{dt} = -\frac{m_{\text{sp}}}{2\rho_d} n_{\text{H}} \sum A_i \left(\frac{8kT}{\pi m_i} \right)^{1/2} \frac{e^{-s_i^2}}{2s_i}) X_i(\epsilon) d\epsilon$$

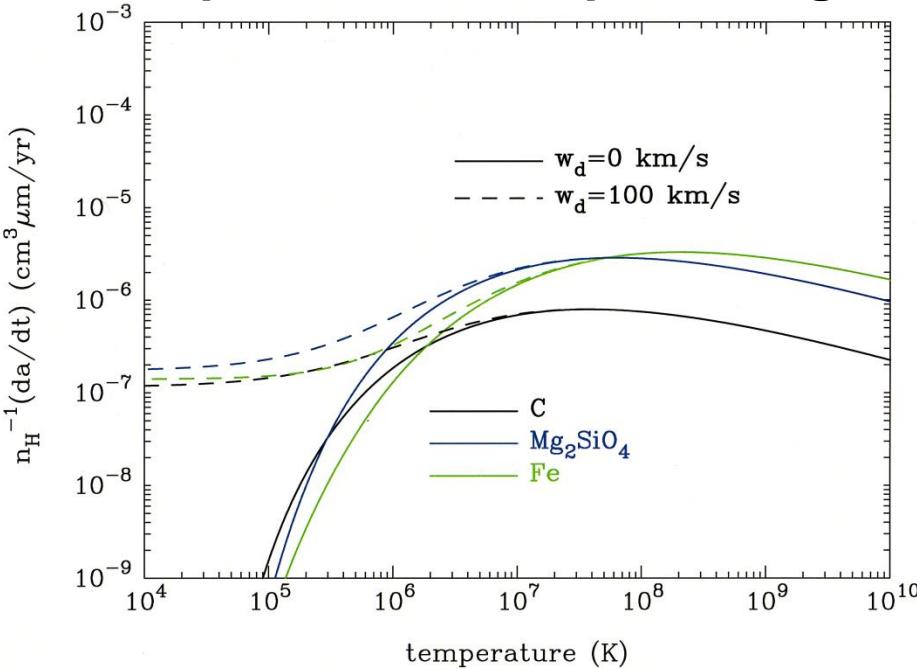
v

$$\times \int \epsilon_i^{1/2} e^{-\epsilon_i} \sinh(2s_i \epsilon_i^{1/2}) Y_i^0(\epsilon_i) d\epsilon_i$$

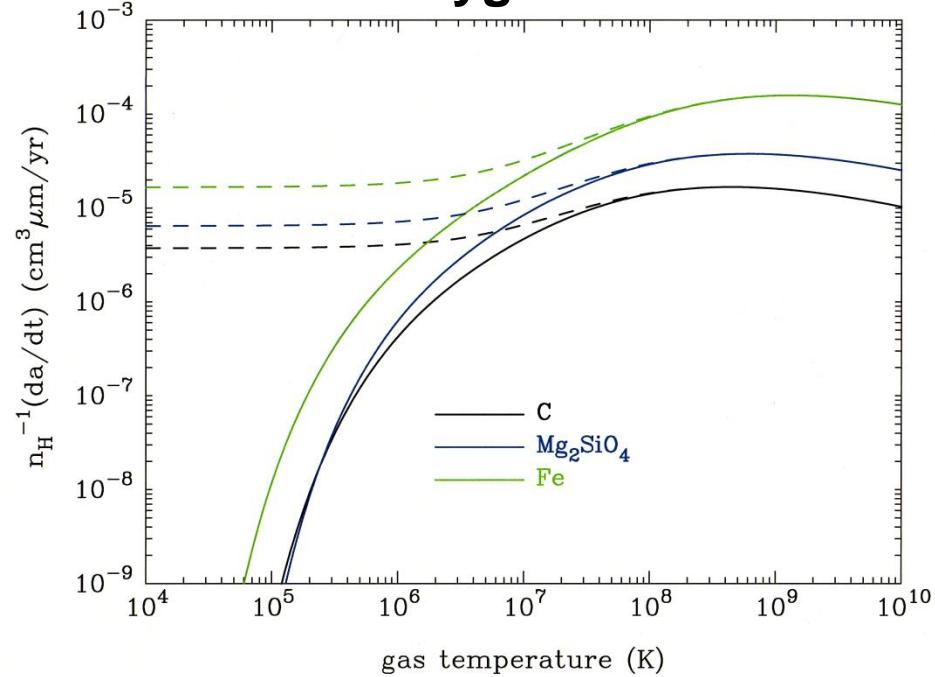
2-2-5. Erosion rate of dust by sputtering

Nozawa+’06, ApJ, 648, 435

for primordial composition gas



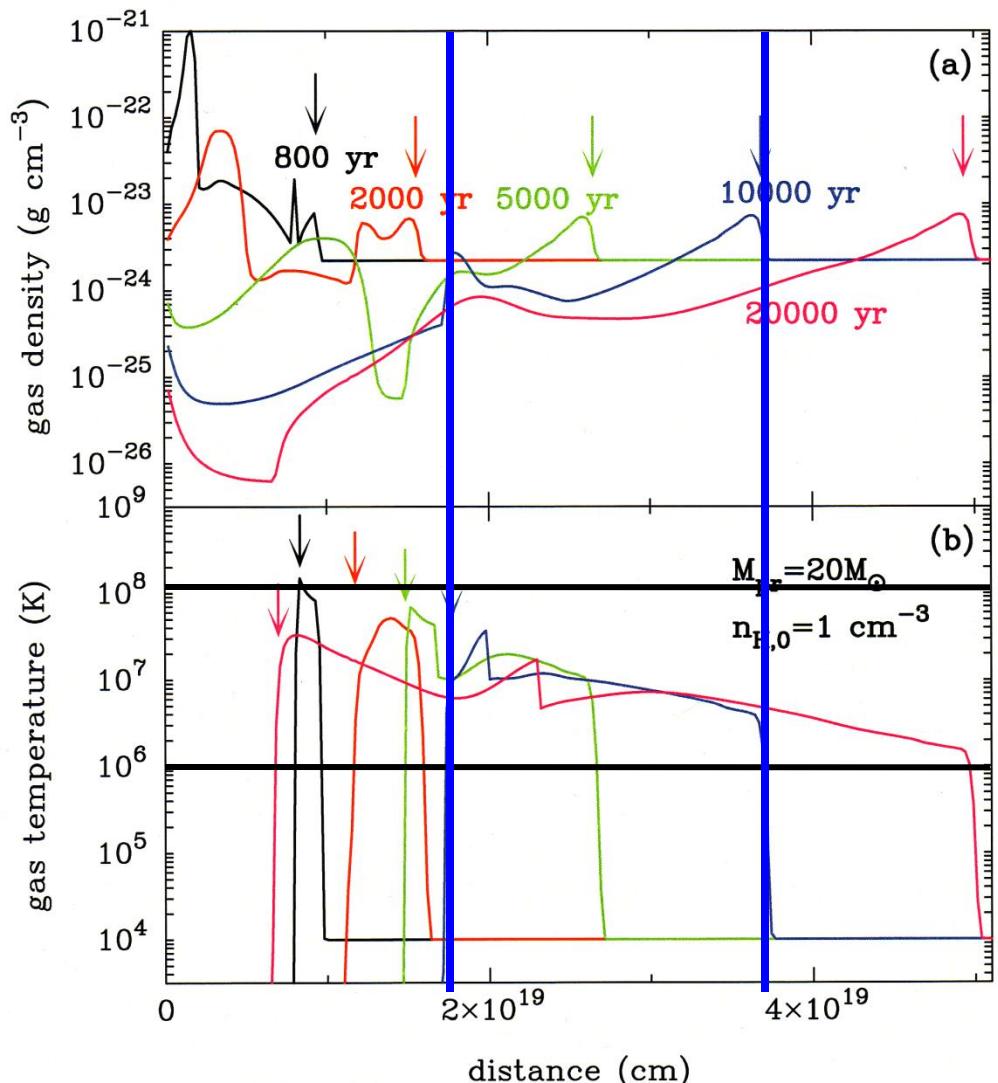
for oxygen ions



- erosion rate by sputtering quickly increases above 10^5 K and peaks at 10^7 - 10^8 K
- erosion rate : $da / dt \sim 10^{-6} n_H \mu\text{m yr}^{-1} \text{ cm}^3$
for the primordial gas (H and He) at $T > 10^6$ K

2-2-6. Temperature and density of gas in SNRs

Nozawa+’07, ApJ, 666, 955



Model : $M_{\text{pr}} = 20 \text{ Msun}$ ($E_{51}=1$)
 $n_{\text{H},0} = 1 \text{ cm}^{-3}$

Downward-pointing arrows:
forward shock in upper panel
reverse shock in lower panel

