

Lecture on Cosmic Dust

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Today's Contents:

- 1) Composition of dust
- 2) Extinction of stellar lights by dust
- 3) Thermal radiation from dust

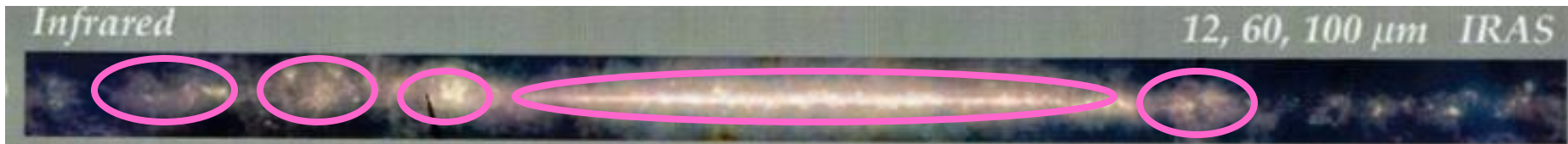
Introduction

- **Cosmic dust**: solid particle with size of a few Å to 1 mm
interplanetary dust, interstellar dust, intergalactic dust

Milky Way (optical)



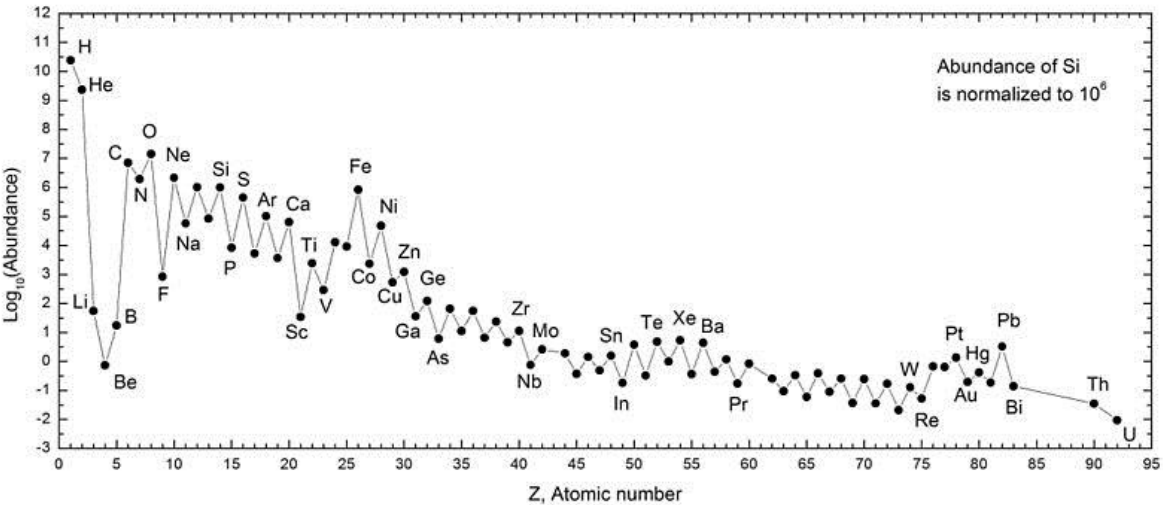
Milky Way (infrared)



Dust grains absorb UV/optical photons and reemit by thermal radiation at IR wavelengths!

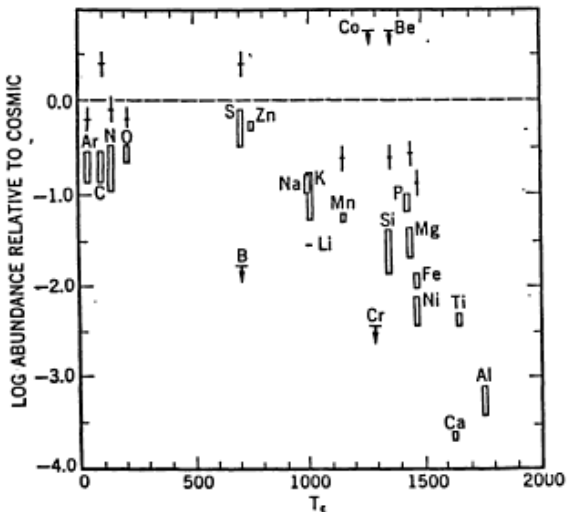
1-1. Hints as to composition of cosmic dust

Dust particles are composed of metals ($N > 5$)
→ what are abundances of metals in space?



Solar elemental abundances
(Asplund+2009, ARAA, 47, 481)

Element	Log10(n)	Ratio to H
H	12.00	1.00
He	10.93	8.51×10^{-2}
O	8.69	4.90×10^{-4}
C	8.43	2.69×10^{-4}
Ne	7.93	8.51×10^{-5}
N	7.83	6.76×10^{-5}
Mg	7.60	3.98×10^{-5}
Si	7.51	3.24×10^{-5}
Fe	7.50	3.16×10^{-5}
S	7.12	1.32×10^{-5}



depletion of elements
(Sakurai 1993)

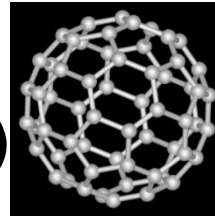
More than 90 % of Si, Mg, Fe, Al, and Ca are depleted in the ISM

1-2. Expected composition of dust (1)

○ Major candidates of cosmic dust

▪ carbonaceous grains (C-based)

- graphite
- amorphous carbon
- diamond, C₆₀ (fullerene)



▪ silicate grains (SiO₄²⁻-based)

- Mg_xFe_(1-x)SiO₃ (pyroxene):

MgSiO₃ (enstatite)

FeSiO₃ (ferrosilite)

- Mg_{2x}Fe_{2(1-x)}SiO₄ (olivine):

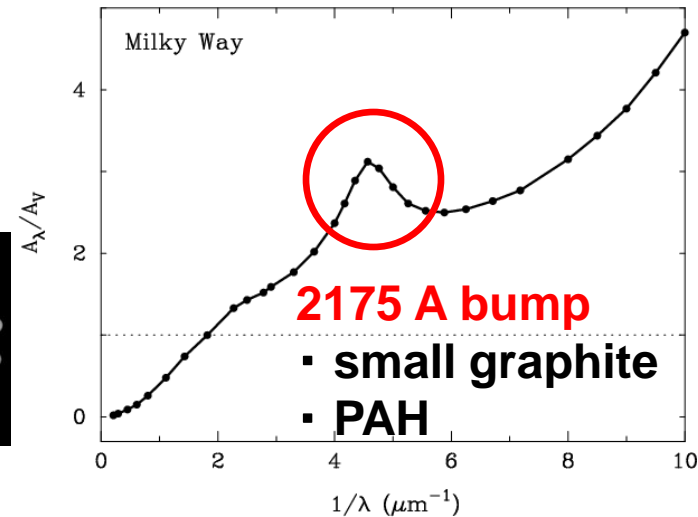
Mg₂SiO₄ (forsterite)

Fe₂SiO₄ (fayalite)

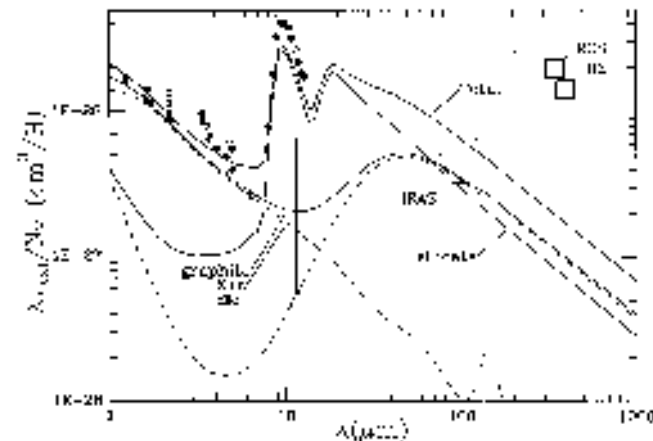
- SiO₂ (silica, quartz)

- **astronomical silicate** (MgFeSiO₄)

extinction curve



IR spectral feature



1-3. Expected composition of dust (2)

○ Minor candidates of dust composition

▪ Iron-bearing dust (Fe-based) ??

Fe (iron), FeO (wustite), Fe₂O₃ (hematite),
Fe₃O₄ (magnetite), FeS (troilite), FeS₂ (pyrite)

▪ Other carbides and oxides

SiC (silicon carbide), Al₂O₃ (corundum), MgO, TiC, ...

