

# 超新星放出ガス中でのダスト形成 と衝撃波中でのダスト破壊

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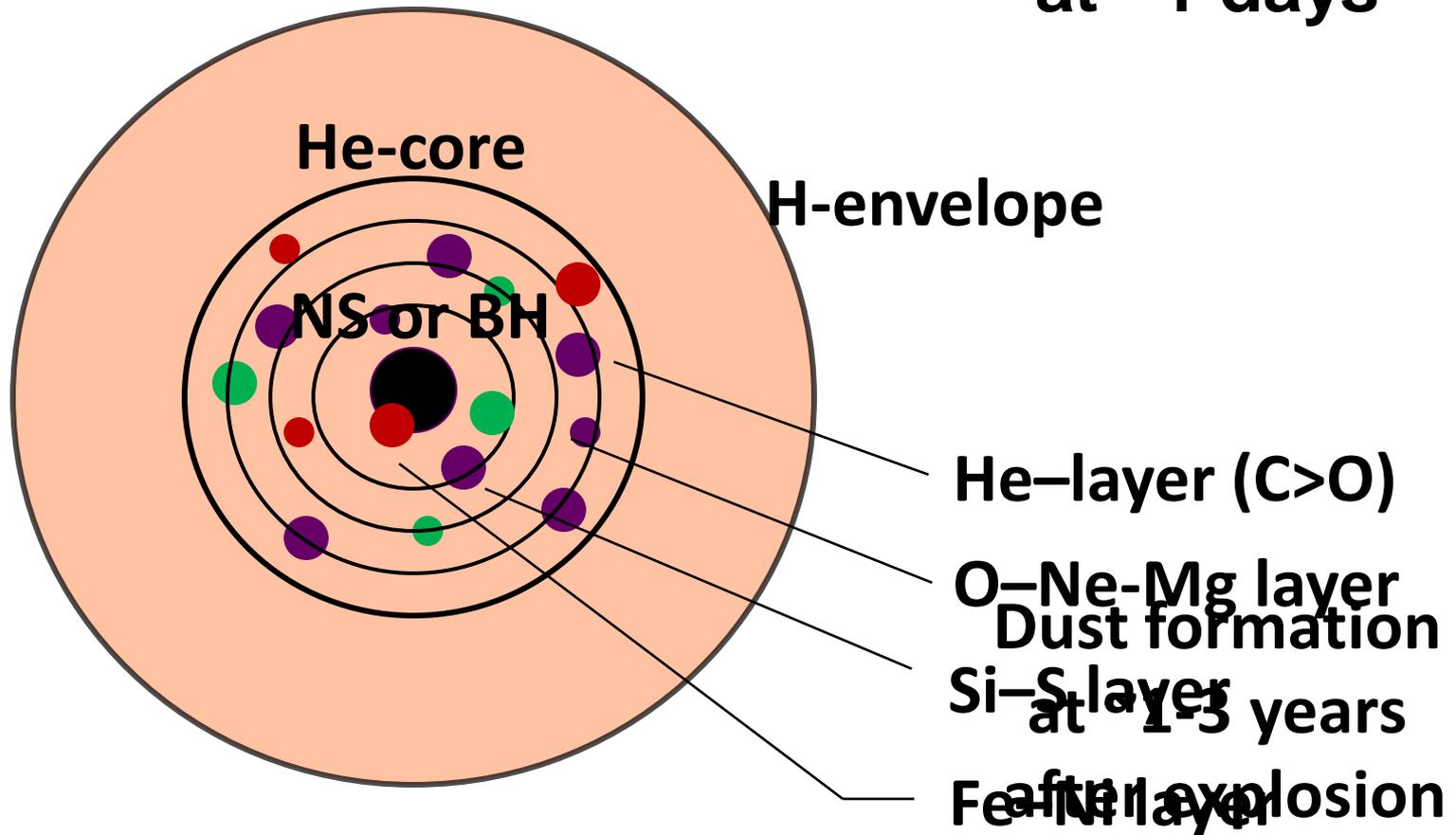
N. Tominaga (Konan Univ.)

# 1. Introduction

- supernovae (SNe) : explosions of massive (early type) stars with  $M_{\text{ZAMS}} > 8 M_{\text{sun}}$
- SNe are important sources of interstellar dust?
  - huge amounts of dust grains ( $>10^8 M_{\text{sun}}$ ) are detected in host galaxies of quasars at redshift  $z > 5$ 
    - $0.1 M_{\text{sun}}$  of dust per SN is needed to explain such massive dust at high- $z$  (e.g. Dwek et al. 2007)
  - contribution of dust mass from AGB stars and SNe
    - $n(\text{AGB stars}) / n(\text{SNe}) \sim 10\text{-}20$
    - $M_{\text{dust}} = 0.01\text{-}0.05 M_{\text{sun}}$  per AGB (Zhukovska & Gail 2008)
    - $M_{\text{dust}} = 0.1\text{-}1.0 M_{\text{sun}}$  per SN (Nozawa et al. 2003; 2007)

# 2. Dust Formation in Pop III SNe

at ~1 days

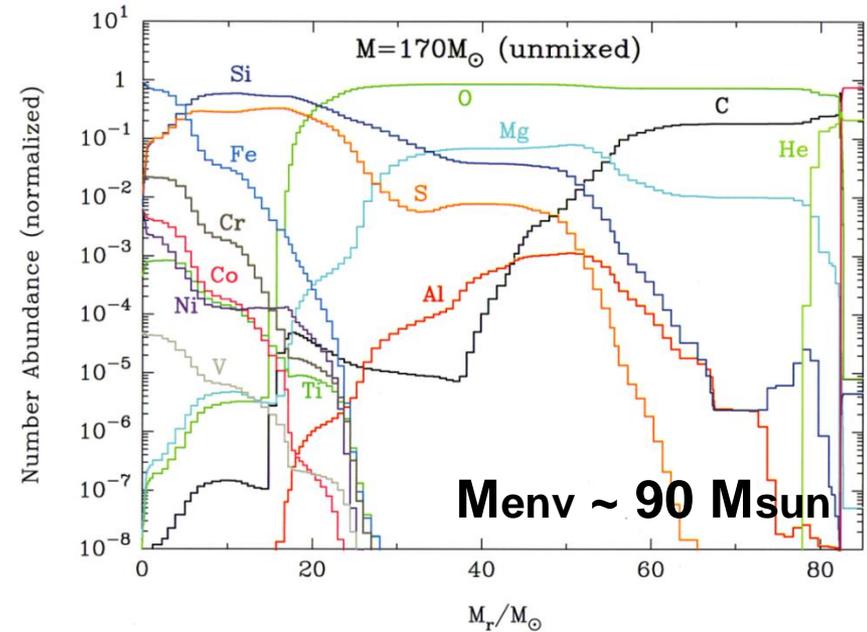
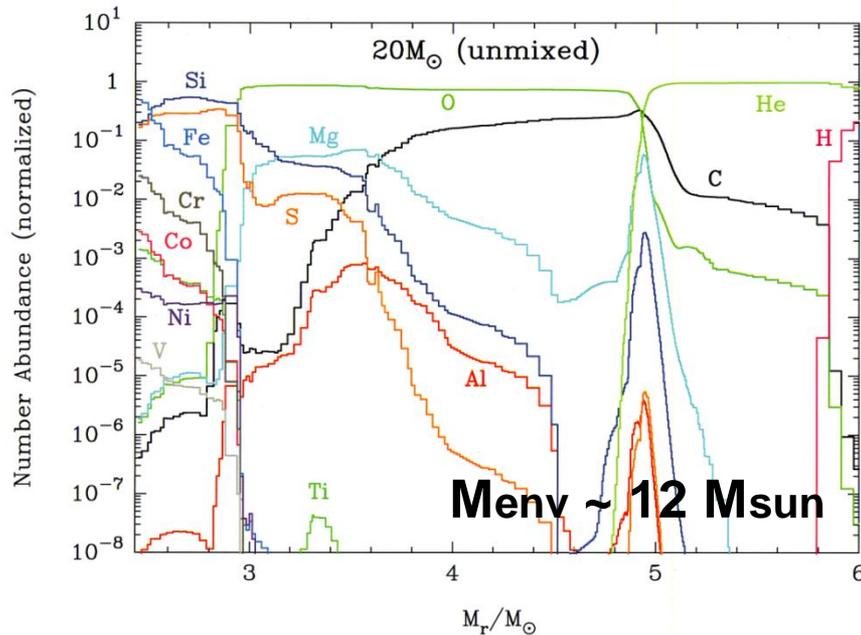


# 2-1. Dust formation in primordial SNe

Nozawa+'03, ApJ, 598, 785

## ○ Population III SNe model (Umeda & Nomoto'02)

- SNe II-P :  $M_{ZAMS} = 13, 20, 25, 30 M_{\text{sun}}$  ( $E_{51}=1$ )
- PISNe :  $M_{ZAMS} = 170 M_{\text{sun}}$  ( $E_{51}=20$ ),  $200 M_{\text{sun}}$  ( $E_{51}=28$ )



- nucleation and grain growth theory (Kozasa & Hasegawa'88)
- complete and no mixing of elements within the He-core
- complete formation of CO and SiO