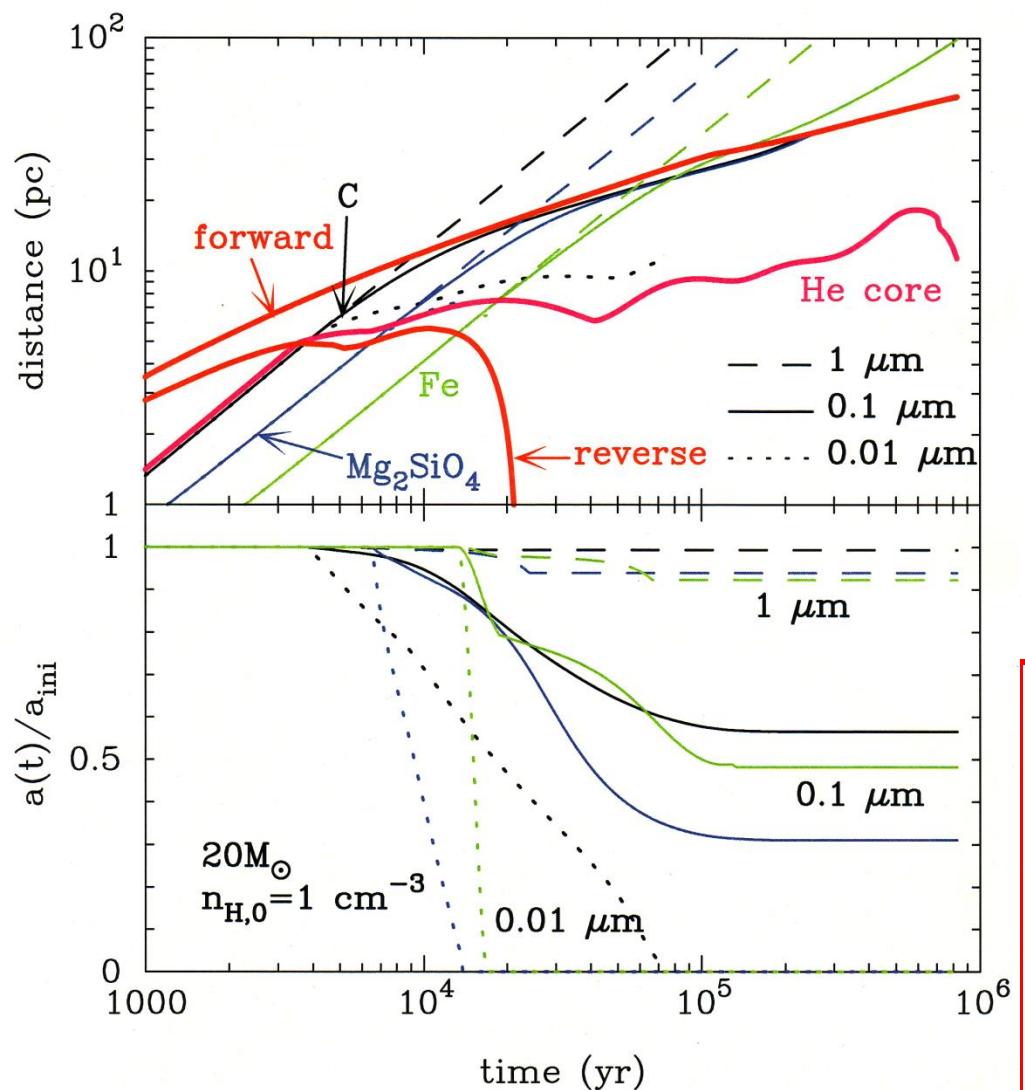
The temperature of the gas swept up by the shocks

$\rightarrow 10^6\text{-}10^8 \text{ K}$

\downarrow
Dust grains residing in the shocked hot gas are eroded by sputtering

2-2-7. Evolution of dust in SNRs

Nozawa+’07, ApJ, 666, 955



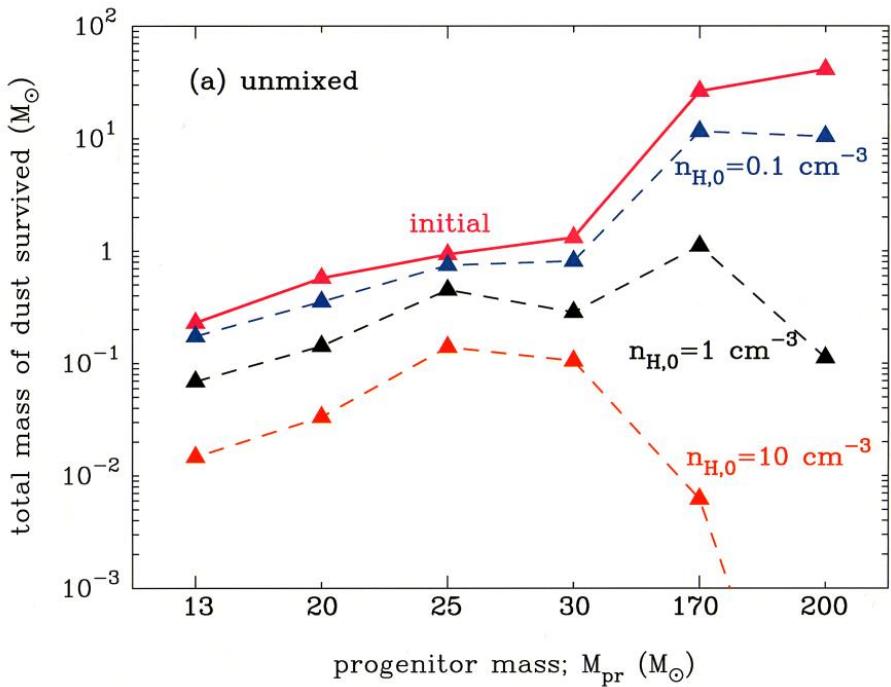
$a_{\text{ini}} = 0.01 \mu\text{m}$ (dotted lines)
→ completely destroyed

$a_{\text{ini}} = 0.1 \mu\text{m}$ (solid lines)
→ trapped in the shell

$a_{\text{ini}} = 1 \mu\text{m}$ (dashed lines)
→ injected into the ISM

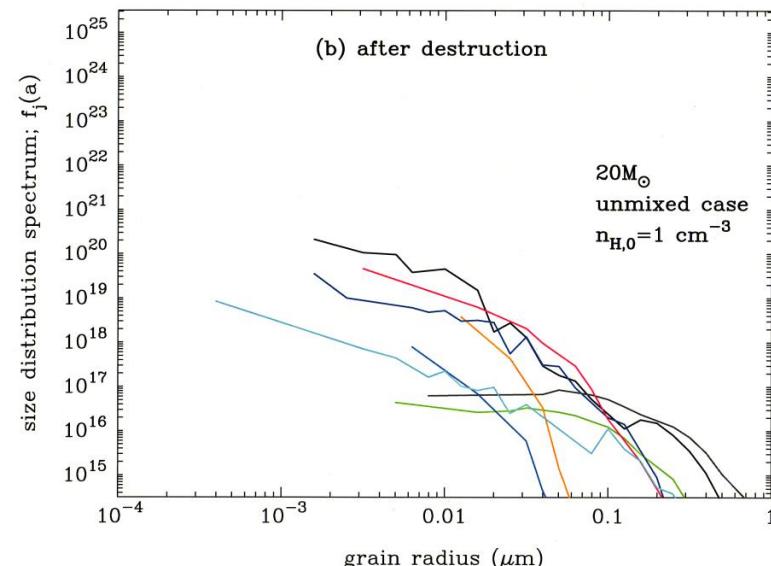
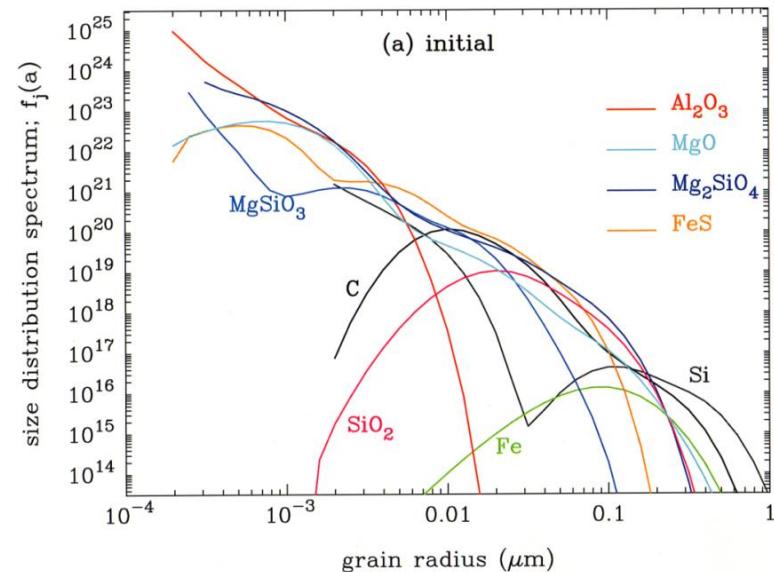
2-2-8. Total mass and size of surviving dust