▪ Ices (appeared in MCs and YSOs)

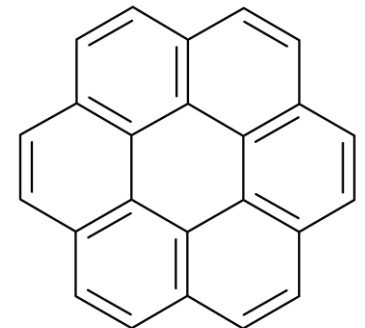
H₂O, CO, CO₂, NH₃, CH₄, CH₃OH, HCN, C₂H₂, ...

▪ large molecules

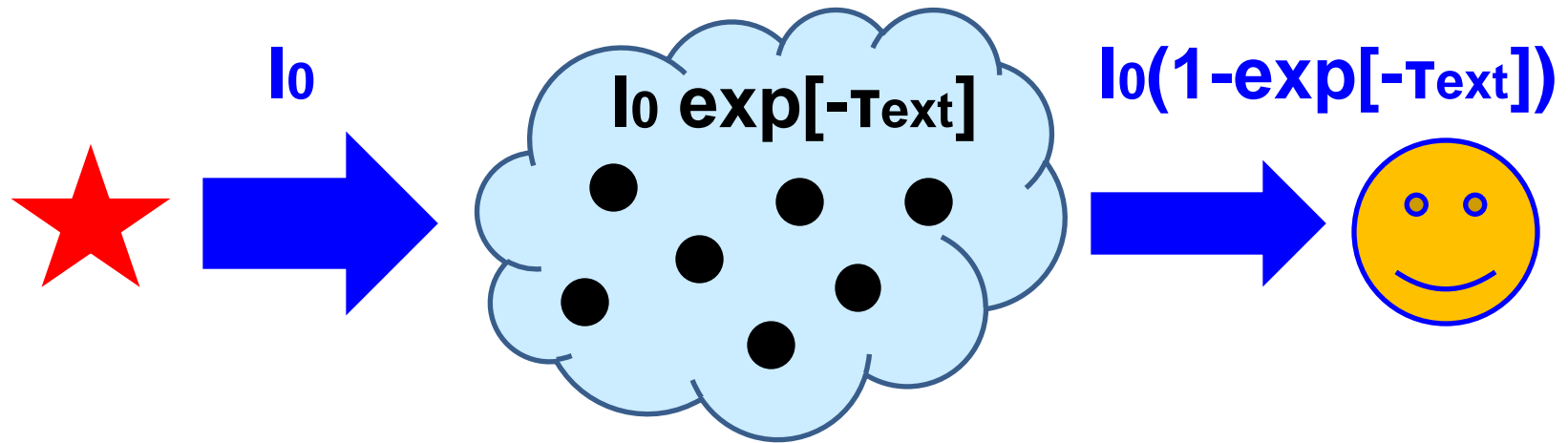
PAH (Polycyclic Aromatic Hydrocarbon):

C₂₄H₁₂ (coronene), etc.

HAC (Hydrogenated Amorphous Carbon)



2-1. Extinction by dust



○ Extinction = absorption + scattering

albedo = scattering / extinction : $w = T_{sca} / T_{ext}$

▪ extinction curve

optically thin, slab-like geometry

bright point source (OB stars, QSOs, GRB afterglow)

▪ attenuation curve

effective extinction including effects of radiative

transfer (Calzetti law for galaxies, **Calzetti+1994, 2000**)

2-2. Optical depth by dust

○ Optical depth produced by dust with a radius “a”

$$\tau_{\text{ext}}(\mathbf{a}, \lambda) = \int d\mathbf{r} n_{\text{dust}}(\mathbf{a}, \mathbf{r}) C_{\text{ext}}(\mathbf{a}, \lambda)$$

$C_{\text{ext}}(\mathbf{a}, \lambda)$: cross section of dust extinction

$$\tau_{\text{ext}}(\mathbf{a}, \lambda) = \int d\mathbf{r} n_{\text{dust}}(\mathbf{a}, \mathbf{r}) \pi a^2 Q_{\text{ext}}(\mathbf{a}, \lambda)$$

$Q_{\text{ext}}(\mathbf{a}, \lambda) = C_{\text{ext}}(\mathbf{a}, \lambda) / \pi a^2$: extinction coefficient

$Q_{\text{ext}}(\mathbf{a}, \lambda) = Q_{\text{abs}}(\mathbf{a}, \lambda) + Q_{\text{sca}}(\mathbf{a}, \lambda)$, $w = Q_{\text{sca}}(\mathbf{a}, \lambda) / Q_{\text{ext}}(\mathbf{a}, \lambda)$

○ Total optical depth produced by dust

$$\tau_{\text{ext}, \lambda} = \int d\mathbf{a} \int d\mathbf{r} f_{\text{dust}}(\mathbf{a}, \mathbf{r}) \pi a^2 Q_{\text{ext}}(\mathbf{a}, \lambda)$$

$f_{\text{dust}}(\mathbf{a}) = dn_{\text{dust}}(\mathbf{a})/da$: size distribution of dust

(= number density of dust with radii between a and $a+da$)

$$A_{\lambda} = -2.5 \log_{10}(\exp[-\tau_{\text{ext}, \lambda}]) = 1.068 \tau_{\text{ext}, \lambda}$$

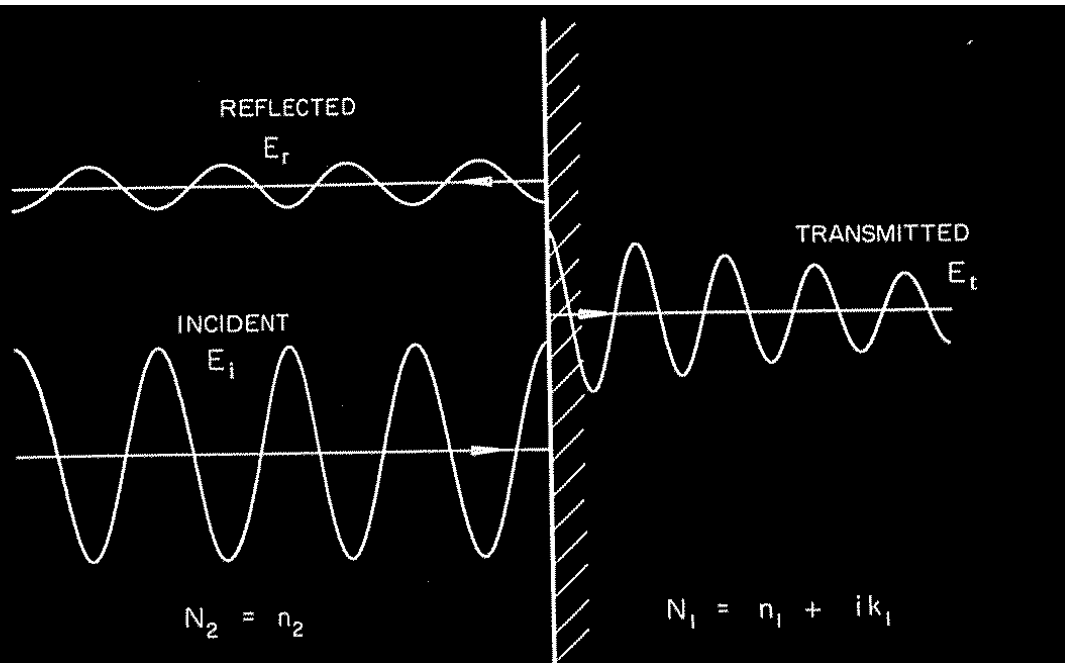
2-3. Evaluation of Q factors

How are Q-factors determined?