# 2-2-1. Calculations of dust formation

- nucleation and grain growth theory taking account of chemical reaction at condensation

(Kozasa & Hasegawa'87)

- key species :  
gas species with the  
least collision frequency  
among reactants



key species controls the  
kinetics of the nucleation  
and grain growth

Dust species	Chemical reactions
Fe <sub>(s)</sub>	Fe <sub>(g)</sub> → Fe <sub>(s)</sub>
FeS <sub>(s)</sub>	Fe <sub>(g)</sub> + S <sub>(g)</sub> → FeS <sub>(s)</sub>
Si <sub>(s)</sub>	Si <sub>(g)</sub> → Si <sub>(s)</sub>
Ti <sub>(s)</sub>	Ti <sub>(g)</sub> → Ti <sub>(s)</sub>
V <sub>(s)</sub>	V <sub>(g)</sub> → V <sub>(s)</sub>
Cr <sub>(s)</sub>	Cr <sub>(g)</sub> → Cr <sub>(s)</sub>
Co <sub>(s)</sub>	Co <sub>(g)</sub> → Co <sub>(s)</sub>
Ni <sub>(s)</sub>	Ni <sub>(g)</sub> → Ni <sub>(s)</sub>
Cu <sub>(s)</sub>	Cu <sub>(g)</sub> → Cu <sub>(s)</sub>
C <sub>(s)</sub>	C <sub>(g)</sub> → C <sub>(s)</sub>
SiC <sub>(s)</sub>	Si <sub>(g)</sub> + C <sub>(g)</sub> → SiC <sub>(s)</sub>
TiC <sub>(s)</sub>	Ti <sub>(g)</sub> + C <sub>(g)</sub> → TiC <sub>(s)</sub>
Al <sub>2</sub> O <sub>3(s)</sub>	2Al <sub>(g)</sub> + 3O <sub>(g)</sub> → Al <sub>2</sub> O <sub>3(s)</sub>
MgSiO <sub>3(s)</sub>	Mg <sub>(g)</sub> + SiO <sub>(g)</sub> + 2O <sub>(g)</sub> → MgSiO <sub>3(s)</sub>
Mg <sub>2</sub> SiO <sub>4(s)</sub>	2Mg <sub>(g)</sub> + SiO <sub>(g)</sub> + 3O <sub>(g)</sub> → Mg <sub>2</sub> SiO <sub>4(s)</sub>
SiO <sub>2(s)</sub>	SiO <sub>(g)</sub> + O <sub>(g)</sub> → SiO <sub>2(s)</sub>
MgO <sub>(s)</sub>	Mg <sub>(g)</sub> + O <sub>(g)</sub> → MgO <sub>(s)</sub>
Fe <sub>3</sub> O <sub>4(s)</sub>	3Fe <sub>(g)</sub> + 4O <sub>(g)</sub> → Fe <sub>3</sub> O <sub>4(s)</sub>
FeO <sub>(s)</sub>	Fe <sub>(g)</sub> + O <sub>(g)</sub> → FeO <sub>(s)</sub>

## 2-2-2. Nucleation rate of dust

### Steady-state nucleation rate

$$J_s(t) = \alpha_s \Omega \left( \frac{2\sigma}{\pi m_1} \right)^{\frac{1}{2}} \Pi c_1^2(t) \exp \left[ -\frac{4}{27} \frac{\mu^3}{(\ln S)^2} \right]$$

### Supersaturation ratio

$$\ln S = \ln \left( \frac{p_1}{\dot{p}_1} \right) = -\frac{1}{kT} (\dot{g}_s - \dot{g}_1) + \ln \left( \frac{p_1}{p_0} \right)$$

$\alpha_s$  : sticking probability of key species ( $\alpha_s = 1$ , in the calculations)

$\Omega$  : volume of the condensate per key species ( $\Omega = 4\pi a_0^3/3$ )

$\sigma$  : surface energy of the condensate

$m_1$  : mass of key species

$c_1(t)$  : number density of key species

$\mu$  :  $\mu \equiv 4\pi a_0^2 \sigma / kT$  ; energy barrier for nucleation

## 2-2-3. Basic equations for dust formation

### ▪ Equation of mass conservation

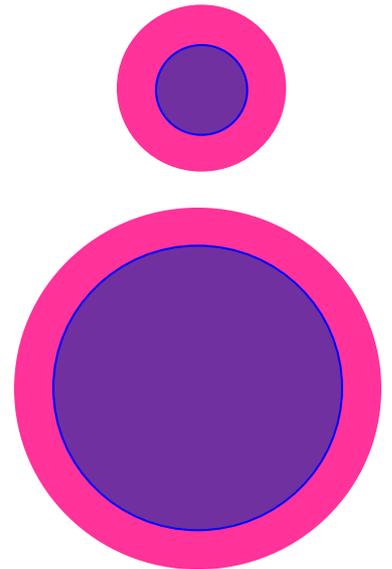
$$c_{10} - c_1 = \int_{t_0}^t J_{n_s}(t') \frac{a^3(t, t')}{a_0^3} dt',$$

### ▪ Equation of grain growth

$$\frac{da}{dt} = s\Omega_0 \left( \frac{kT}{2\pi m_1} \right)^{\frac{1}{2}} c_1 \left( 1 - \frac{1}{S} \right),$$



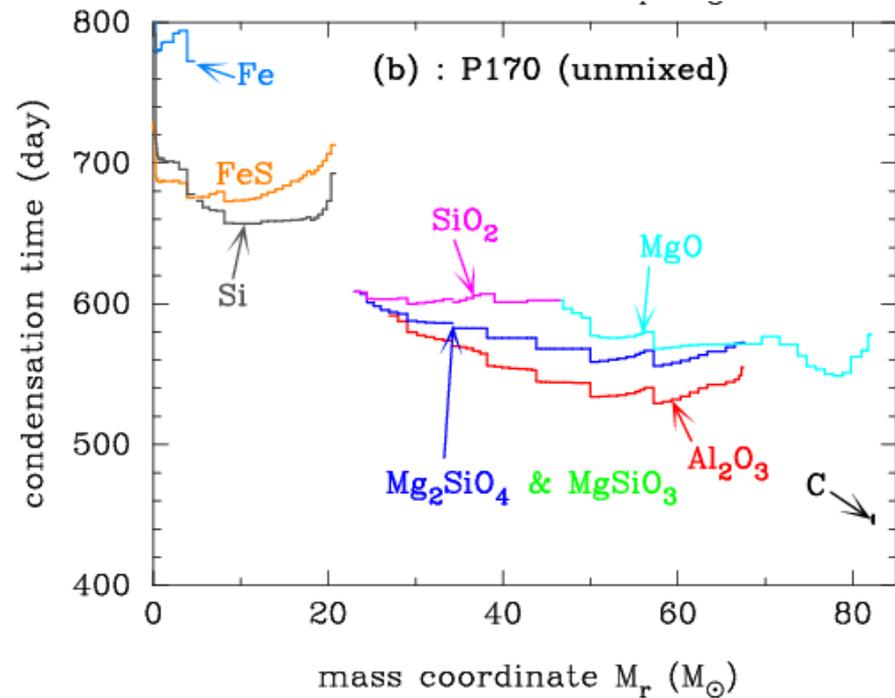
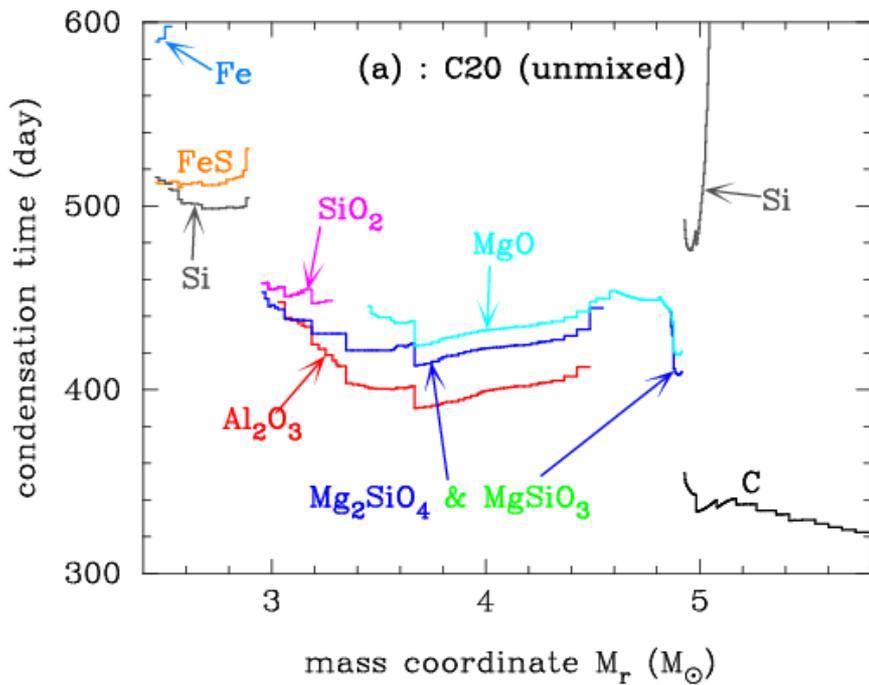
$$\frac{dV}{dt} = s\Omega_0 4\pi a^2 \left( \frac{kT}{2\pi m_1} \right)^{\frac{1}{2}} c_1 \left( 1 - \frac{1}{S} \right),$$