Nozawa+’07, ApJ, 666, 955



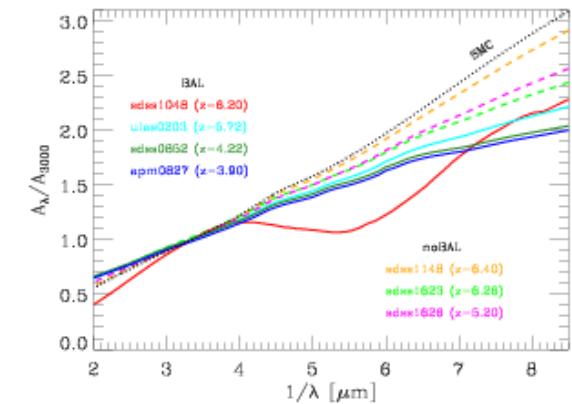
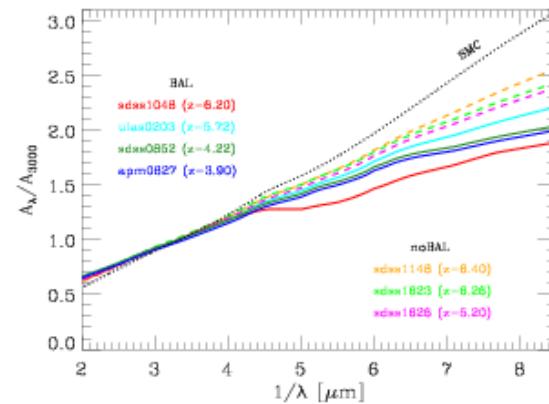
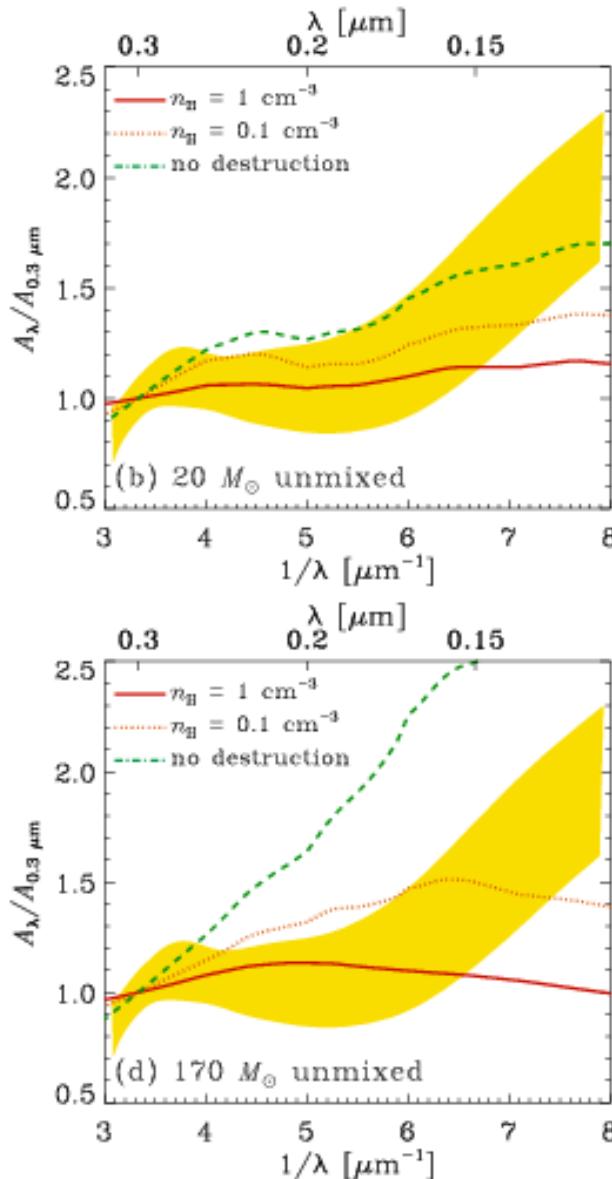
total dust mass surviving the destruction in Type II-P SNRs; $0.08\text{-}0.8 \text{ M}_{\odot}$ ($n_{\text{H},0} = 0.1\text{-}1 \text{ cm}^{-3}$)

size distribution of surviving dust is dominated by large grains ($> 0.01 \mu\text{m}$)

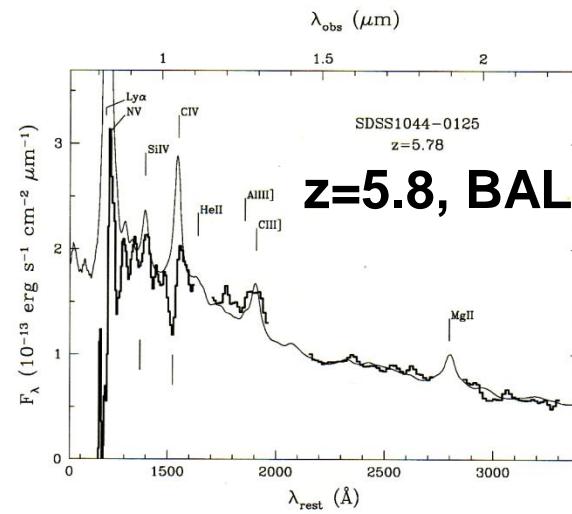


2-2-9. Flattened extinction curves at high-z

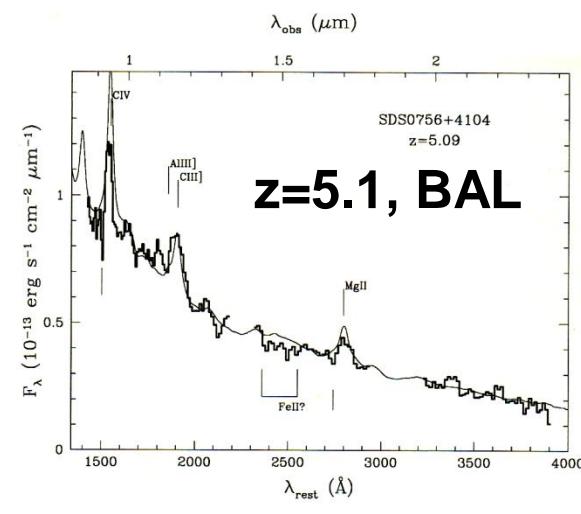
Hirashita, TN,+’08, MNRAS, 384, 1725



Gallerani+’10, A&A, 523, 85



$z=5.8$, BAL



$z=5.1$, BAL

Maiolino+’04, A&A, 420, 889

2-3. Summary of dust production in Pop III SNe

- The fate of newly formed dust within primordial SNRs strongly depends on the initial radii and compositions.
- The size distribution of dust surviving the destruction in SNRs is weighted to relatively large size ($> 0.01 \mu\text{m}$).
- The total mass of surviving dust grains decreases with increasing the ambient gas density
 - for $n_{\text{H},0} = 0.1\text{-}1 \text{ cm}^{-3}$
 - SNe II-P $\rightarrow M_{\text{dust}} = 0.1\text{-}0.8 \text{ Msun}$
 - PISNe $\rightarrow M_{\text{dust}} = 0.1\text{-}15 \text{ Msun}$
- Extinction curves in the early universe are expected to be flat if SNe II-P are main sources of dust at high z .

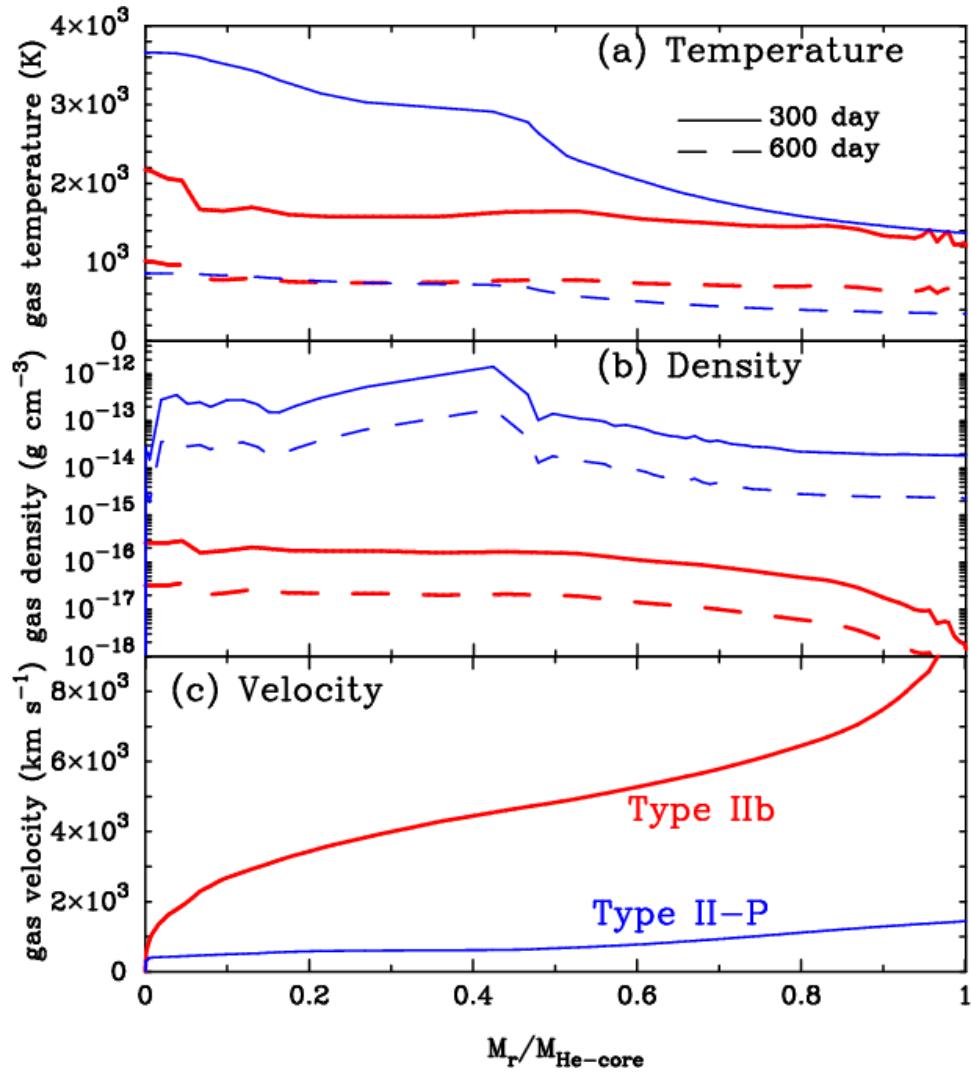
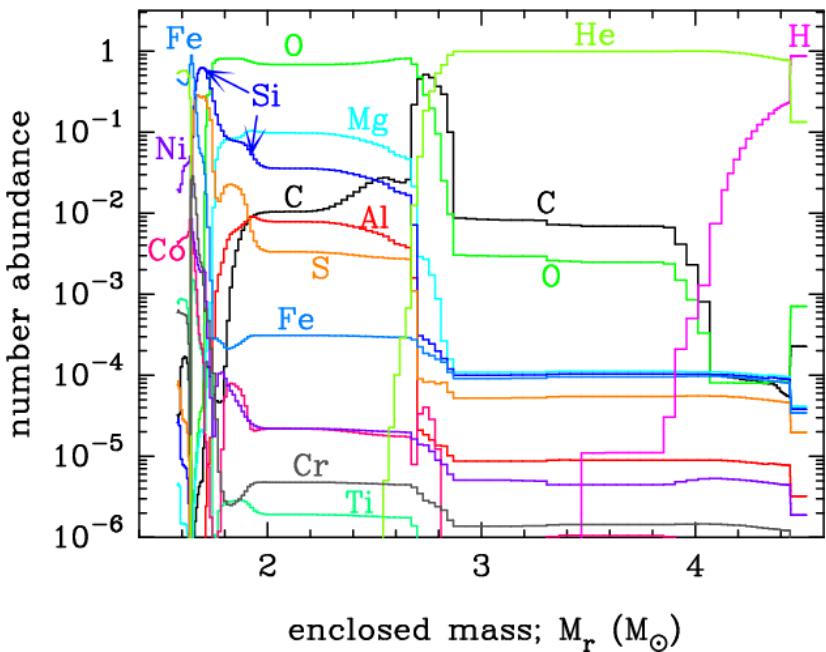
3. Formation of dust grains in various types of SNe