→ using the Mie (scattering) theory

○ Mie solution (Bohren & Huffmann 1983)

describes the scattering of electromagnetic radiation by a sphere (solving Maxwell equations)



refractive index

$$m(\lambda) = n(\lambda) + i k(\lambda)$$

n, k : optical constant

dielectric permeability

$$\epsilon = \epsilon_1 + i \epsilon_2$$

$$\epsilon_1 = n^2 - k^2$$

$$\epsilon_2 = 2nk$$

2-4. Behaviors of Q factors

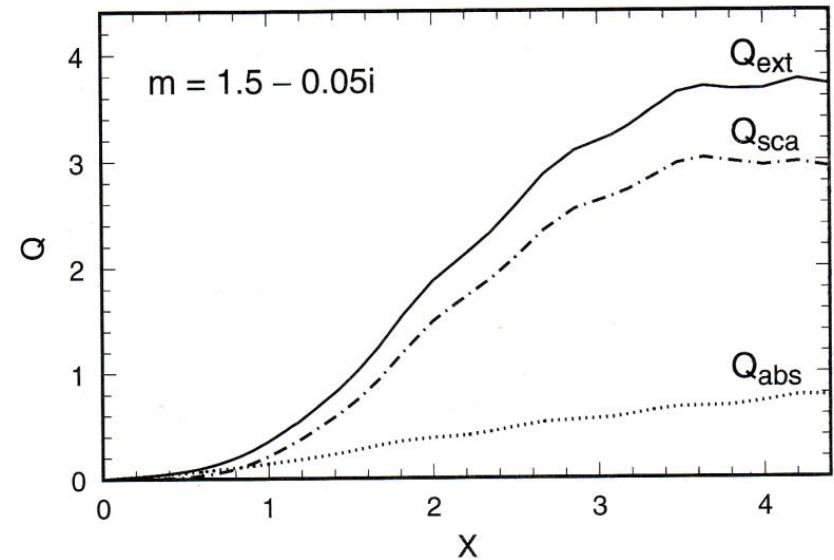
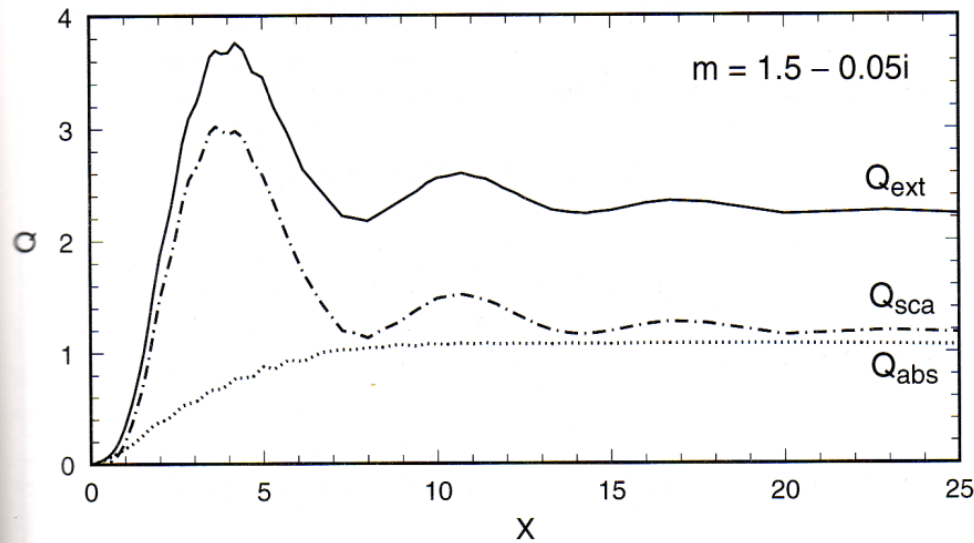
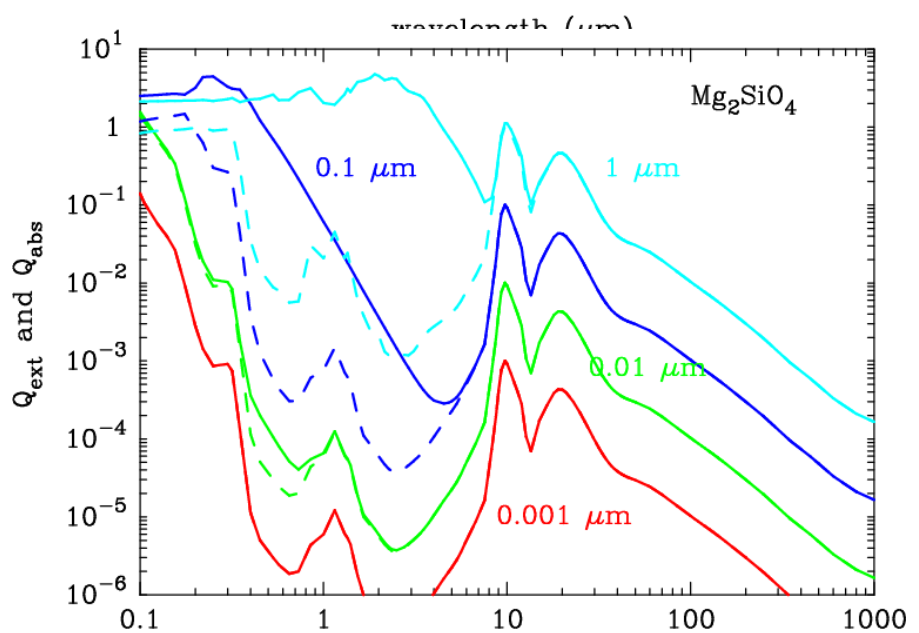
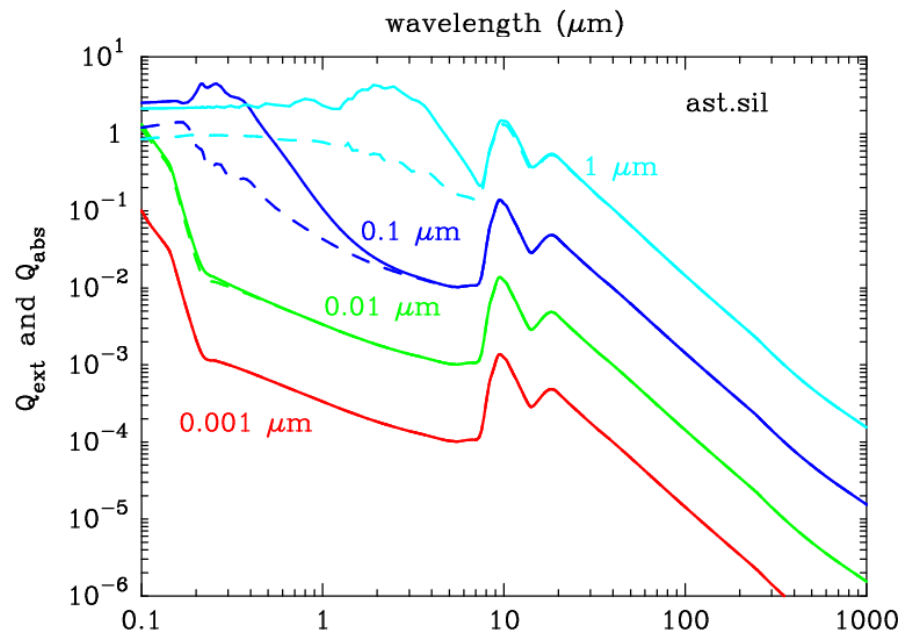
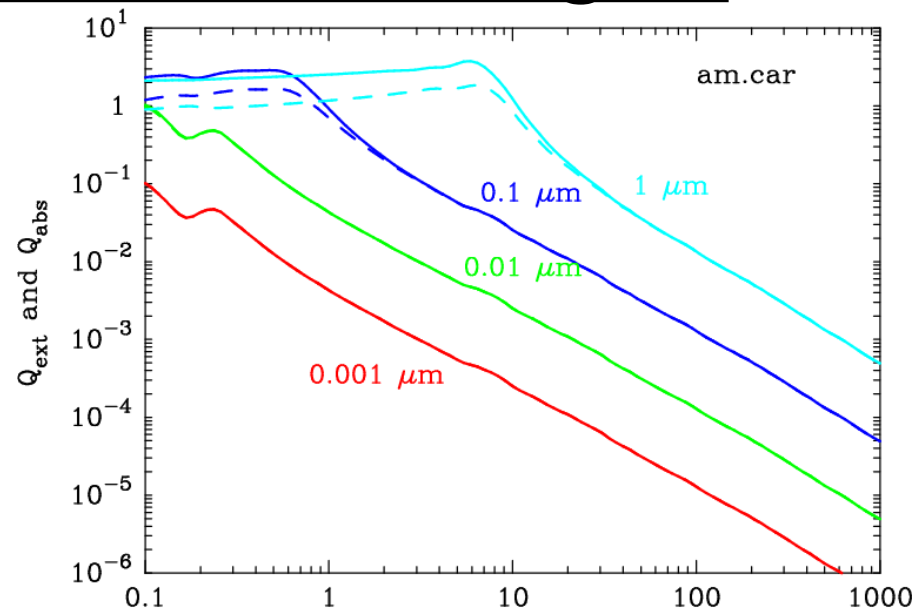
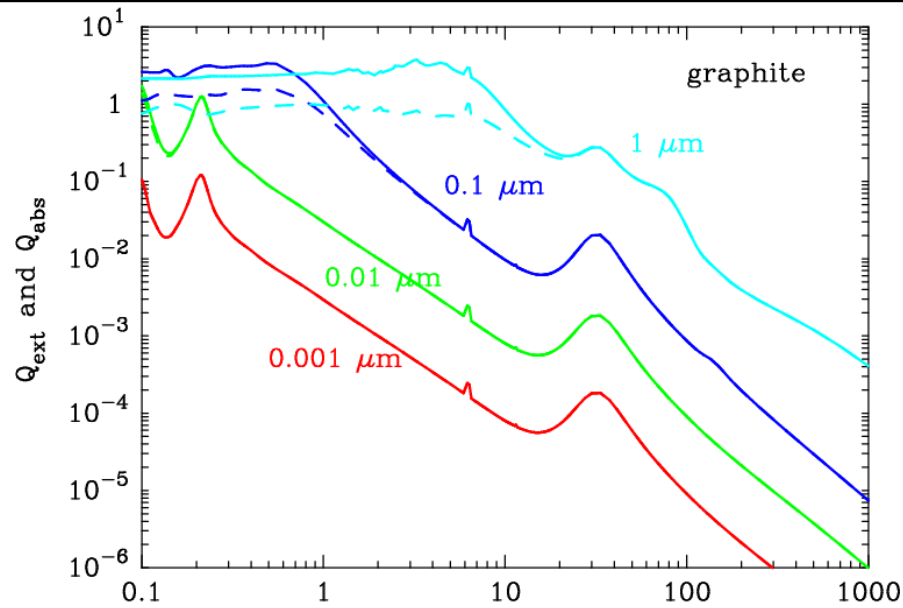