**Growth rate is independent of grain radius**

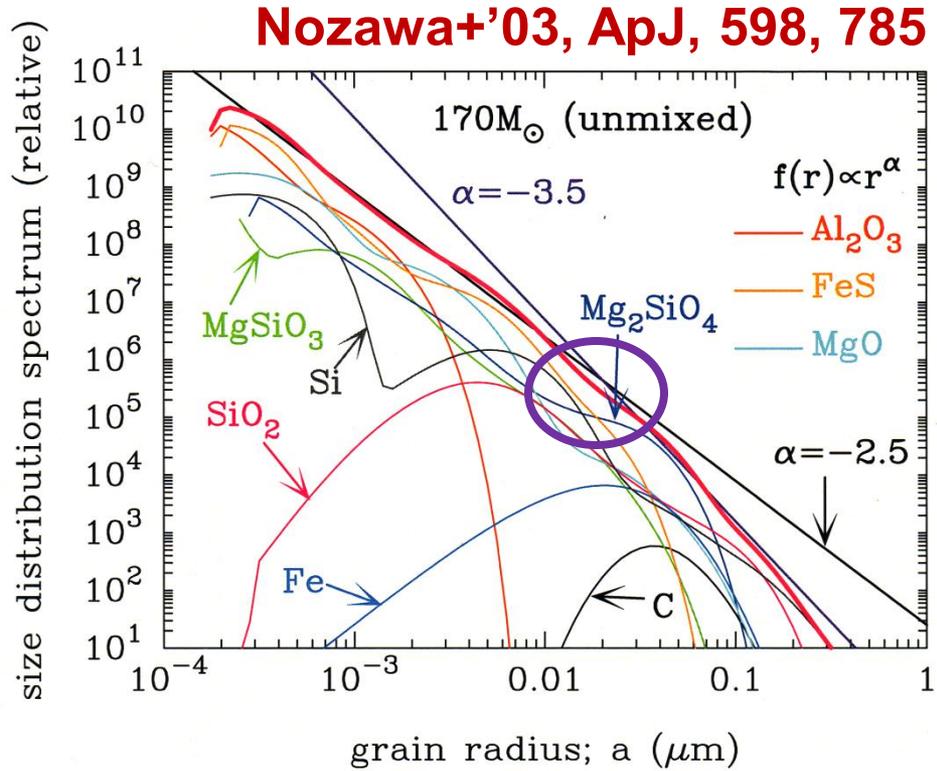
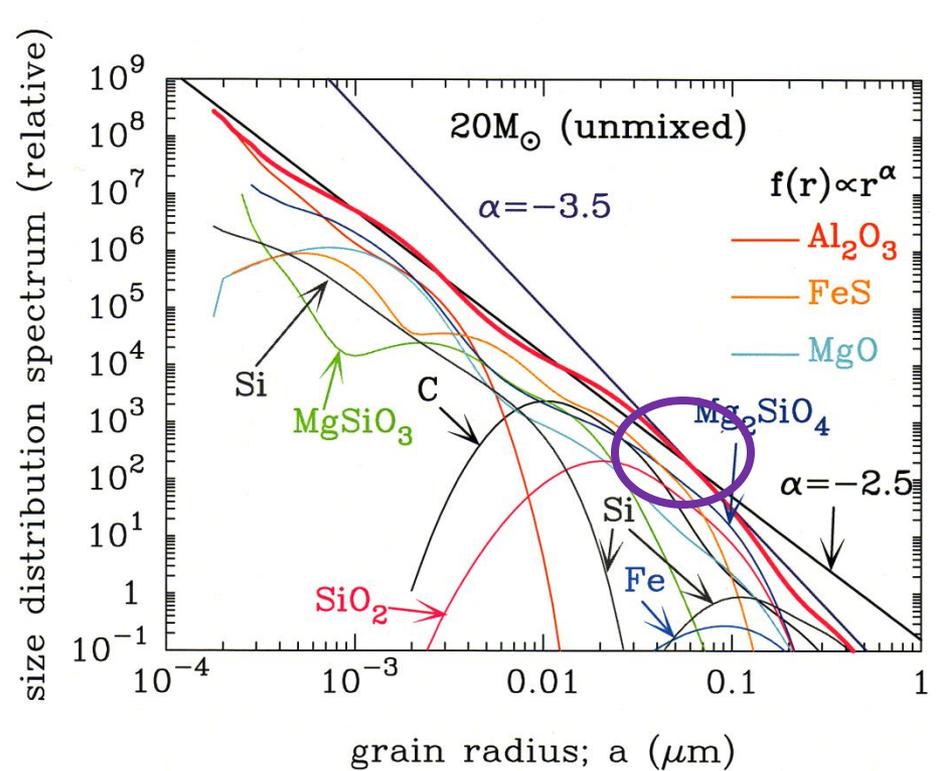
# 2-3-1. Dust formed in primordial SNe

Nozawa+'03, ApJ, 598, 785



- Various dust species (C,  $\text{MgSiO}_3$ ,  $\text{Mg}_2\text{SiO}_4$ ,  $\text{SiO}_2$ ,  $\text{Al}_2\text{O}_3$ , MgO, Si, FeS, Fe) form in the unmixed ejecta, according to the elemental composition of gas in each layer
- The condensation time: **300-600 days** for SNe II-P  
**400-800 days** for PISNe

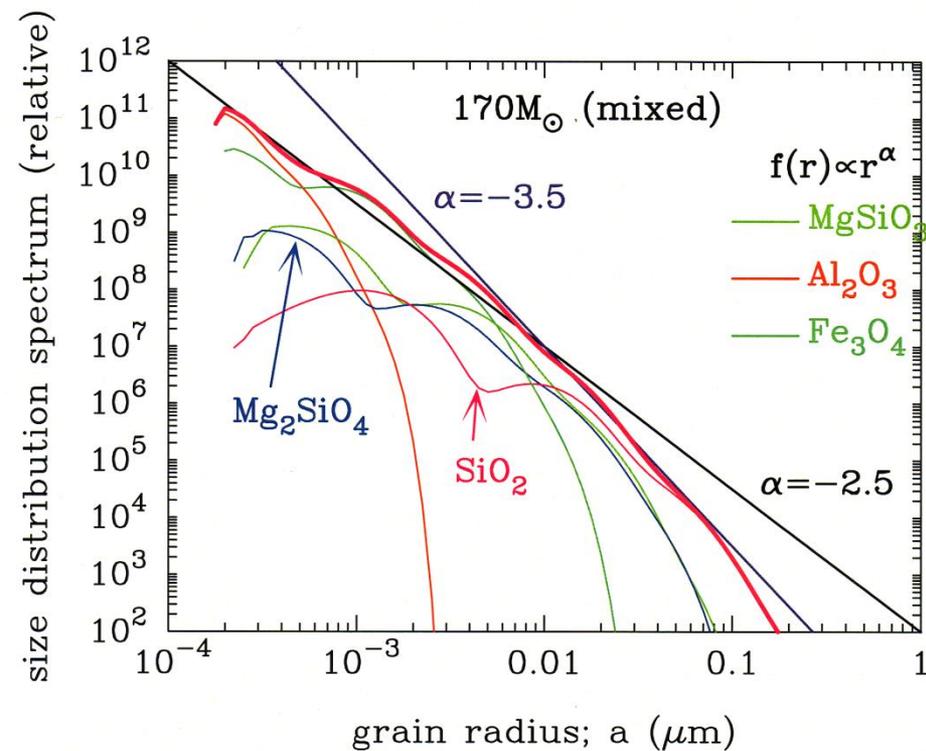
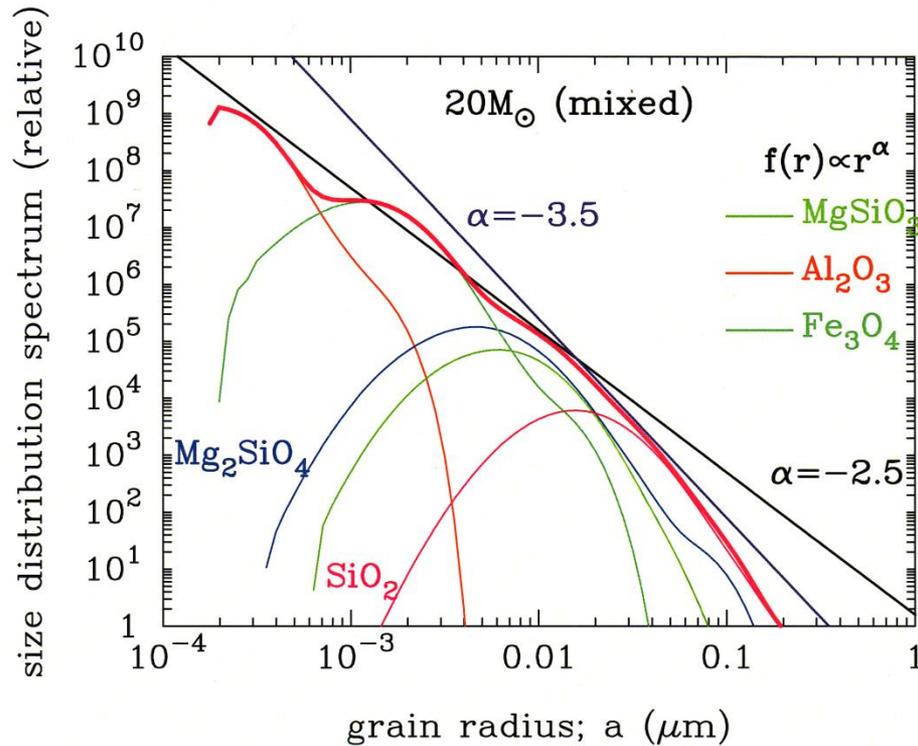
# 2-3-2. Size distribution of newly formed dust