3-1-1. Dust formation in Type IIb SN

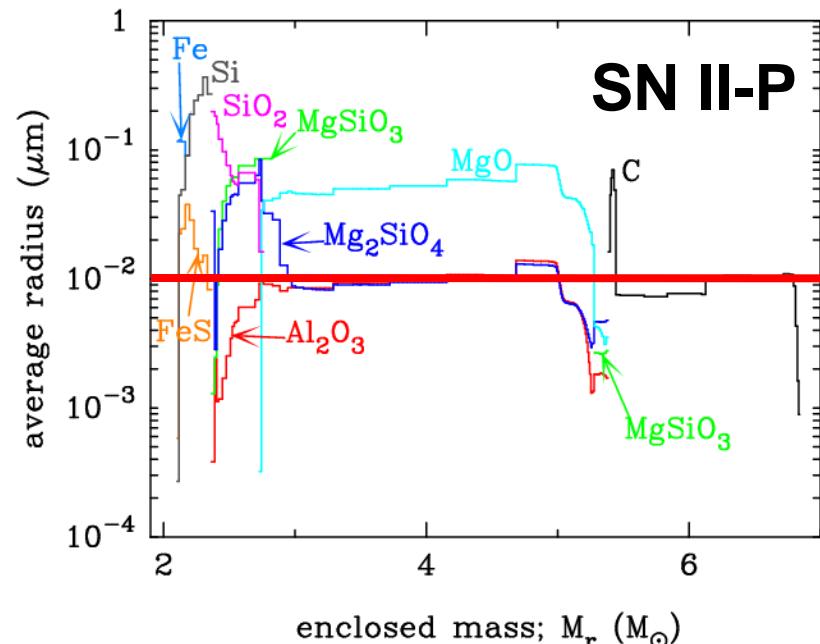
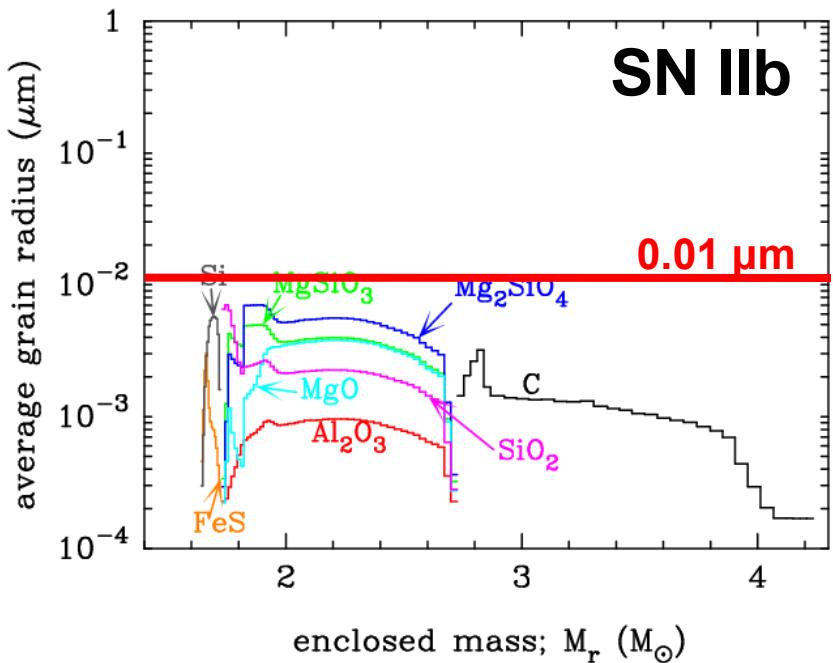


O SN IIb model (SN1993J-like model)

- $M_{\text{ej}} = 2.94 \text{ M}_{\odot}$
- $M_{\text{ZAMS}} = 18 \text{ M}_{\odot}$
- $M_{\text{H-env}} = 0.08 \text{ M}_{\odot}$
- $E_{51} = 1$
- $M(^{56}\text{Ni}) = 0.07 \text{ M}_{\odot}$



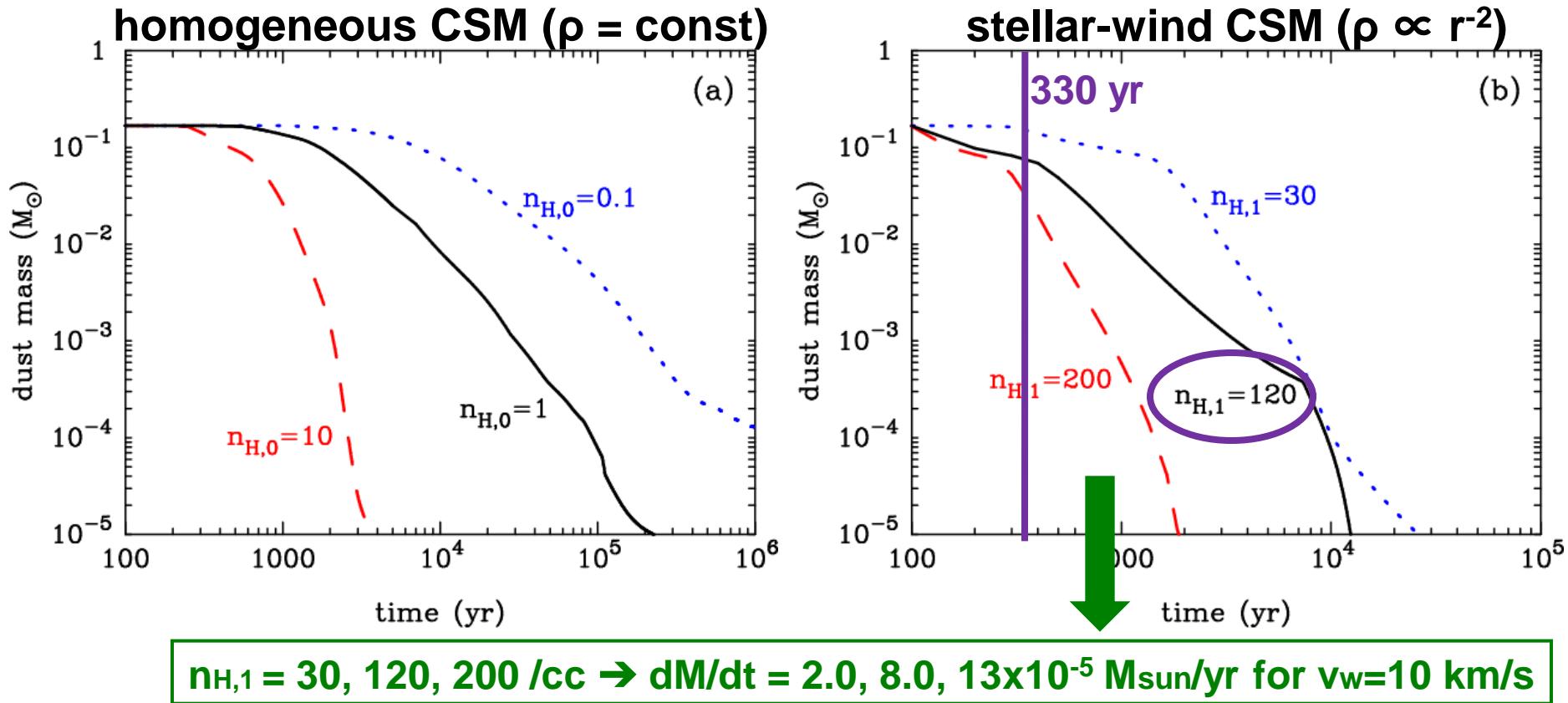
3-1-2. Dependence of dust radii on SN type



- condensation time of dust
300-700 d after explosion
- total mass of dust formed
 - **0.167 M_{\odot} in SN IIb**
 - **$0.1-1 \text{ M}_{\odot}$ in SN II-P**

- the radius of dust formed in H-stripped SNe is small
 - **SN IIb without massive H-env $\rightarrow a_{\text{dust}} < 0.01 \mu\text{m}$**
 - **SN II-P with massive H-env $\rightarrow a_{\text{dust}} > 0.01 \mu\text{m}$**