Figure 3.1. Results of Mie theory calculations for spherical grains of refractive index $m = 1.5 - 0.05i$. Efficiency factors Q_{ext} , Q_{sca} and Q_{abs} are plotted against the dimensionless size parameter $X = 2\pi a/\lambda$.

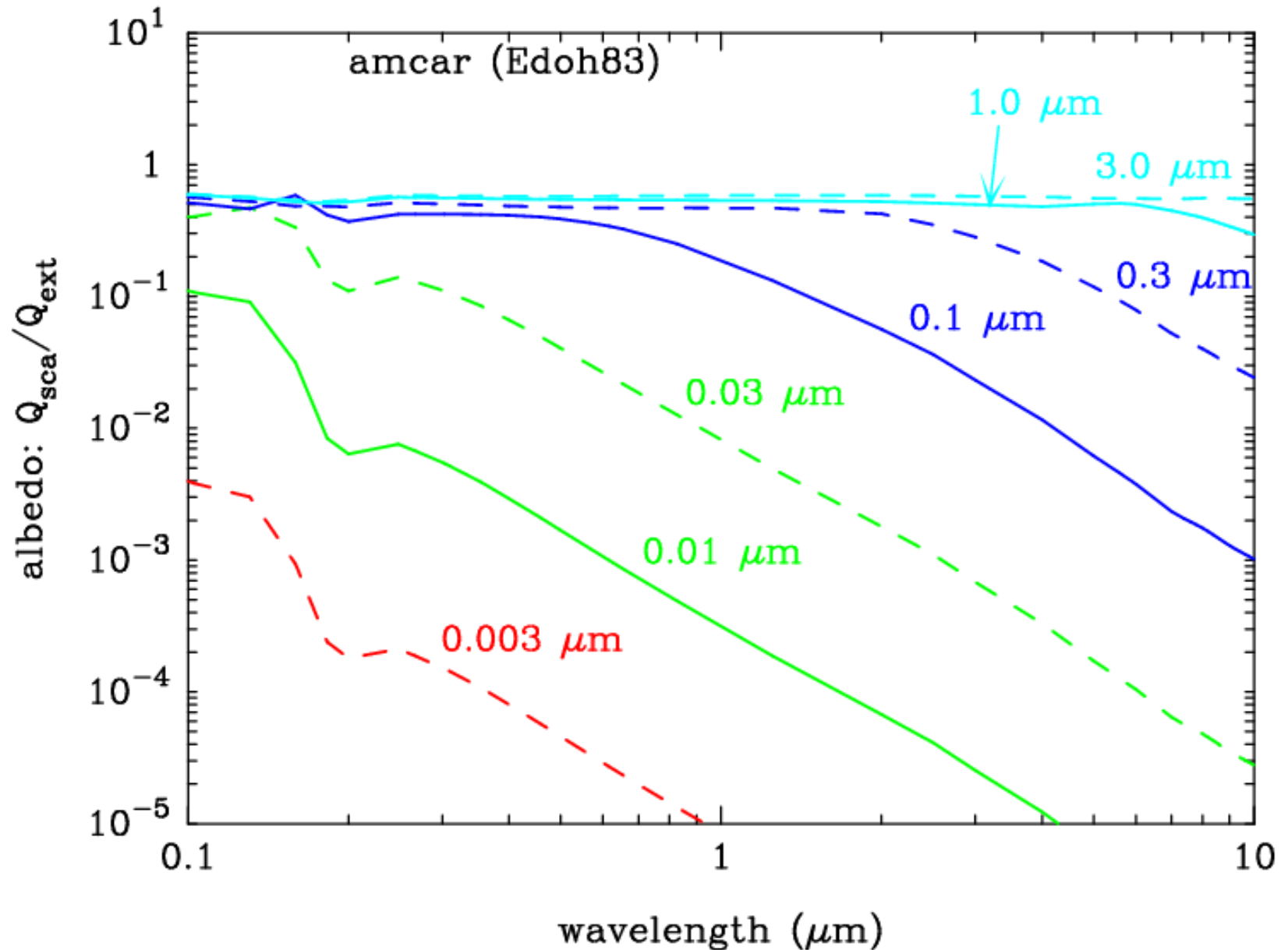
size parameter : $x = 2 \pi a / \lambda$

- $x \sim 3-4$ ($a \sim \lambda$) $\rightarrow Q_{sca}$ has a peak with $Q_{sca} > \sim 2$
- $x \gg 1$ ($a \gg \lambda$) $\rightarrow Q_{ext} = Q_{sca} + Q_{abs} \sim 2$
- $x \ll 1$ ($a \ll \lambda$, Rayleigh limit)
 - $\rightarrow Q_{sca} \propto x^4 \propto a^4$
 - $Q_{abs} \propto x \propto a$
 - $\rightarrow Q_{ext} = Q_{abs} \propto a$