Nozawa+'03, ApJ, 598, 785

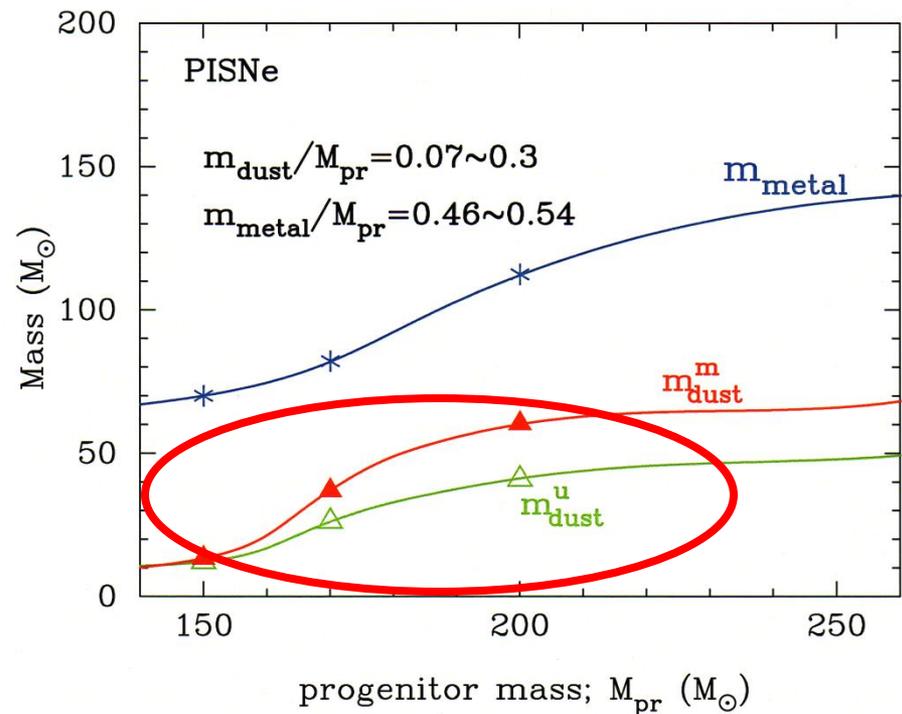
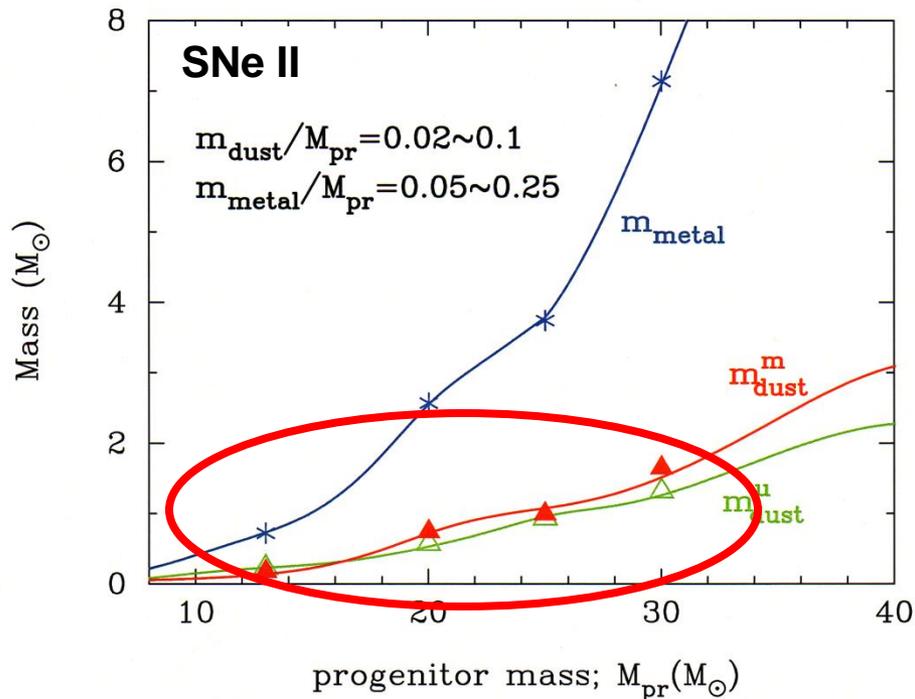
- C, SiO<sub>2</sub>, and Fe grains have lognormal size distribution, while the other grains have power-law size distribution
- The composition and size distribution of dust formed are almost independent of type of supernova  
 (average dust radius is smaller for PISNe than SNe II-P)

## 2-3-3. Size distribution of dust in mixed cases



- Because oxygen is rich in the mixed ejecta, only silicates (MgSiO<sub>3</sub>, Mg<sub>2</sub>SiO<sub>4</sub>, SiO<sub>2</sub>) and oxides (Fe<sub>3</sub>O<sub>4</sub>, Al<sub>2</sub>O<sub>3</sub>) form
- The size distribution of each dust species except for Al<sub>2</sub>O<sub>3</sub> is lognormal-like

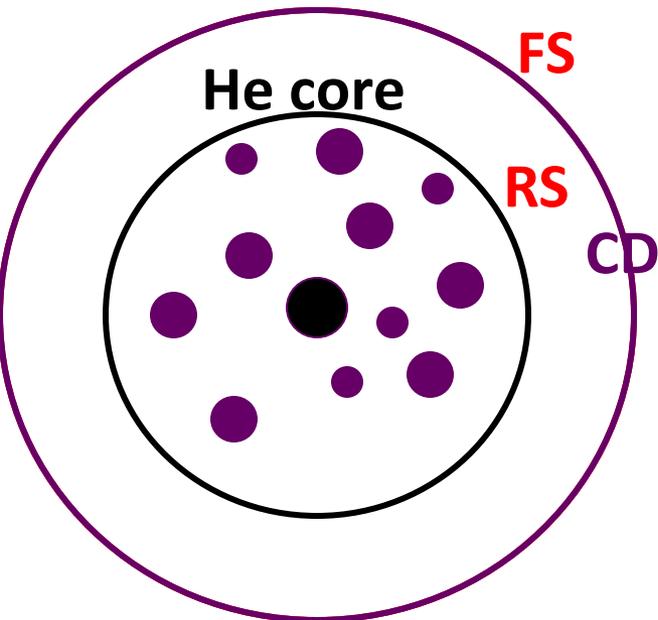
# 2-3-4. Total mass of dust formed



- **Total dust mass increases with increasing MzAMS**  
SNe II:  $M_{\text{dust}} = 0.1\text{-}2 M_{\text{sun}}$ ,  $f_{\text{dep}} = M_{\text{dust}} / M_{\text{metal}} = 0.2\text{-}0.3$   
PISNe:  $M_{\text{dust}} = 10\text{-}60 M_{\text{sun}}$ ,  $f_{\text{dep}} = M_{\text{dust}} / M_{\text{metal}} = 0.3\text{-}0.5$
- **Dust mass for the mixed case is generally larger than for the unmixed case**

# 3. Dust Evolution in SNRs

$$T = (1-2) \times 10^4 \text{ K}$$
$$n_{\text{H},0} = 0.1-1 \text{ cm}^{-3}$$



# 3-1-1. Time evolution of SNRs

- Basic equations (spherical symmetry)

$$\frac{\partial \rho}{\partial t} + \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \rho v) = 0$$

$$\frac{\partial}{\partial t} (\rho v) + \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \rho v^2) = -\frac{\partial P}{\partial r}$$

$$\begin{aligned} \frac{\partial}{\partial t} \left( \frac{\rho v^2}{2} + \frac{P}{\gamma - 1} \right) + \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \left[ \frac{\rho v^2}{2} + \frac{\gamma P}{\gamma - 1} \right] v \right) \\ = -(n_e n_H \Lambda_{\text{gas}}(T) + \Lambda_{\text{ic}}(T) + \Lambda_{\text{d}}(n_H, T)) \end{aligned}$$

$\Lambda_{\text{gas}}(T)$  : cooling function of gas by the atomic process

(Sutherland & Dopita 1993; Smith et al. 2001)

$\Lambda_{\text{ic}}(T)$  : inverse Compton cooling (Ikeuchi & Ostriker 1986)

$$\Lambda_{\text{ic}}(T) = 5.41 \times 10^{-32} (1+z)^4 n_e (T/10^4 \text{ K}) \text{ (we adopt } z = 20)$$

$\Lambda_{\text{d}}(n_H, T)$  : cooling of gas through thermal emission of dust

- numerical code : flux splitting method (van Albada et al. 1982)