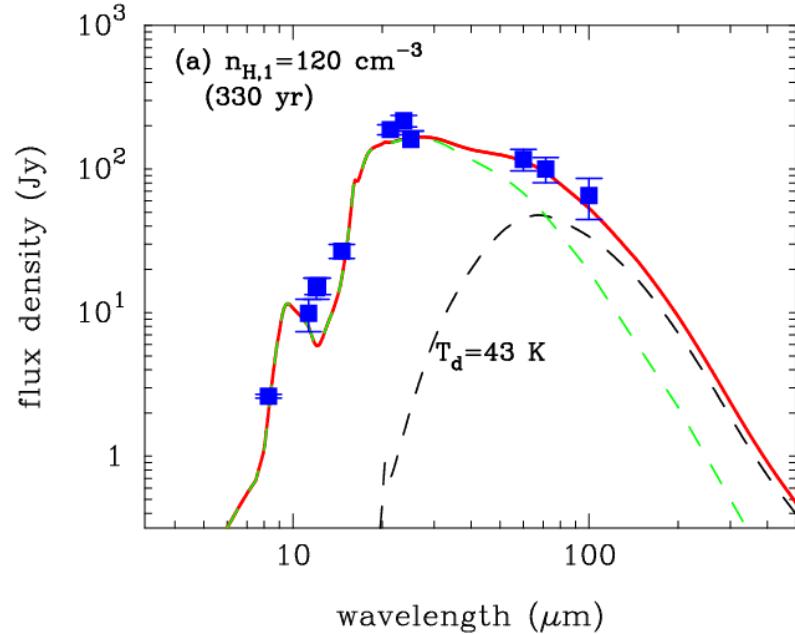
3-1-3. Destruction of dust in Type IIb SNR



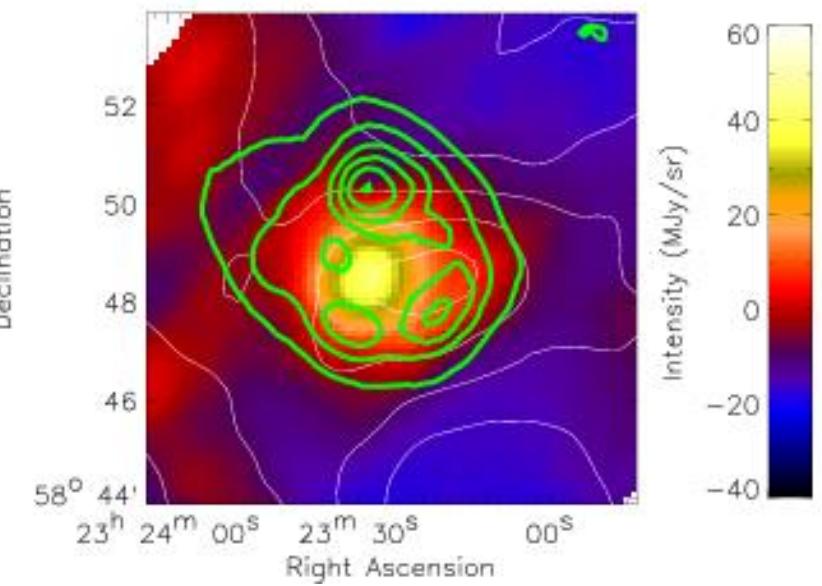
Almost all newly formed grains are destroyed in shocked gas within the SNR for CSM gas density of $n_H > 0.1 / \text{cc}$

- small radius of newly formed dust
- early arrival of reverse shock at dust-forming region

3-1-4. IR emission from dust in Cas A SNR



AKARI corrected 90 μm image

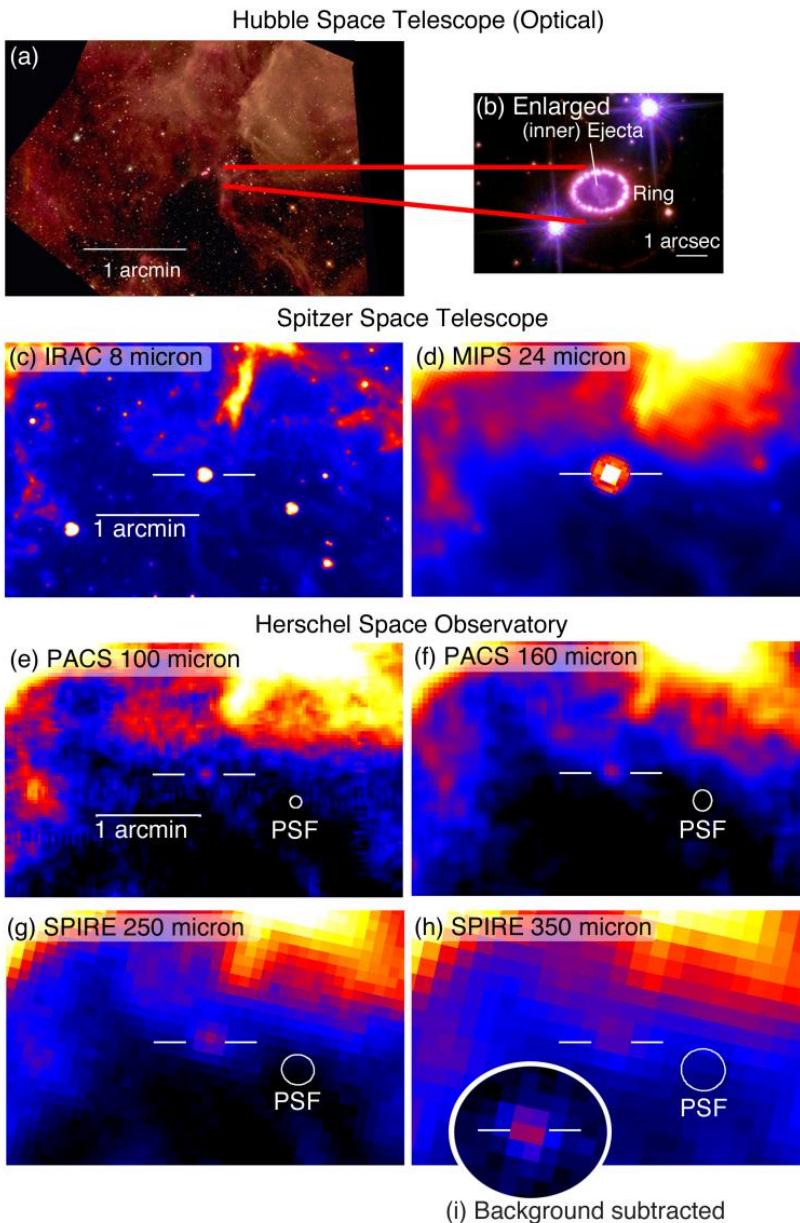


- total mass of dust formed
 $M_{\text{dust}} = 0.167 \text{ Msun}$
- shocked dust : **0.095 Msun**
 $M_{\text{d,warm}} = 0.008 \text{ Msun}$
- unshocked dust :
 $M_{\text{d,cool}} = 0.072 \text{ Msun}$
with $T_{\text{dust}} \sim 40 \text{ K}$

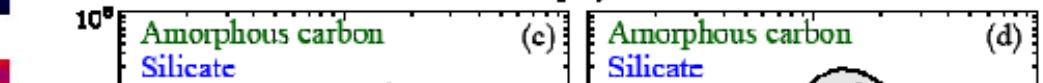
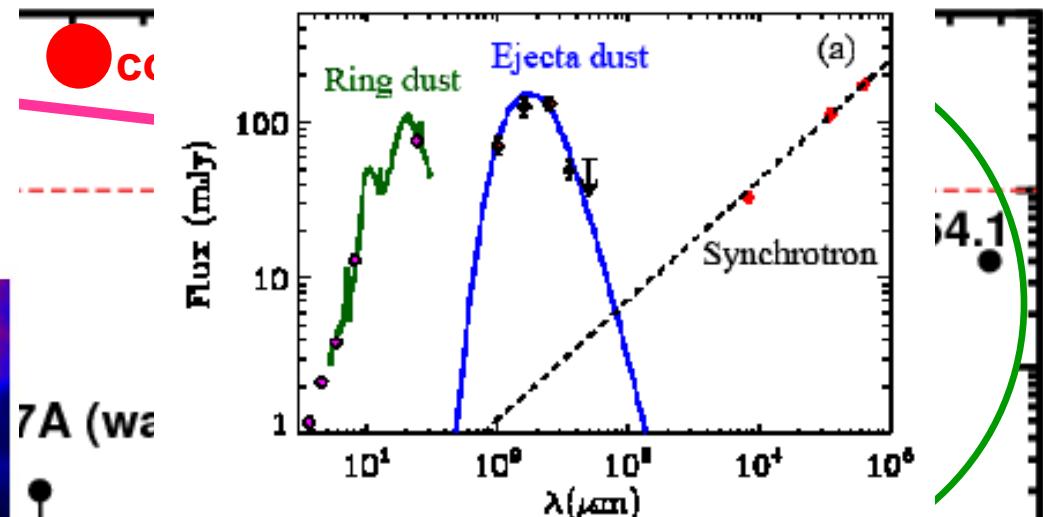
AKARI observation
 $M_{\text{d,cool}} = 0.03-0.06 \text{ Msun}$
 $T_{\text{dust}} = 33-41 \text{ K}$
(Sibthorpe+'10)

Herschel observation
 $M_{\text{d,cool}} = 0.075 \text{ Msun}$
 $T_{\text{dust}} \sim 35 \text{ K}$ (Barlow+'10)