2-5. Q factors as a function of wavelenghts

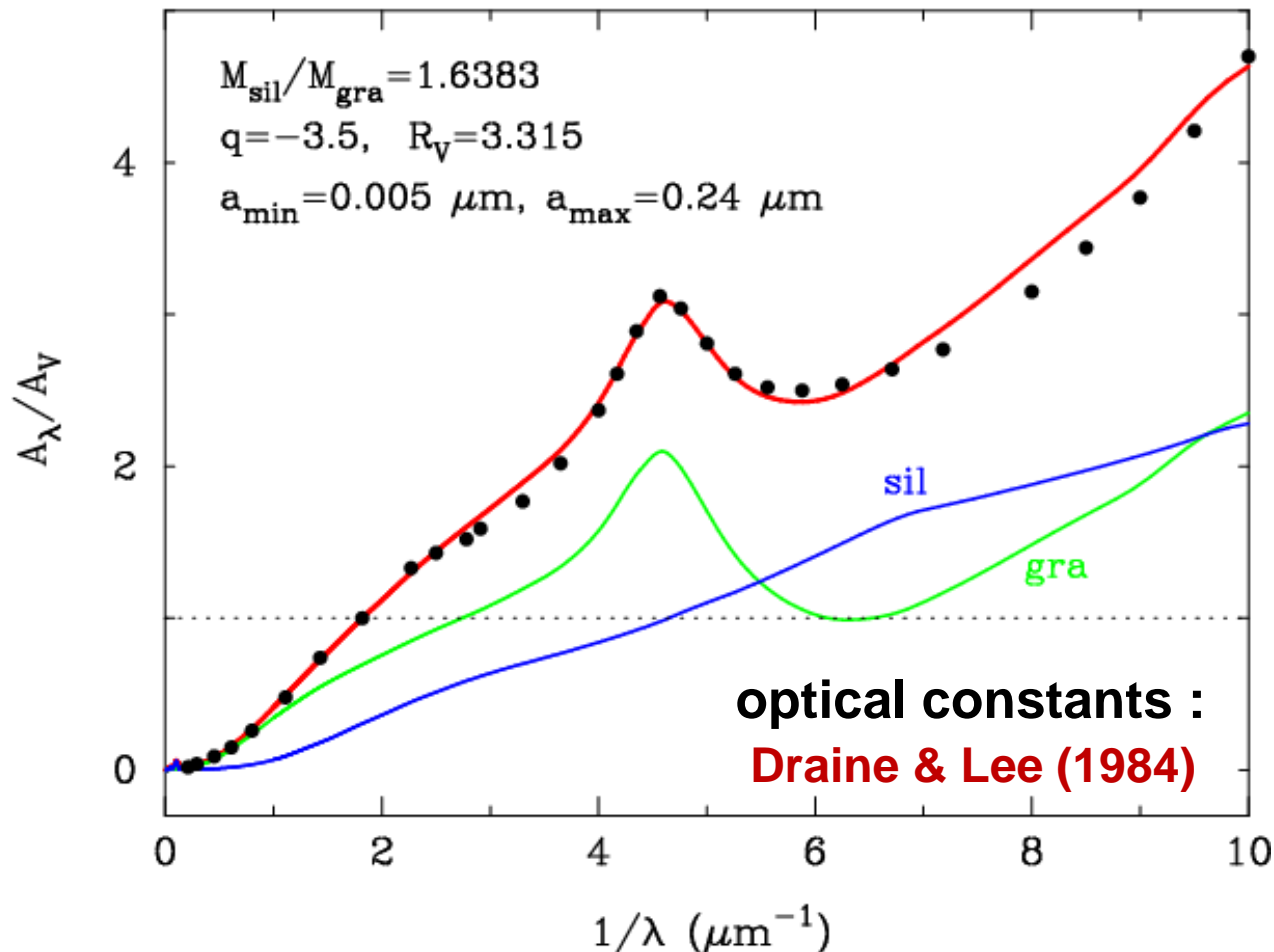


2-6. Dependence of albedo on grain radius

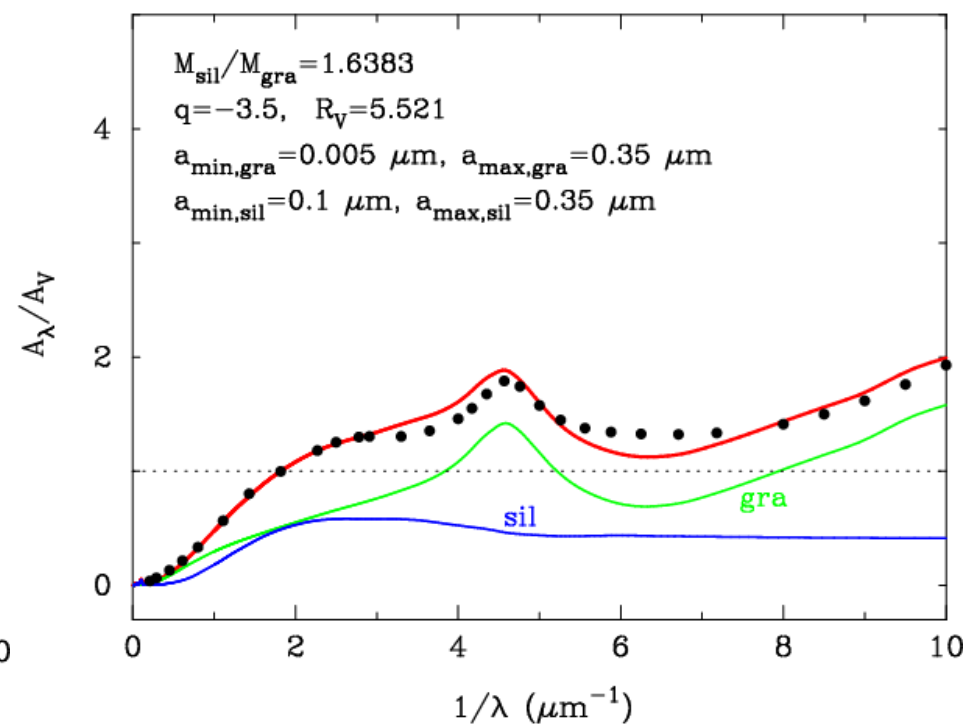
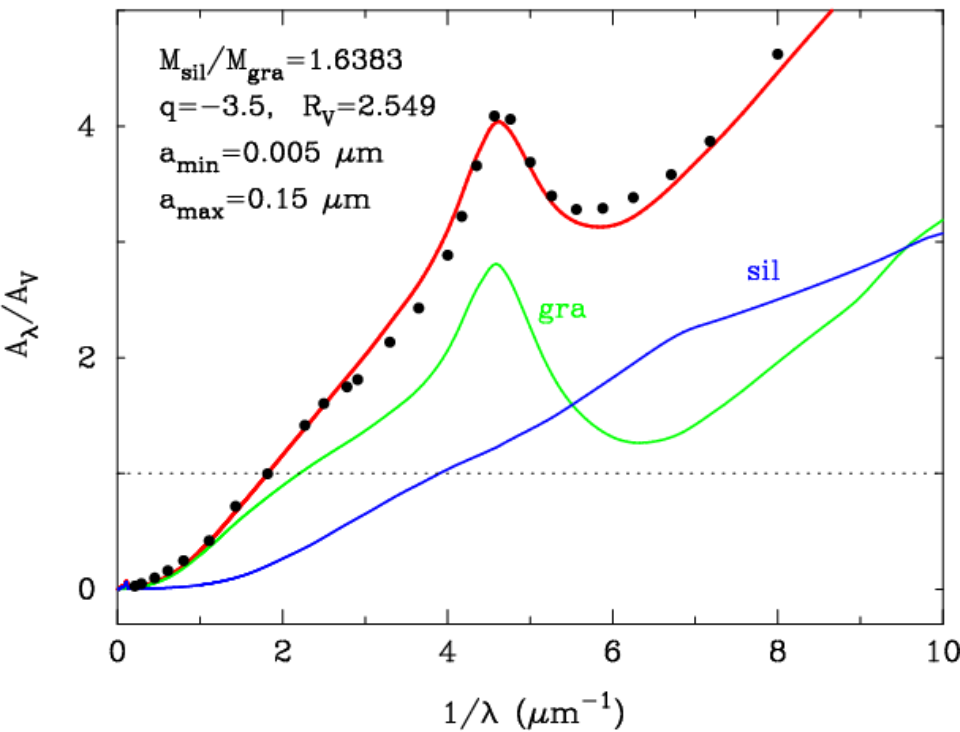


2-7. Classical interstellar dust model in MW

- MRN dust model (Mathis, Rumpl, & Nordsieck 1977)
 - dust composition : silicate (MgFeSiO₄) & graphite (C)
 - size distribution : power-law distribution
 $n(a) \propto a^{-q}$ with $q=3.5$, $0.005 \mu\text{m} \leq a \leq 0.25 \mu\text{m}$