## 3-1-2. Initial condition for hydro calculations

### ▪ Hydrodynamical model of SNe (Umeda & Nomoto'02)

- SNe II :  $M_{\text{pr}}=13, 20, 25, 30 M_{\text{sun}}$  ( $E_{51}=1$ )
- PISNe :  $M_{\text{pr}}=170$  ( $E_{51}=20$ ),  $200 M_{\text{sun}}$  ( $E_{51}=28$ )

### ▪ The ambient medium (homogeneous)

- gas temperature :  $T = 10^4 \text{ K}$
- gas density :  $n_{\text{H},0} = 0.1, 1, \text{ and } 10 \text{ cm}^{-3}$

### ▪ Dust Model

- initial size distribution and spatial distribution of dust  
→ results of dust formation calculations
- treating as a test particle

The calculation is performed from 10 yr up to  $\sim 10^6$  yr

# 3-2-1. Dynamics of dust

- deceleration of dust due to drag force (Baines et al. 1965)

$$\frac{dw_d}{dt} = \frac{F_{\text{drag}}}{m_d} = -\frac{3n_H kT}{2a\rho_d} \sum_i A_i G_i(s_i) \quad (w_d : \text{relative velocity})$$

$\rho_d$  ; mass density of a grain

$A_i$  ; the number abundance of gas species  $i$  normalized by  $n_H$

$$G_i(s_i) = \left( s_i^2 + 1 - \frac{1}{4s_i^2} \right) \text{erf}(s_i) + \left( s_i + \frac{1}{2s_i} \right) \frac{e^{-s_i^2}}{\sqrt{\pi}}$$

↓

$$G_i(s_i) \approx \frac{8s_i}{3\sqrt{\pi}} \left( 1 + \frac{9\pi}{64}s_i^2 \right)^{\frac{1}{2}} \quad (\text{Draine \& Salpeter 1979})$$

where  $s_i^2 = m_i w_d^2 / 2kT$

## 3-2-2. Erosion rate of dust by sputtering

- dust destruction by sputtering (e.g., Dwek, Foster & Vancura 1996)

$$\frac{da}{dt} = -\frac{m_{\text{sp}}}{4\pi a^2 \rho_d} \sum_i \mathcal{R}(Y_i(E))$$

$Y_i(E) = 2Y_i^0(E)$  ; the angle-averaged sputtering yield

$m_{\text{sp}}$  ; average mass of the sputtered atoms

- rate equation over a modified Maxwellian distribution of gas taken account of relative velocity of dust to gas (e.g., Shull 1978)

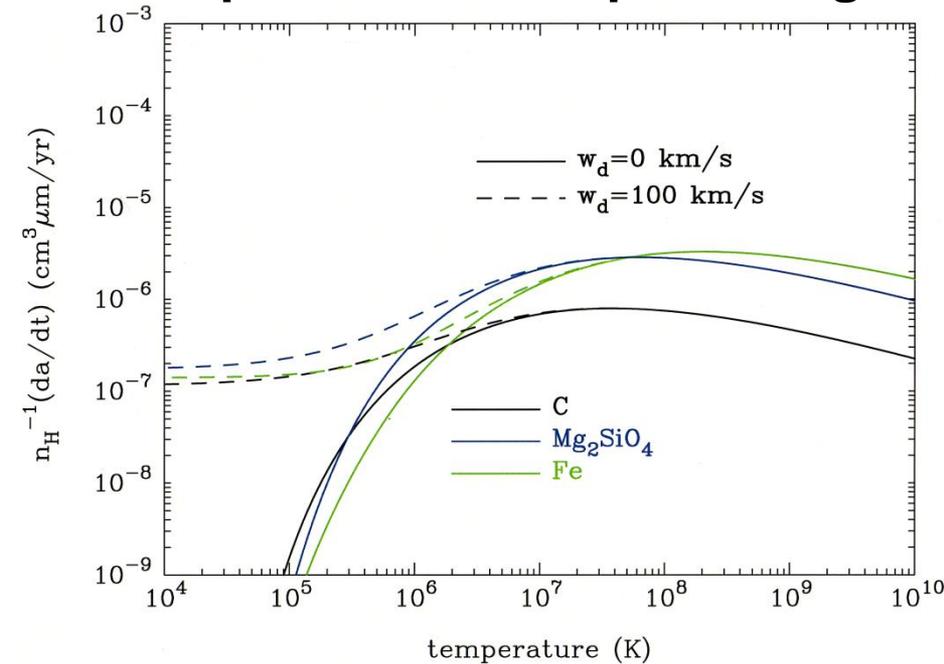
$$\mathcal{R}(X_i(\epsilon)) = n_{\text{H}} A_i \pi a^2 \left( \frac{8kT}{\pi m_i} \right)^{\frac{1}{2}} \frac{e^{-s_i^2}}{2s_i} \int \sqrt{\epsilon} e^{-\epsilon} \sinh(2s_i \sqrt{\epsilon}) X_i(\epsilon) d\epsilon$$

where  $\epsilon = E/kT$

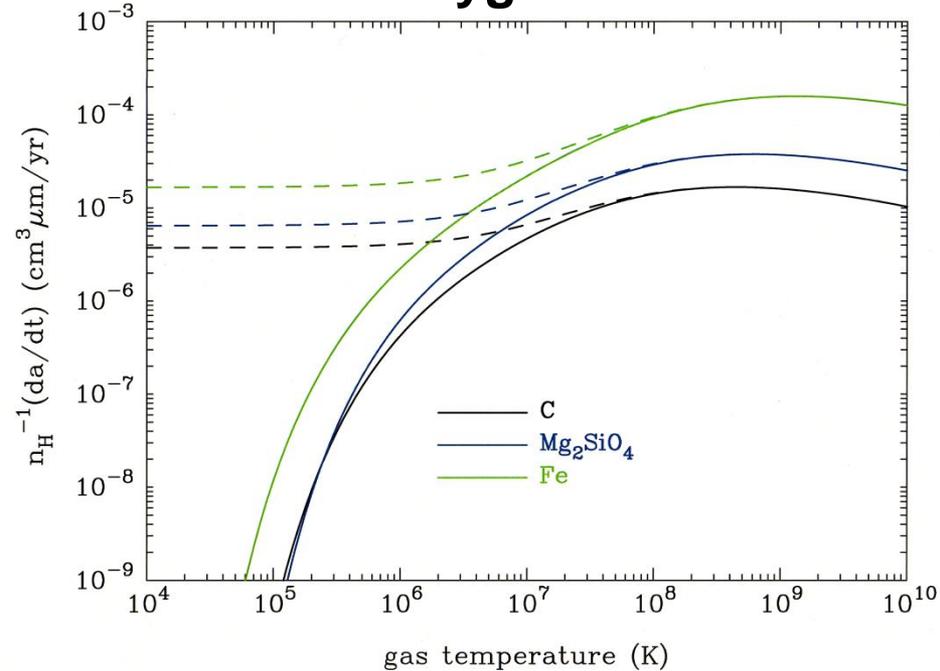
# 3-2-3. Erosion rate of dust by sputtering