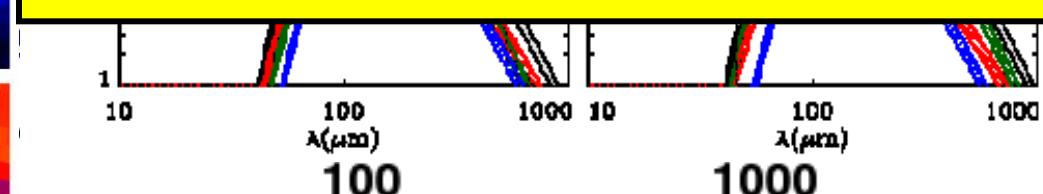
3-2-1. Missing-dust problem in CCSNe



Matsuura, ..., TN, et al. 2011, Science

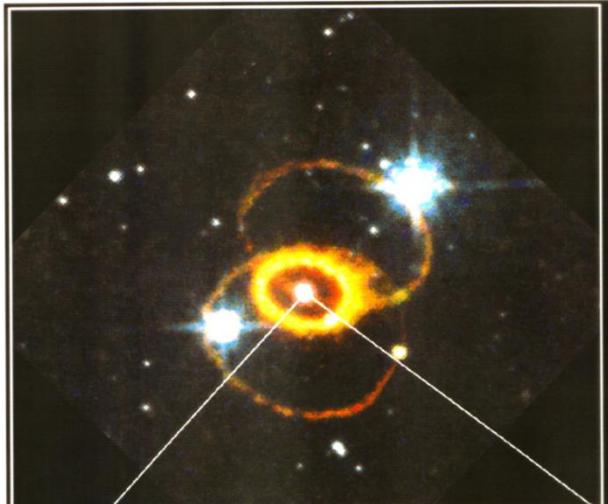
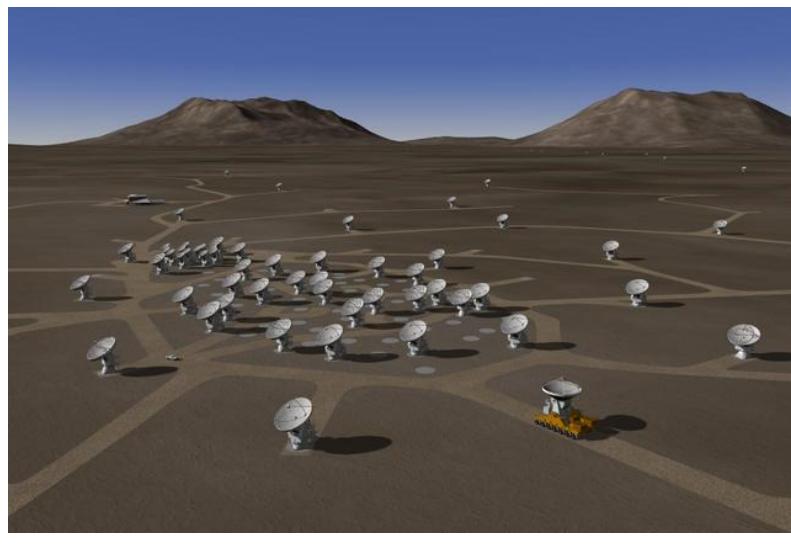


Herschel detects cool (20K) dust of 0.4-0.7 Msun toward SN 1987A!



Notice in the mass of dust between observations and theoretical predictions!!

3-2-2. Resolving cold dust in SN87A with ALMA

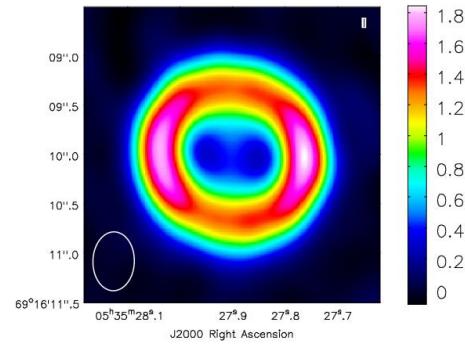


Supernova 1987A

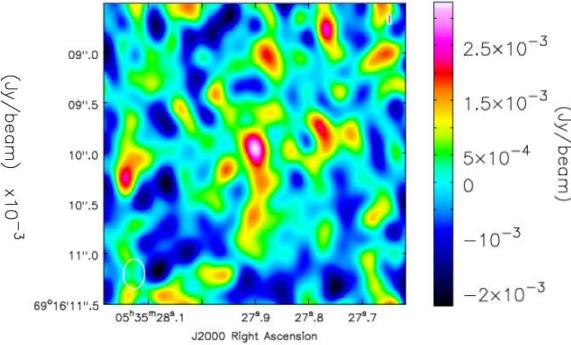
PRC97-03 • ST Scl OPO • January 14, 1997
J. Pun (NASA/GSFC), R. Kirshner (CfA) and NASA

ALMA Cycle 0 Proposal
'Detecting cool dust in SN1987A'
(PI: Nozawa)

Band 7 (850 μm)



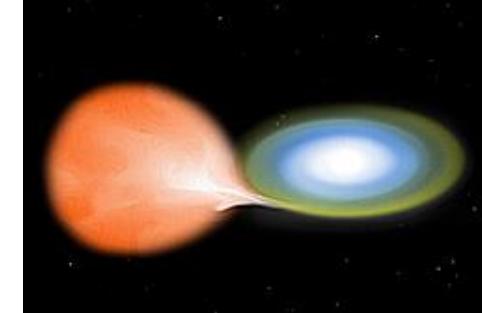
Band 9 (450 μm)



0.1 M_{\odot} of silicate $\rightarrow 5\sigma$ detection

This proposal is ranked in the highest priority to be observed !

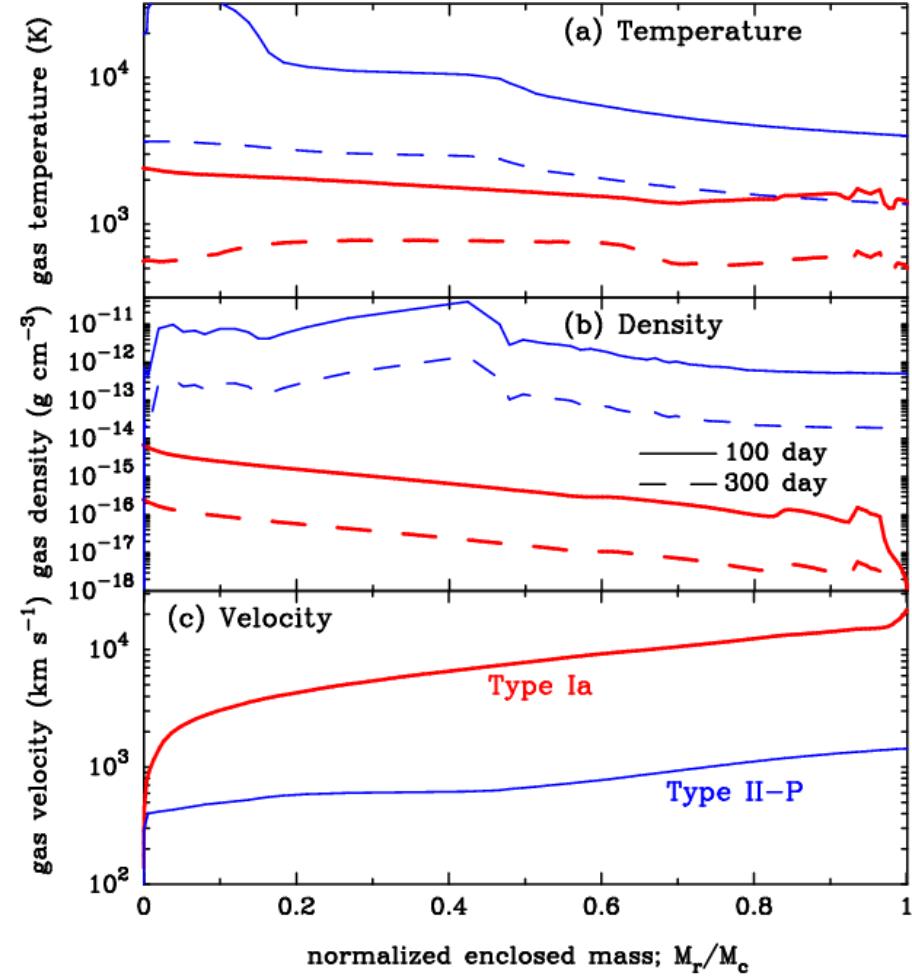
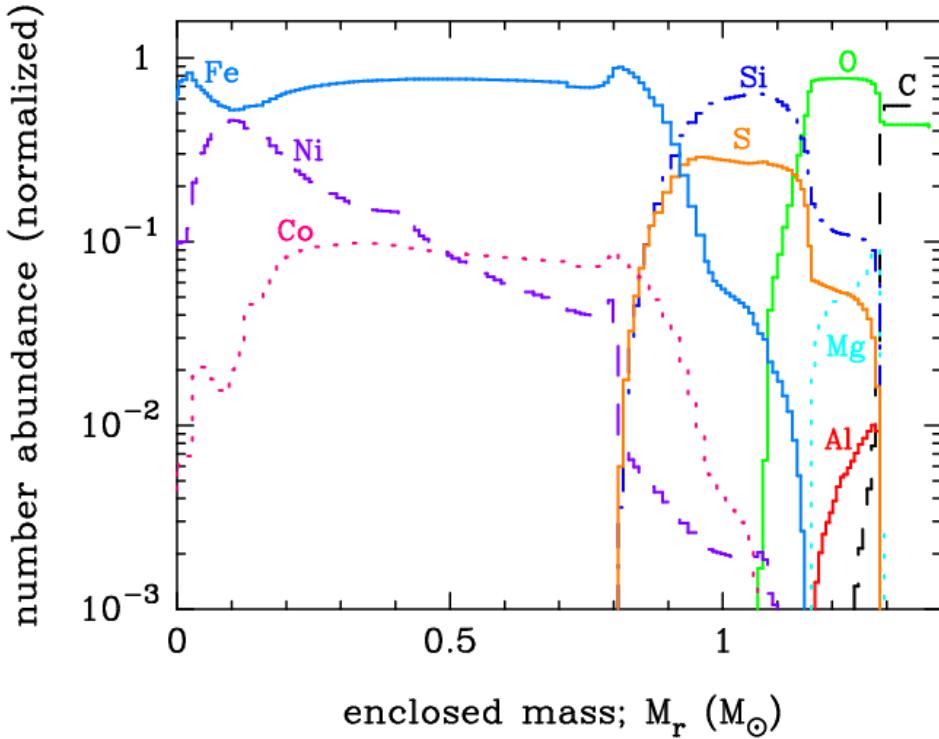
3-3-1. Dust formation in Type Ia SN