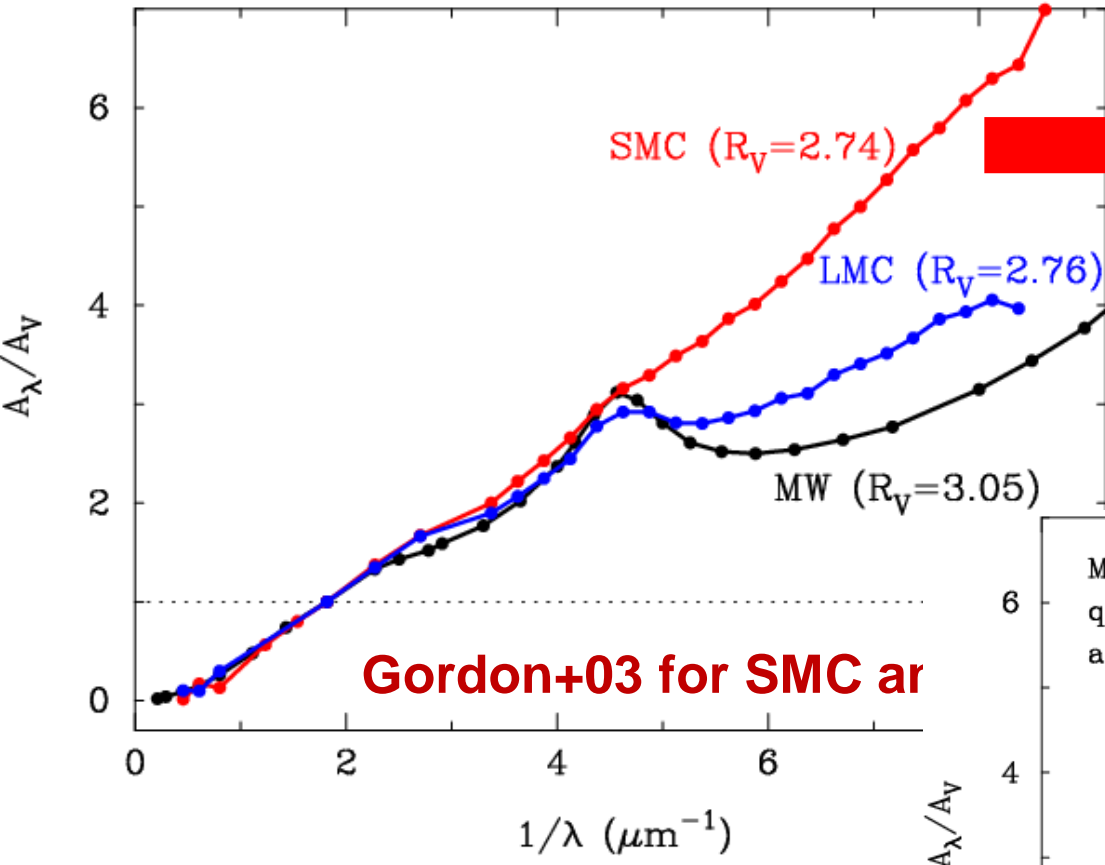


2-8. Dependence of extinction on grain size



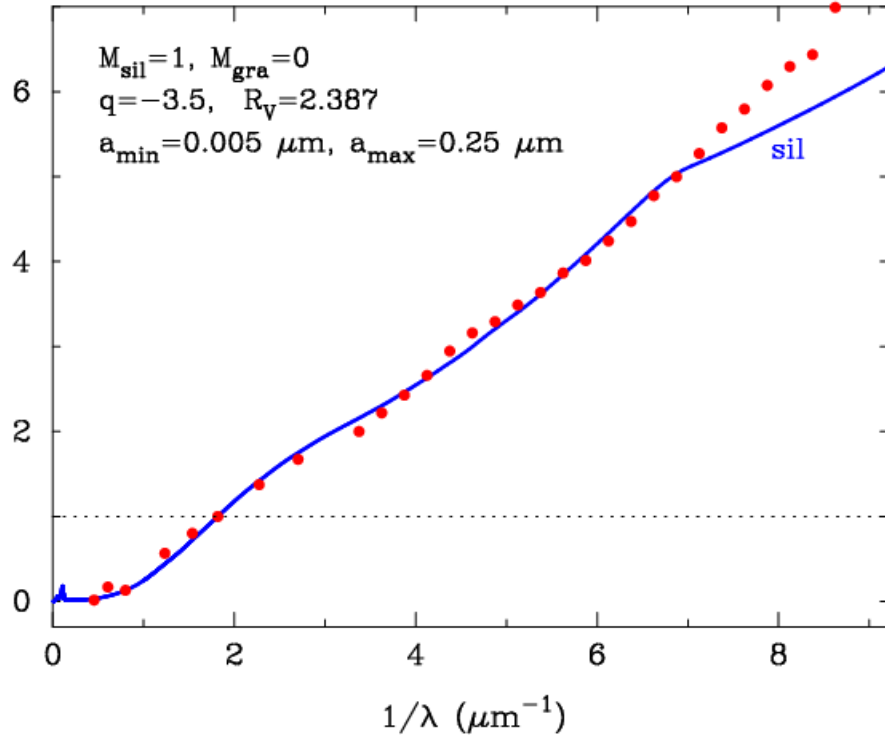
- R_V provides a rough estimate of grain size
 - smaller $R_V \rightarrow$ steeper curve \rightarrow smaller grain radius
 - larger $R_V \rightarrow$ flatter curve \rightarrow larger grain radius

2-9. Extinction curves in Magellanic Clouds



no 2175 A bump
steep far-UV rise

SMC extinction curve can be explained by the MRN model without graphite (Pei 1992)



- SMC extinction curve
- implying different properties
- being used as an ideal analog
- cf. extinction curves of red

3-1. Thermal emission from dust grains

○ Luminosity density emitted by dust grains

$$\begin{aligned}L_{\lambda}(\mathbf{a}) &= 4 N_{\text{dust}}(\mathbf{a}) C_{\text{emi}}(\mathbf{a},\lambda) \pi B_{\lambda}(T_{\text{dust}}[\mathbf{a}]) \\ &= 4\pi a^2 N_{\text{dust}}(\mathbf{a}) Q_{\text{abs}}(\mathbf{a},\lambda) \pi B_{\lambda}(T_{\text{dust}}[\mathbf{a}])\end{aligned}$$

$$C_{\text{emi}}(\mathbf{a},\lambda) = \pi a^2 Q_{\text{abs}}(\mathbf{a},\lambda)$$

Kirchhoff law: $Q_{\text{emi}}(\mathbf{a},\lambda) = Q_{\text{abs}}(\mathbf{a},\lambda)$

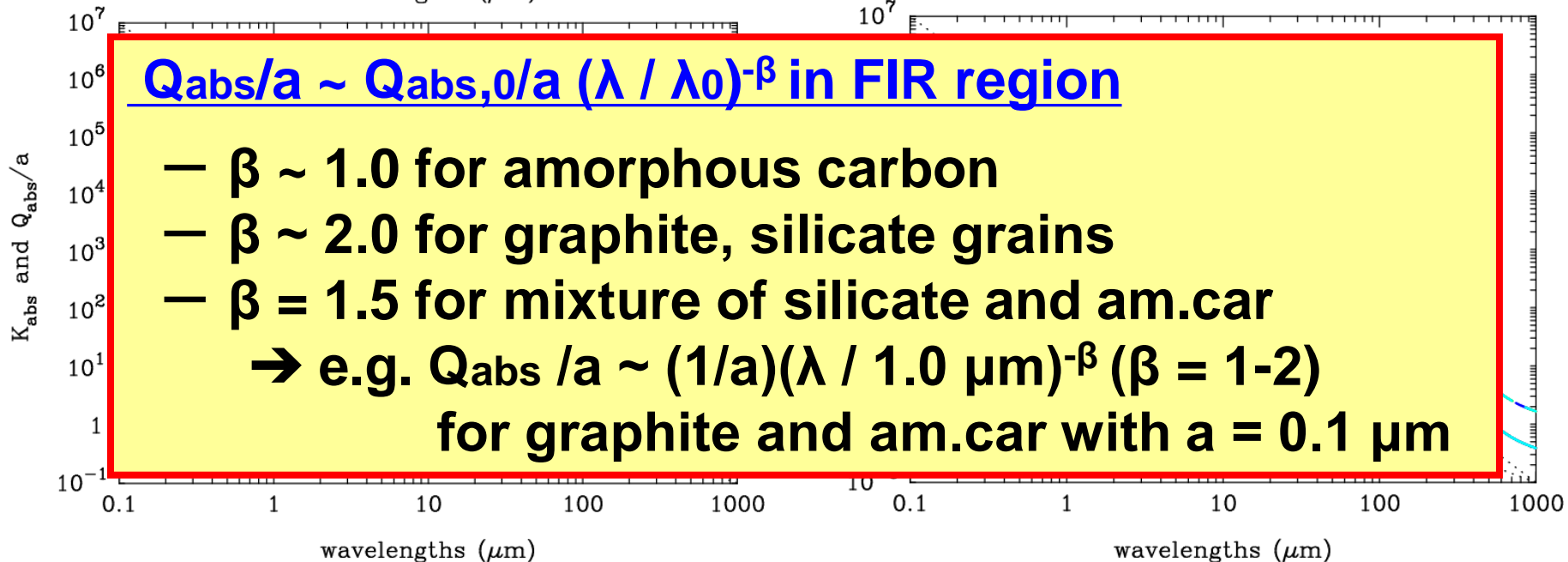
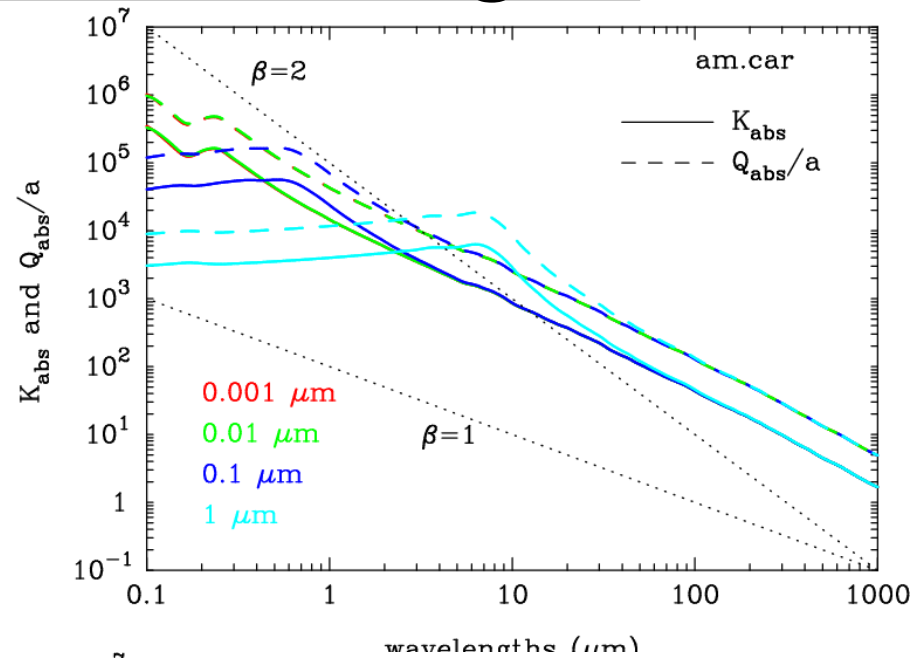
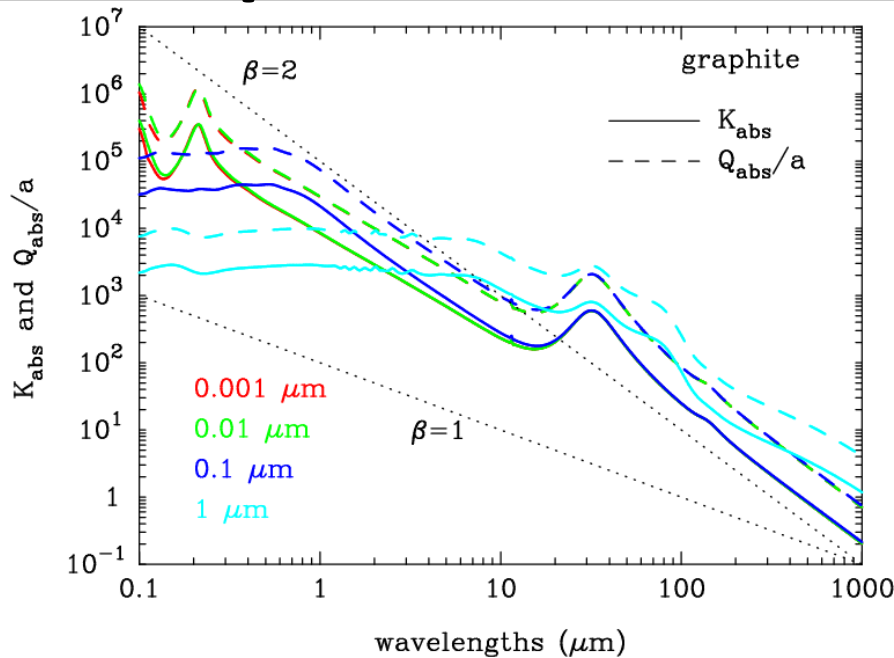
○ at IR wavelengths ($Q_{\text{abs}} \propto a$ for $a \ll \lambda$)