Nozawa+'06, ApJ, 648, 435

for primordial composition gas



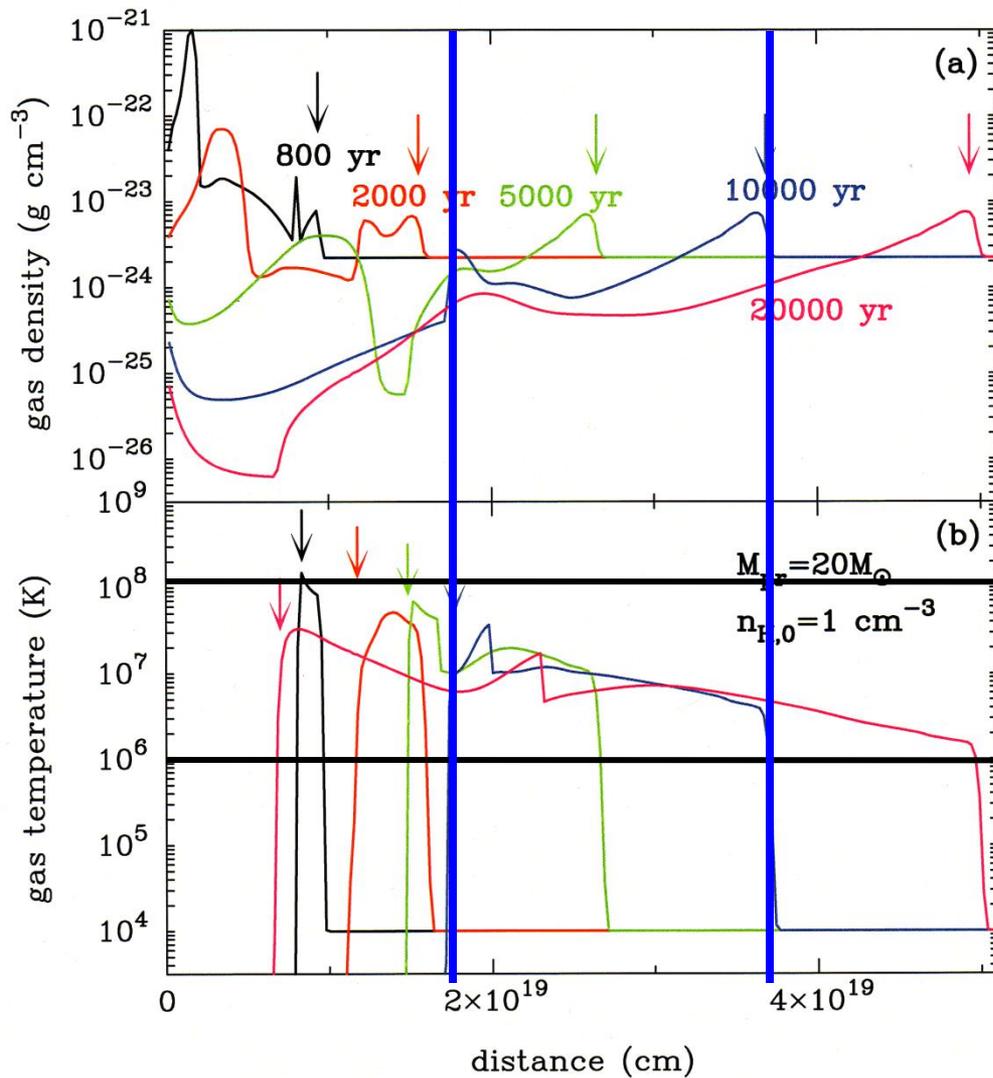
for oxygen ions



- erosion rate by sputtering quickly increases above  $10^5$  K and peaks at  $10^7$  -  $10^8$  K
- erosion rate :  $da / dt \sim 10^{-6} n_H \mu\text{m yr}^{-1} \text{cm}^3$   
for the primordial gas (H and He) at  $T > 10^6$  K

# 3-3-1. Temperature and density of gas in SNRs

Nozawa+'07, ApJ, 666, 955



Model :  $M_{\text{pr}} = 20 M_{\text{sun}} (E_{51}=1)$   
 $n_{\text{H},0} = 1 \text{ cm}^{-3}$

Downward-pointing arrows:  
forward shock in upper panel  
reverse shock in lower panel

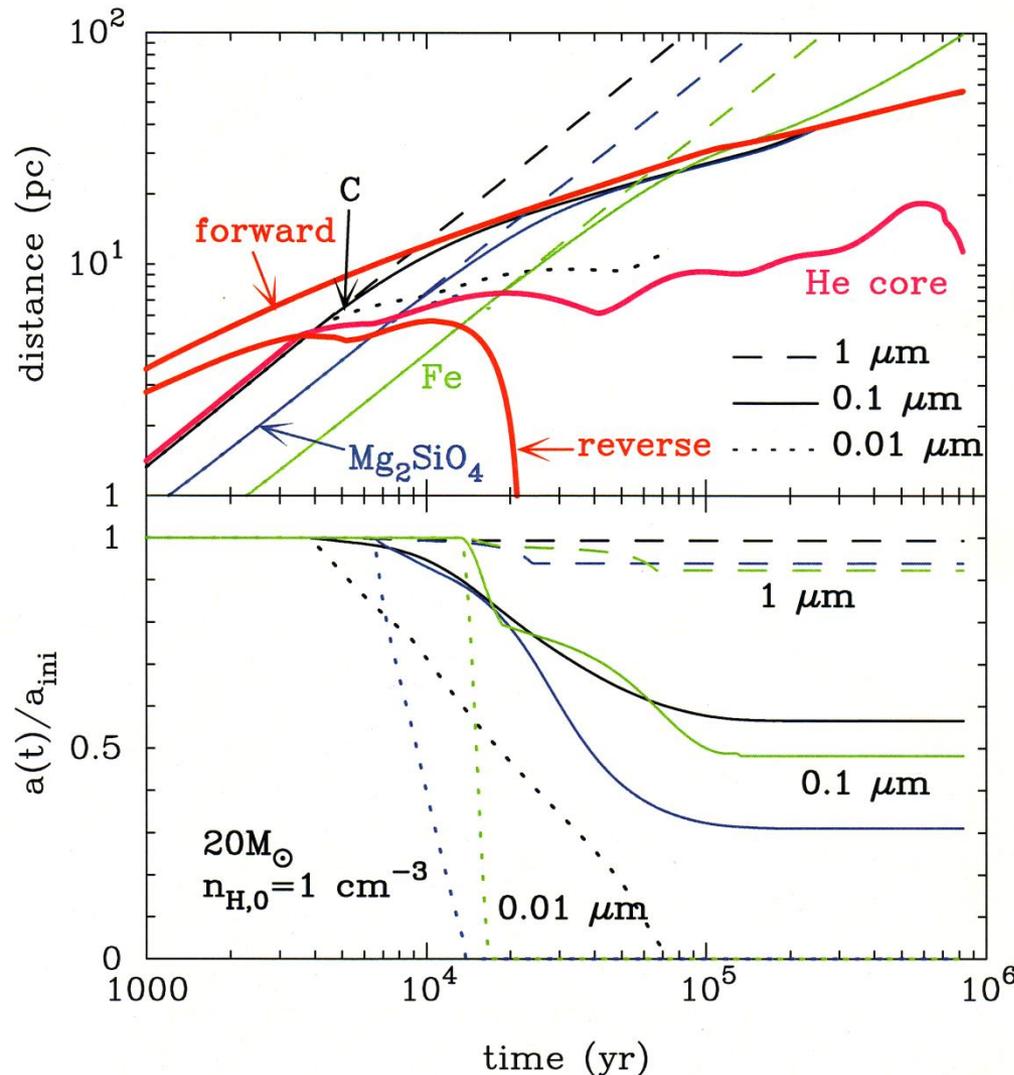
The temperature of the gas swept up by the shocks

→  $10^6$ - $10^8$  K

↓  
Dust grains residing in the shocked hot gas are eroded by sputtering

# 3-3-2. Evolution of dust in SNRs

Nozawa+'07, ApJ, 666, 955



Model :  $M_{pr} = 20 M_{\text{sun}} (E_{51}=1)$   
 $n_{H,0} = 1 \text{ cm}^{-3}$

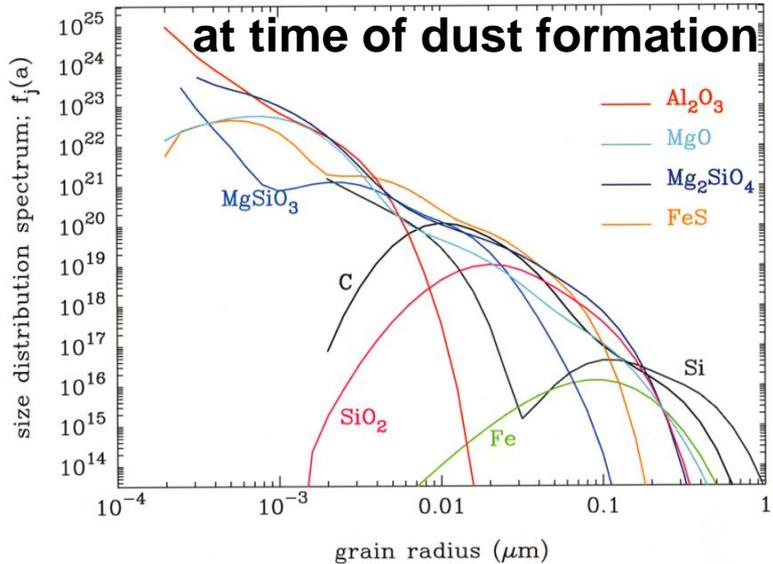
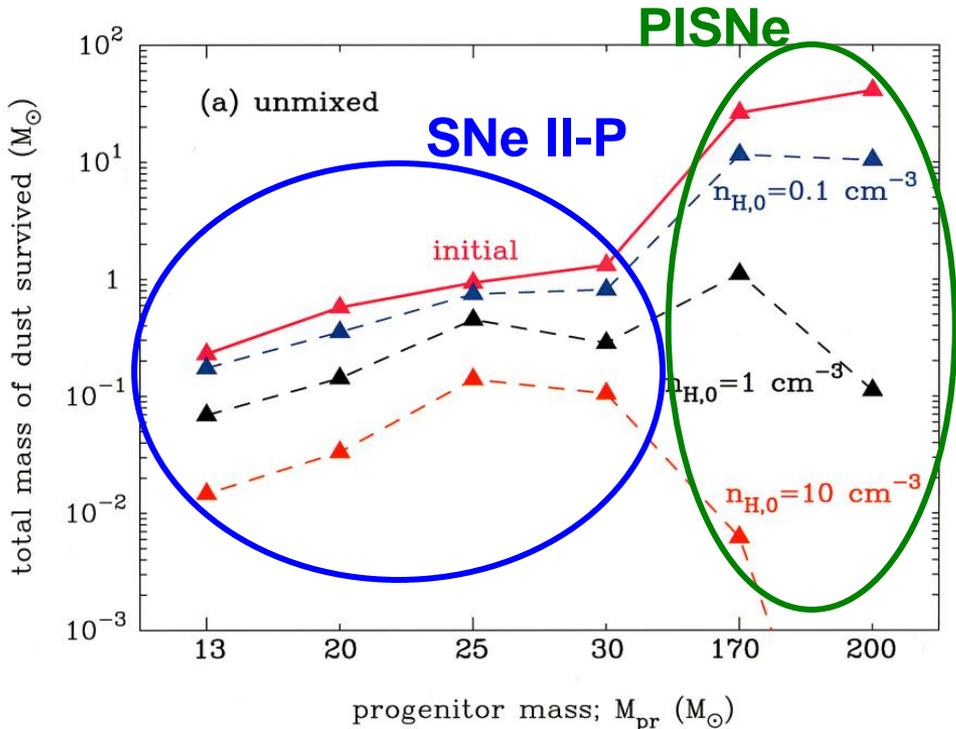
Dust grains in the He core collide with reverse shock at  $(3-13) \times 10^3 \text{ yr}$