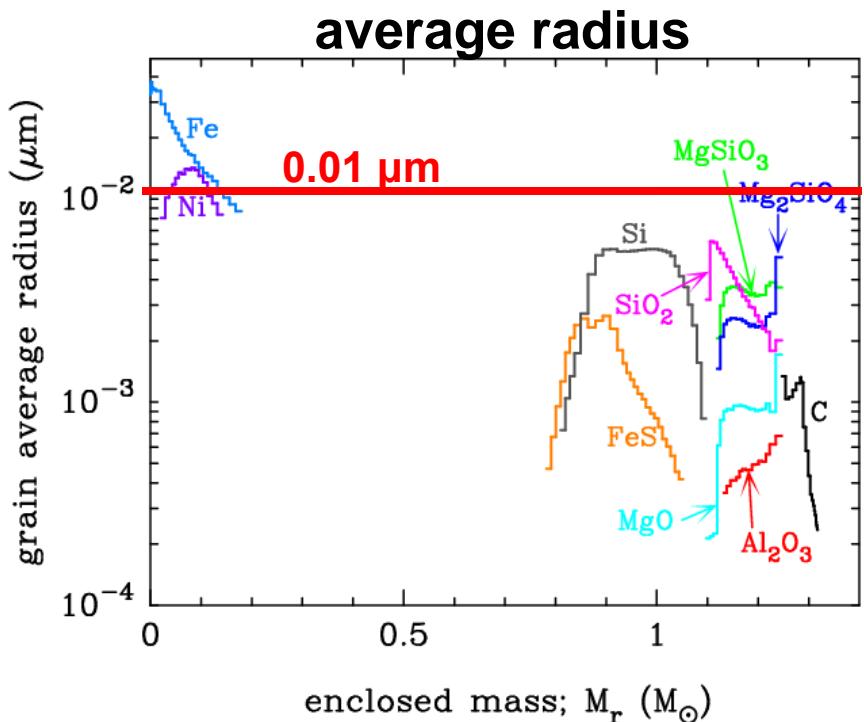
O Type Ia SN model

W7 model (C-deflagration) (Nomoto+’84; Thielemann+’86)

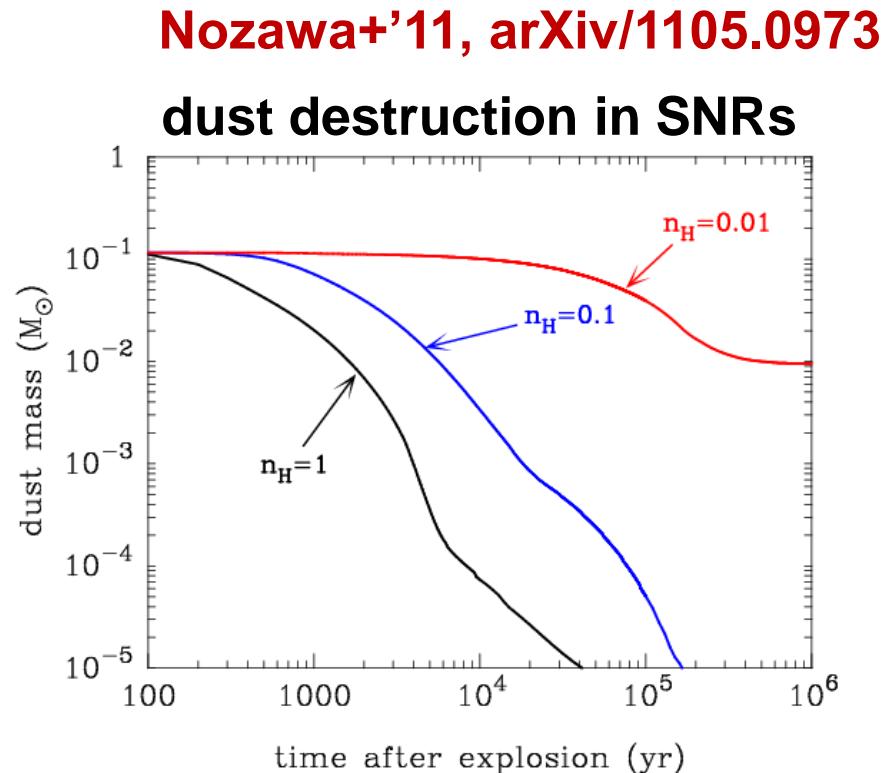
- $M_{\text{ej}} = 1.38 \text{ M}_{\odot}$
- $E_{51} = 1.3$
- $M(^{56}\text{Ni}) = 0.6 \text{ M}_{\odot}$



3-3-2. Dust formation and evolution in SNe Ia



- condensation time :
100-300 days
- average radius of dust :
 $a_{\text{ave}} < \sim 0.01 \mu\text{m}$
- total dust mass :
 $M_{\text{dust}} = 0.1\text{-}0.2 M_\odot$

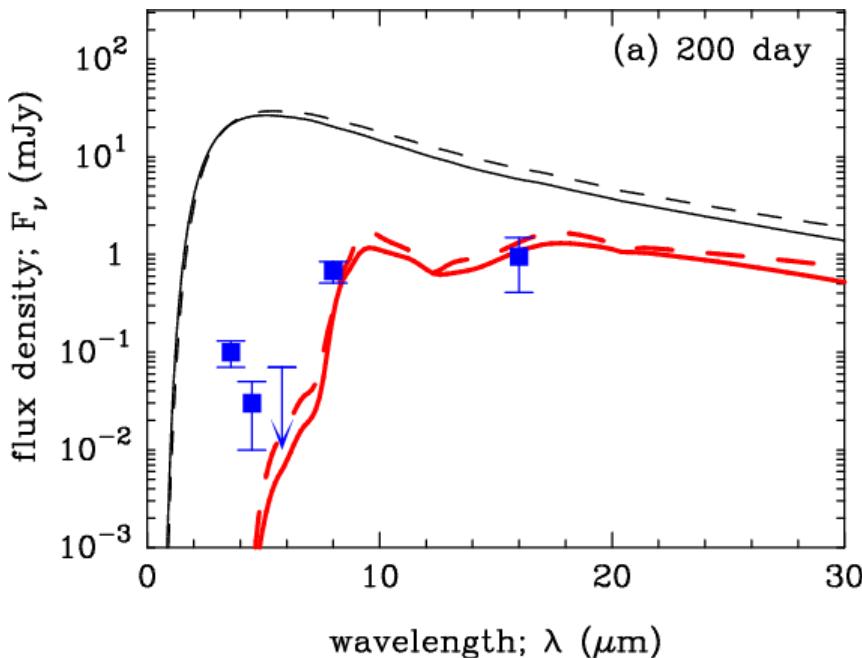


newly formed grains are completely destroyed for ISM density of $n_H > 0.1 \text{ cm}^{-3}$
→ SNe Ia are unlikely to be major sources of dust

3-3-3. Carbon dust and outermost layer

- There has been no evidence for dust formation in SNe Ia
→ Formation of massive carbon dust does not match the observations

Observational data : SN 2005df at day 200 and 400 (Gerardy+'07)



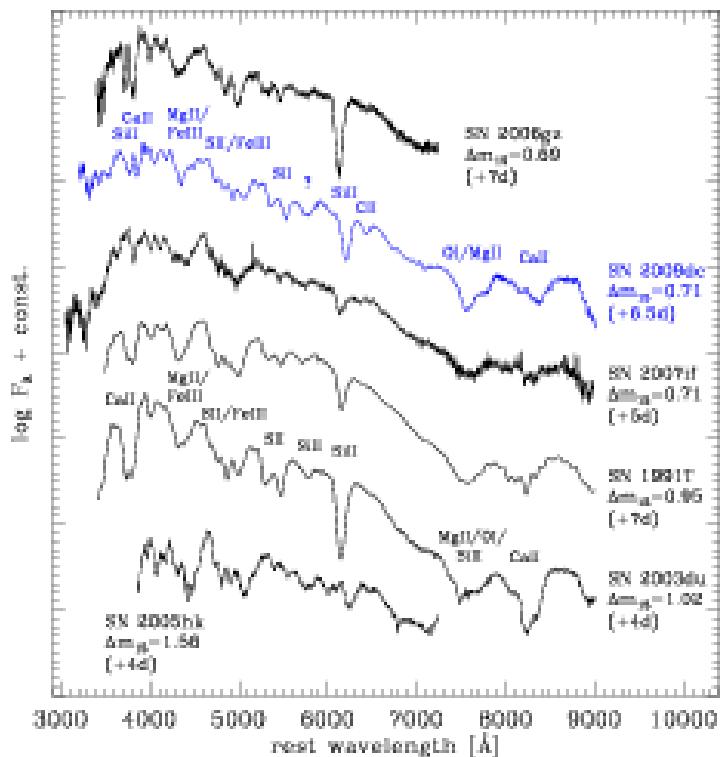
Nozawa+'11, accepted

- massive unburned carbon ($\sim 0.05 \text{ M}_\odot$) in deflagration
 - change of composition of WD by He-shell flash
 - burning of carbon by a delayed detonation