$$\begin{aligned}L_{\lambda}(\mathbf{a}) &= 4 N_{\text{dust}}(\mathbf{a}) (4\pi\rho a^3/3) (3Q_{\text{abs}}[\mathbf{a},\lambda]/4\rho a) \\ &\quad \pi B_{\lambda}(T_{\text{dust}}[\mathbf{a}]) \\ &= 4 M_{\text{dust}} K_{\text{abs}}(\lambda) \pi B_{\lambda}(T_{\text{dust}}[\mathbf{a}])\end{aligned}$$

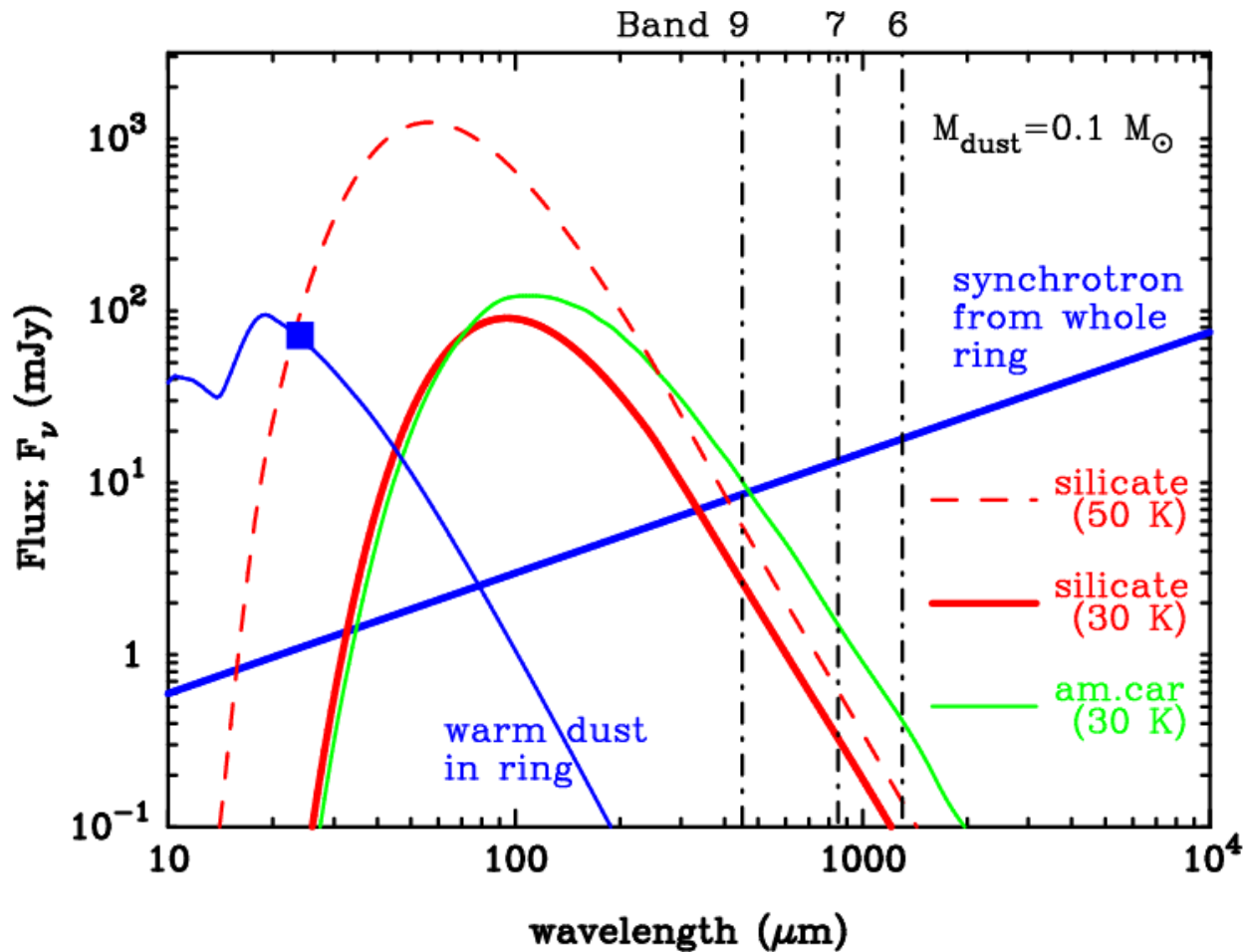
$K_{\text{abs}}(\mathbf{a},\lambda) = 3Q_{\text{abs}}/4\rho a$: mass absorption coefficient