**The evolution of dust heavily depends on the initial radius and composition**

- $a_{ini} = 0.01 \mu\text{m}$  (dotted lines)  
→ completely destroyed
- $a_{ini} = 0.1 \mu\text{m}$  (solid lines)  
→ trapped in the shell
- $a_{ini} = 1 \mu\text{m}$  (dashed lines)  
→ injected into the ISM

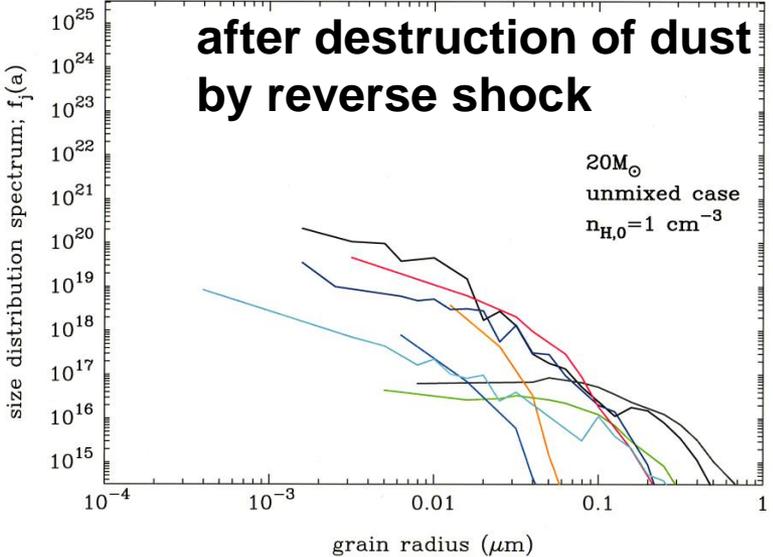
# 3-3-3. Dust mass and size ejected from SN II-P

Nozawa+'07, ApJ, 666, 955



**total dust mass surviving the destruction in Type II-P SNRs; 0.07-0.8 M\_sun ( $n_{H,0} = 0.1-1 \text{ cm}^{-3}$ )**

**size distribution of dust after RS destruction is dominated by large grains ( $> 0.01 \text{ micrometers}$ )**



## 3-4. Summary of dust production in Pop III SNe

- Various species of dust form in the unmixed ejecta  
(almost all Fe, Mg, and Si atoms are locked in dust)
- The fate of newly formed dust within primordial SNRs strongly depends on the initial radii and compositions.
- The size distribution of dust surviving the destruction in SNRs is weighted to relatively large size ( $> 0.01 \mu\text{m}$ ).
- The total mass of dust injected into the ISM decreases with increasing the ambient gas density

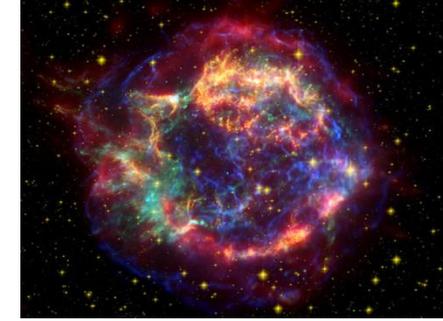
for  $n_{\text{H},0} = 0.1\text{-}1 \text{ cm}^{-3}$

SNe II-P  $\rightarrow M_{\text{dust}} = 0.1\text{-}0.8 M_{\text{sun}}$

PISNe  $\rightarrow M_{\text{dust}} = 0.1\text{-}15 M_{\text{sun}}$

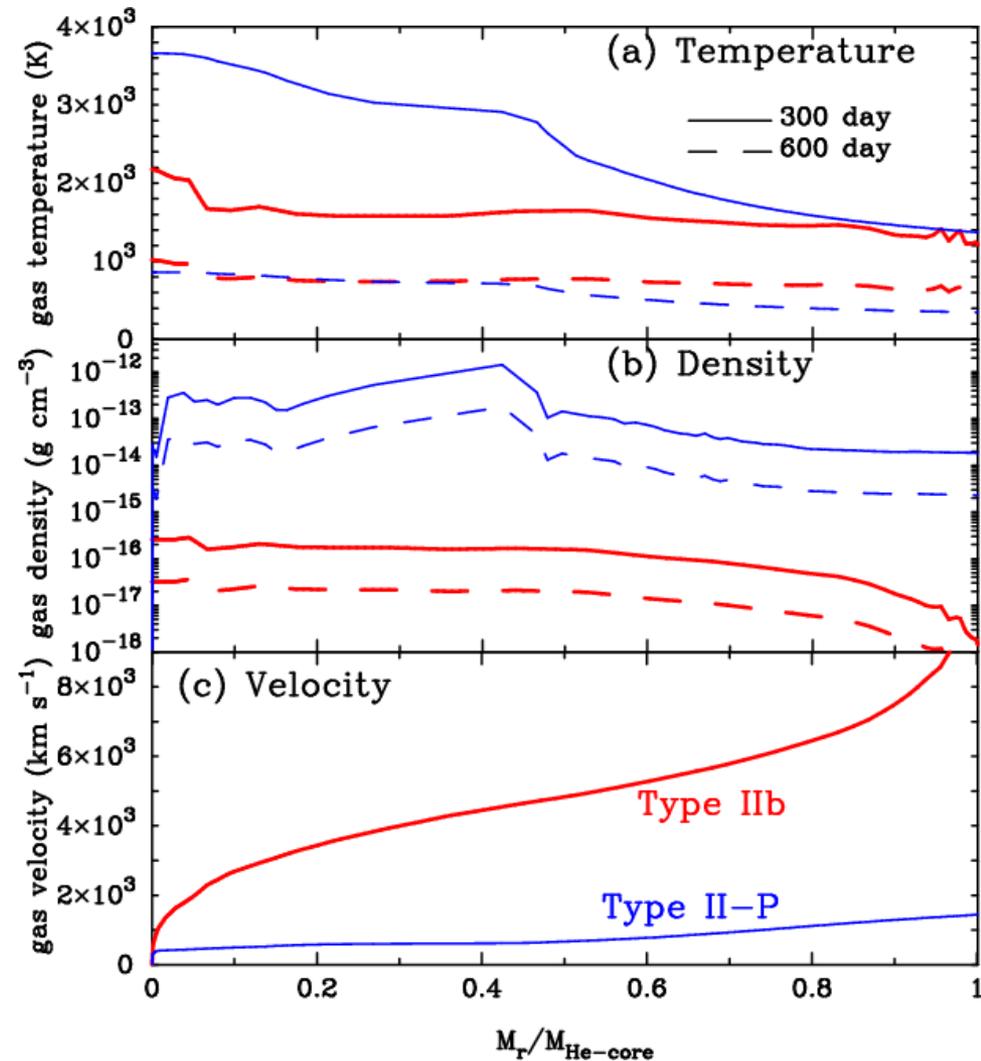
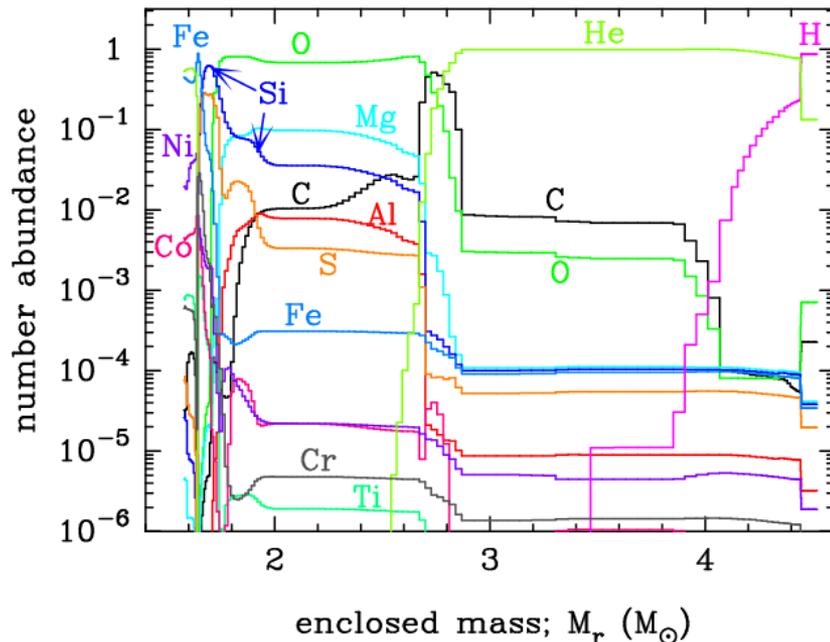
### **3. Formation of dust grains in various types of SNe**