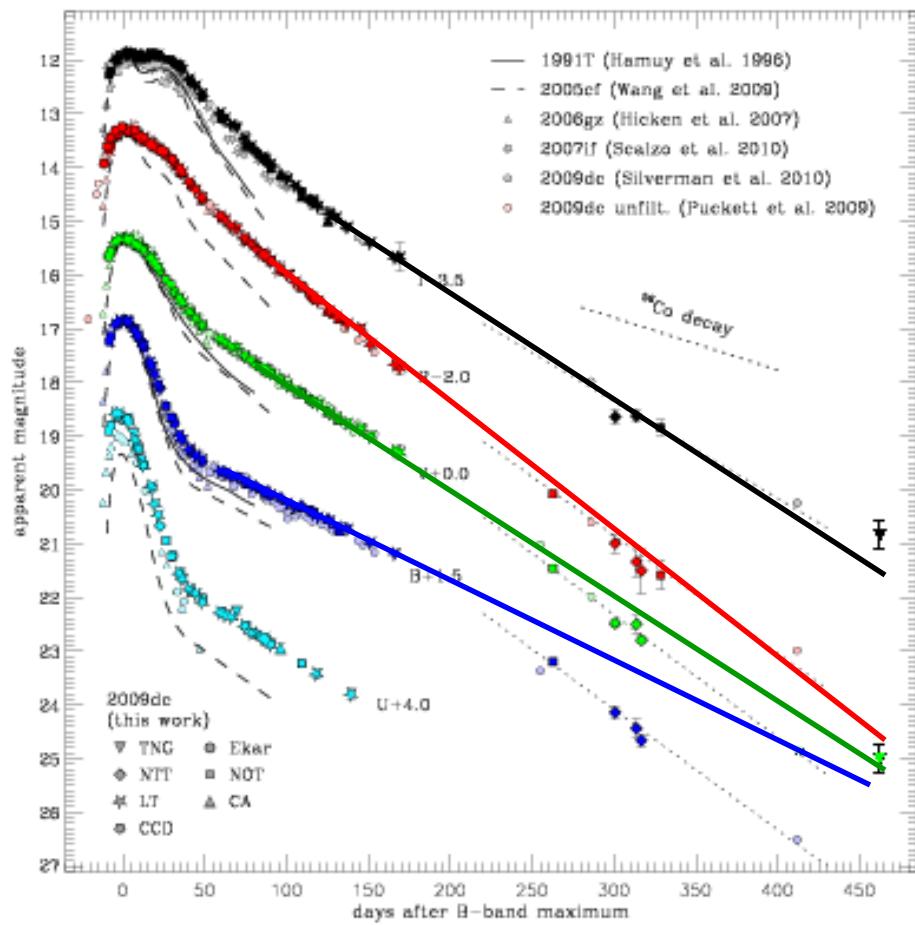
observationally estimated carbon mass in SNe Ia :
 $M_C < 0.01 \text{ M}_\odot$
(Marion+'06; Tanaka+'08)

3-3-4. Dust formation in super-Chandra SNe?

- super-Chandra SNe :
 $M(56\text{Ni}) \sim 1.0 \text{ M}_{\odot}$
- detection of CII line**
- presence of massive unburned carbon



SN 2009dc, Tarbenberger+’10



enhanced fading at ~200 day
→ formation of carbon dust?

4. Implication and Summary

4-1. Implication on evolution history of dust (1)

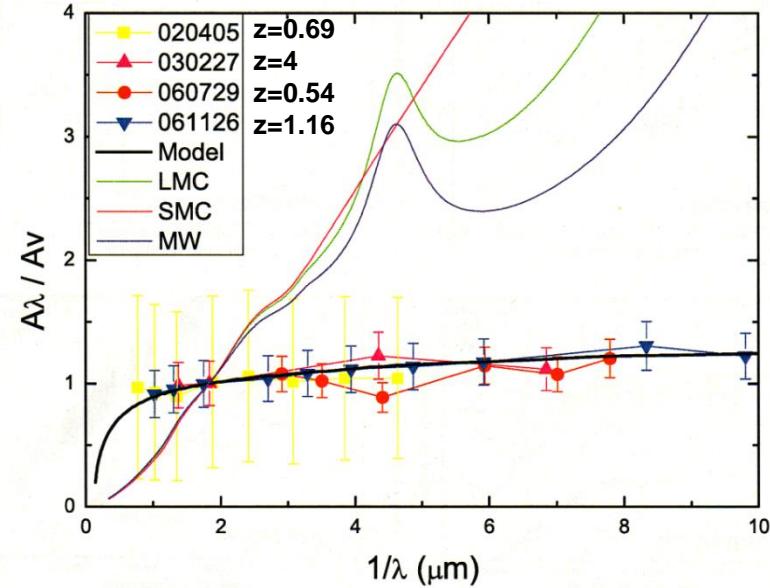
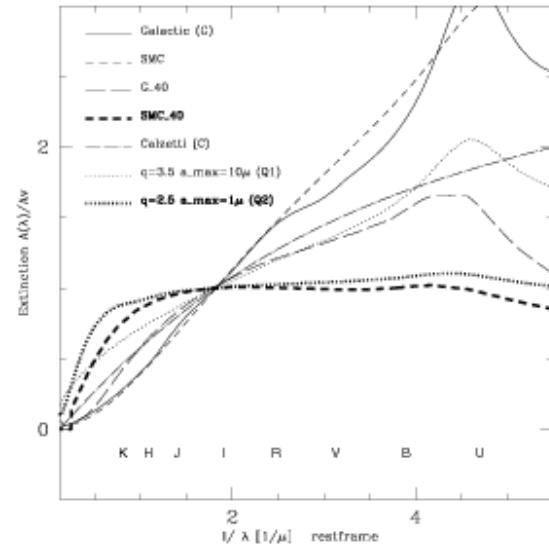
O metal-poor (high-z or starburst) galaxies

- massive stars (SNe) are dominate
- mass loss of massive stars would be less efficient

→ Type II-P SNe might be major sources of dust

- average radius of dust is relatively large ($> 0.01 \mu\text{m}$)
- grain growth makes grain size larger

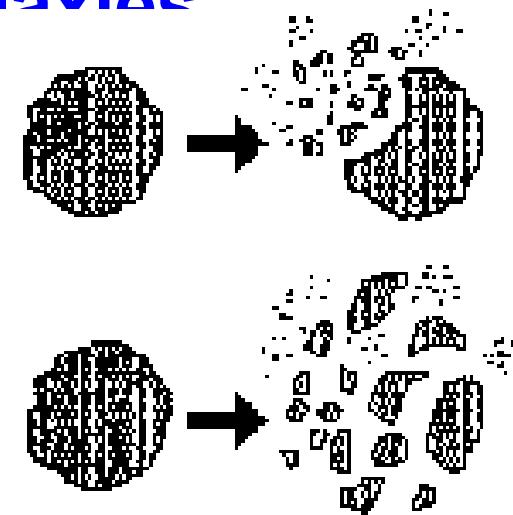
→ dust extinction curve might be gray



4-2. Implication on evolution history of dust (2)

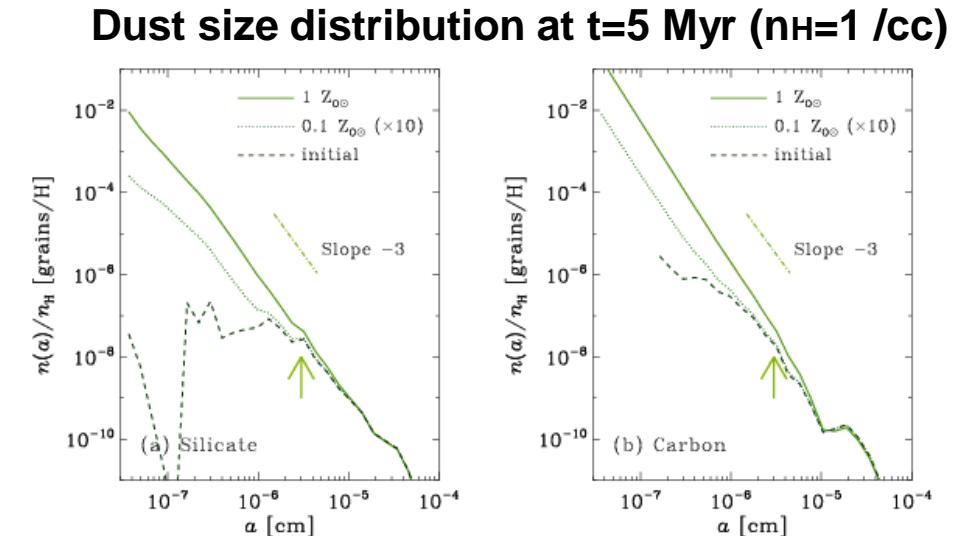
O metal-rich (low-z or Milky Way) galaxies

- low-mass stars are dominate
- mass loss of massive stars would be
- SNe (I Ib, I b/c, Ia) might be minor sour
- dust from AGB stars may also be larg
- How are small dust grains produced?



grain shattering

- warm ionized medium (WIM)
relative velocity of dust in
turbulence : 1-20 km/s
- grain shattering is efficient
in WIM at t=5 Myr if metallicity
is solar and more
- the production of small grains
by shattering steepens the
extinction curve



4-3. Summary of this talk

- SNe are important sources of dust?
 - maybe, Yes in the early universe
(at least, to serve the seeds for grain growth in the ISM)
- Size of newly formed dust depends on types of SNe
 - H-retaining SNe (Type II-P) : $a_{ave} > 0.01 \mu\text{m}$
 - H-stripped SNe (Type I Ib/Ic and Ia) : $a_{ave} < 0.01 \mu\text{m}$
 - dust is almost completely destroyed in the SNRs
 - H-stripped SNe may be poor producers of dust
- Our model treating dust formation and evolution self-consistently can reproduce IR emission from Cas A
- Mass of dust in SNe must be dominated by cool dust
 - FIR and submm observations of SNe are essential