→ IR emission is derived given M_{dust} , K_{abs} , and T_{dust}

3-2. Dependence of Q/a on wavelengths



3-3. Example of thermal emission from dust

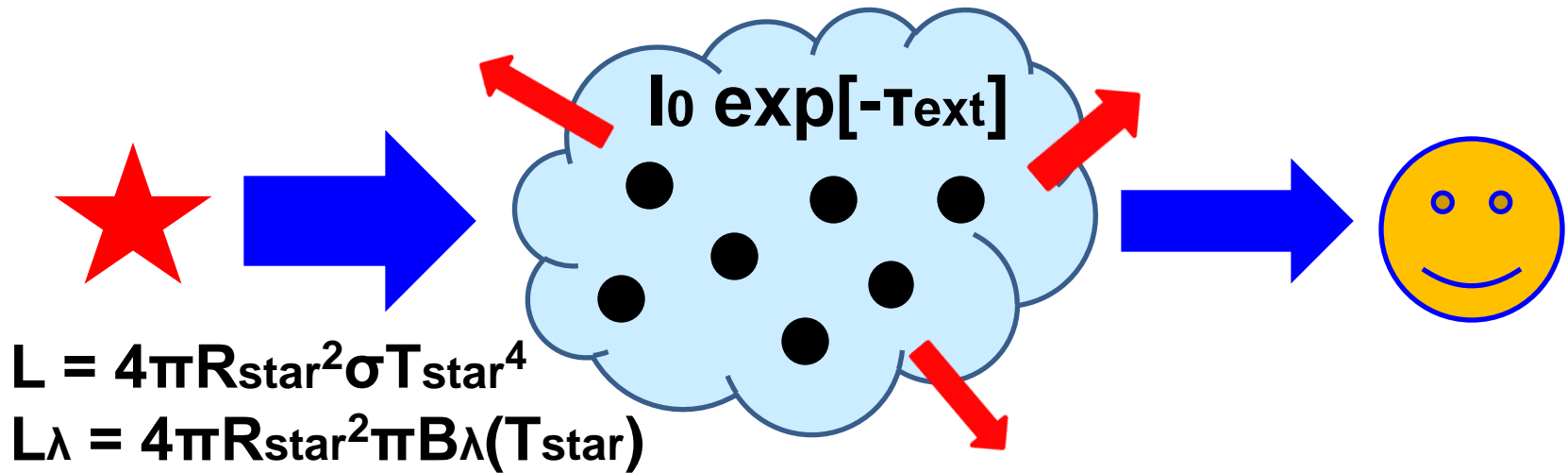


○ at IR wavelengths ($Q_{\text{abs}} \propto a$ for $a \ll \lambda$)

$$L_\lambda(a) = 4 M_{\text{dust}} K_{\text{abs}}(\lambda) \pi B_\lambda(T_{\text{dust}}[a])$$

→ IR emission is derived given M_{dust} , K_{abs} , and T_{dust}

3-4. Reemission of stellar light by dust



How is dust temperature determined?

→ approaching from conservation of energy

○ energy reemitted by dust = energy absorbed by dust

time during which dust is thermalized

extremely short $< 10^{-5}$ s

3-5. Luminosity of dust emission

○ Luminosity emitted by a dust grain

$$\begin{aligned} L_{\text{emi}} &= \int d\lambda \, 4\pi a^2 Q_{\text{abs}}(a, \lambda) \pi B_{\lambda}(T_{\text{dust}}[a]) \\ &= 4\pi a^2 \sigma T_{\text{dust}}^4 \langle Q_{\text{abs}}(T_{\text{dust}}) \rangle \end{aligned}$$

$\langle Q_{\text{abs}}(T_{\text{dust}}) \rangle$: plank-mean absorption coefficient
 $= \int Q_{\text{abs}}(a, \lambda) \pi B_{\lambda}(T_{\text{dust}}[a]) d\lambda / \int \pi B_{\lambda}(T_{\text{dust}}[a]) d\lambda$

○ Stellar luminosity absorbed by a dust grain

$$\begin{aligned} L_{\text{abs}} &= \int d\lambda \, \pi a^2 Q_{\text{abs}}(a, \lambda) F_{\lambda}(T_{\text{star}}) \\ &= (R_{\text{star}}/D)^2 \pi a^2 \sigma T_{\text{star}}^4 \langle Q_{\text{abs}}(T_{\text{star}}) \rangle \\ &= L_{\text{star}}/4\pi D^2 \pi a^2 \langle Q_{\text{abs}}(T_{\text{star}}) \rangle \end{aligned}$$

a star is assumed to be a blackbody:

$$F_{\lambda}(T_{\text{star}}) = L_{\lambda}/4\pi D^2 = (R_{\text{star}}/D)^2 \pi B_{\lambda}(T_{\text{star}})$$

3-6. Temperature of dust grains (1)

○ $L_{\text{abs}} = L_{\text{emi}}$ (dust is assumed to be in thermal equilibrium)

$$(R_{\text{star}}/D)^2 T_{\text{star}}^4 \langle Q_{\text{abs}}(T_{\text{star}}) \rangle = 4T_{\text{dust}}^4 \langle Q_{\text{abs}}(T_{\text{dust}}) \rangle$$

or

$$L_{\text{star}}/4\pi D^2 \langle Q_{\text{abs}}(T_{\text{star}}) \rangle = 4\sigma T_{\text{dust}}^4 \langle Q_{\text{abs}}(T_{\text{dust}}) \rangle$$

○ analytical solutions, assuming $\langle Q_{\text{abs}}(T_{\text{star}}) \rangle = 1$ and

$$Q_{\text{abs}} = Q_{\text{abs},0} (\lambda / \lambda_0)^{-\beta}$$

$$\rightarrow \langle Q_{\text{abs}}(T_{\text{dust}}) \rangle = Q_{\text{abs},0} \lambda_0^\beta (15/\pi^4) (kT_{\text{dust}}/hc)^\beta \zeta(\beta+4) \Gamma(\beta+4)$$

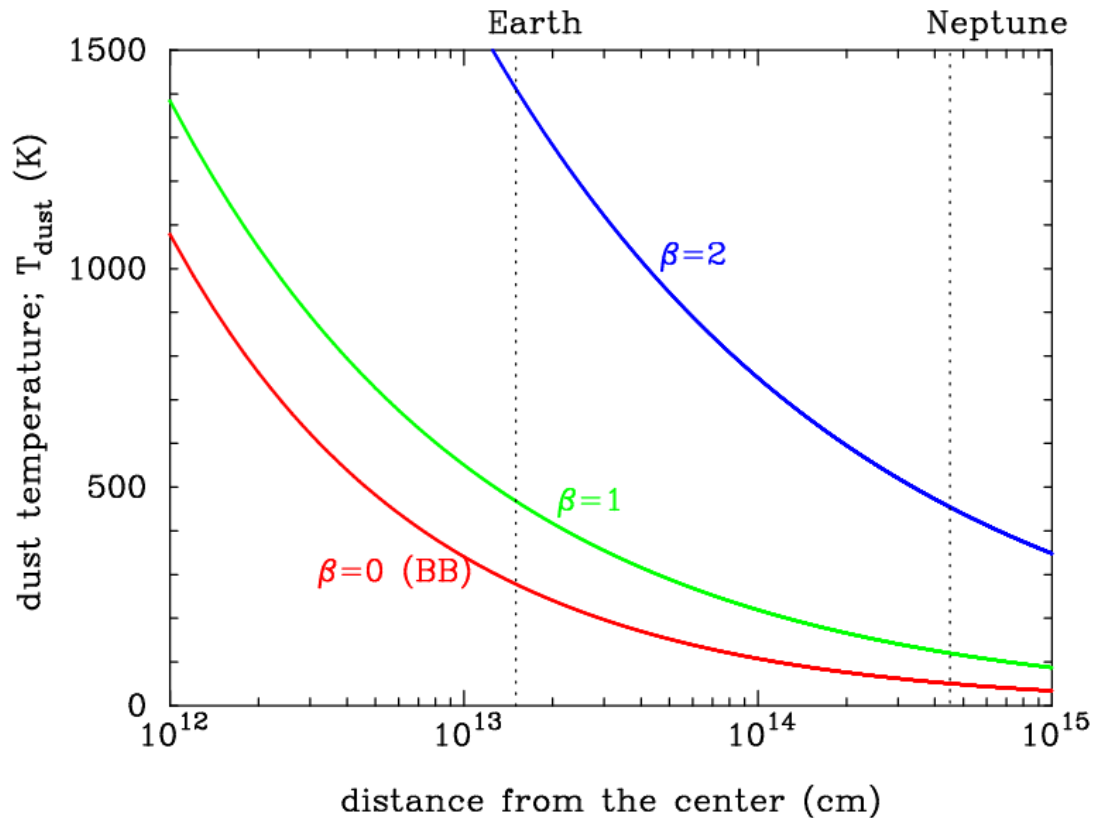
$$(\beta=0) \quad L_{\text{star}}/4\pi D^2 = 4Q_{\text{abs},0} \sigma T_{\text{dust}}^4$$

$$(\beta=1) \quad L_{\text{star}}/4\pi D^2 = 4Q_{\text{abs},0} (k\lambda_0/hc) \sigma T_{\text{dust}}^5 (373.3/\pi^4)$$

$$(\beta=2) \quad L_{\text{star}}/4\pi D^2 = 4Q_{\text{abs},0} (k\lambda_0/hc)^2 \sigma T_{\text{dust}}^6 (\pi^2/63)$$

3-7. Temperature of dust grains (2)

- Dust temperatures as functions of D (distance) and β for $L_{\text{star}} = L_{\text{sun}} = 3.85 \times 10^{33} \text{ erg s}^{-1}$, $Q_{\text{abs},0} = 1$, $\lambda_0 = 1.0 \text{ } \mu\text{m}$



Dust temperature is higher for higher β
→ Higher β causes lower efficiency of radiative cooling