# 5-1. Dust formation in Type IIb SN

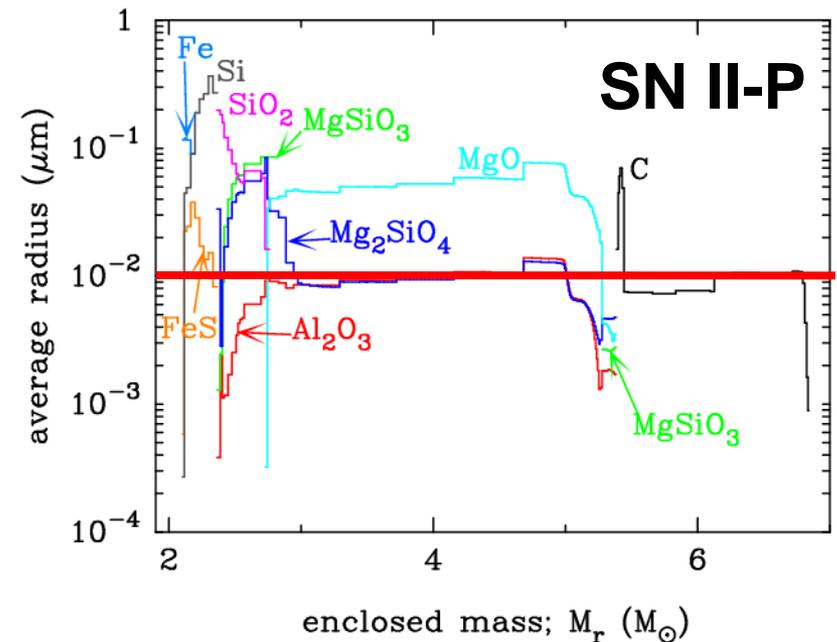
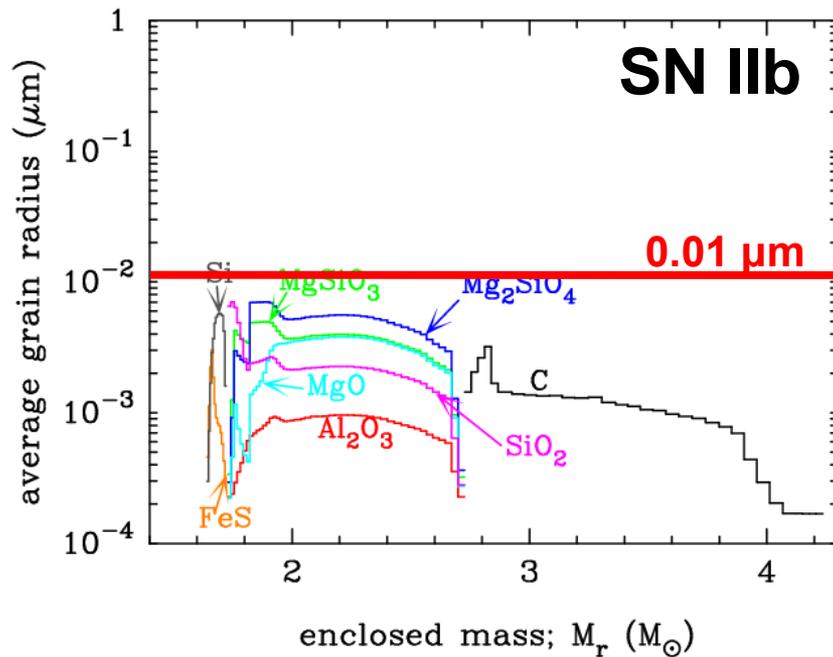


## ○ SN IIb model (SN1993J-like model)

- $M_{\text{eje}} = 2.94 M_{\text{sun}}$   
 $M_{\text{ZAMS}} = 18 M_{\text{sun}}$   
 $M_{\text{H-env}} = 0.08 M_{\text{sun}}$
- $E_{51} = 1$
- $M(^{56}\text{Ni}) = 0.07 M_{\text{sun}}$



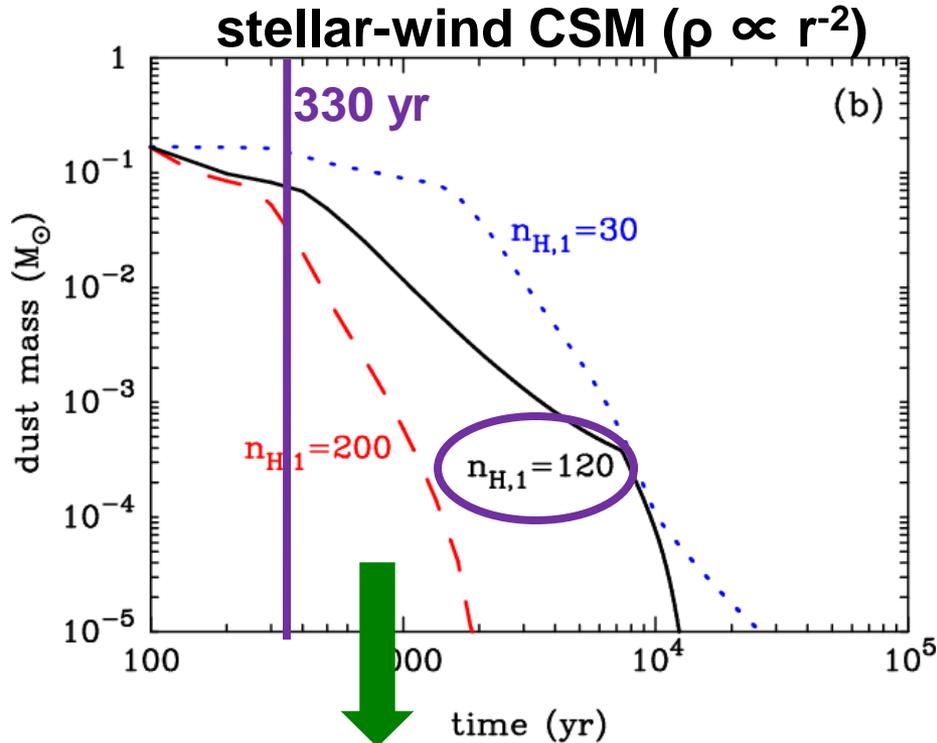
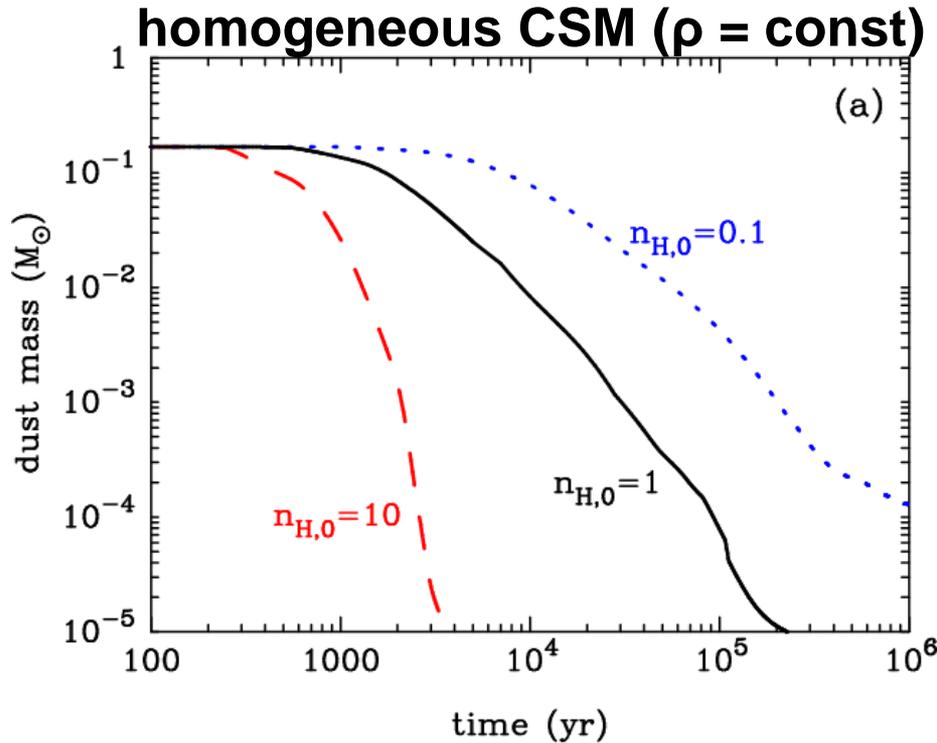
# 5-2. Dependence of dust radii on SN type



- condensation time of dust **300-700 d** after explosion
- total mass of dust formed
  - **$0.167 M_{\text{sun}}$**  in SN IIb
  - **$0.1-1 M_{\text{sun}}$**  in SN II-P

- the radius of dust formed in H-stripped SNe is small
  - **SN IIb without massive H-env**  $\rightarrow a_{\text{dust}} < 0.01 \mu\text{m}$
  - **SN II-P with massive H-env**  $\rightarrow a_{\text{dust}} > 0.01 \mu\text{m}$

# 5-3. Destruction of dust in Type IIb SNR

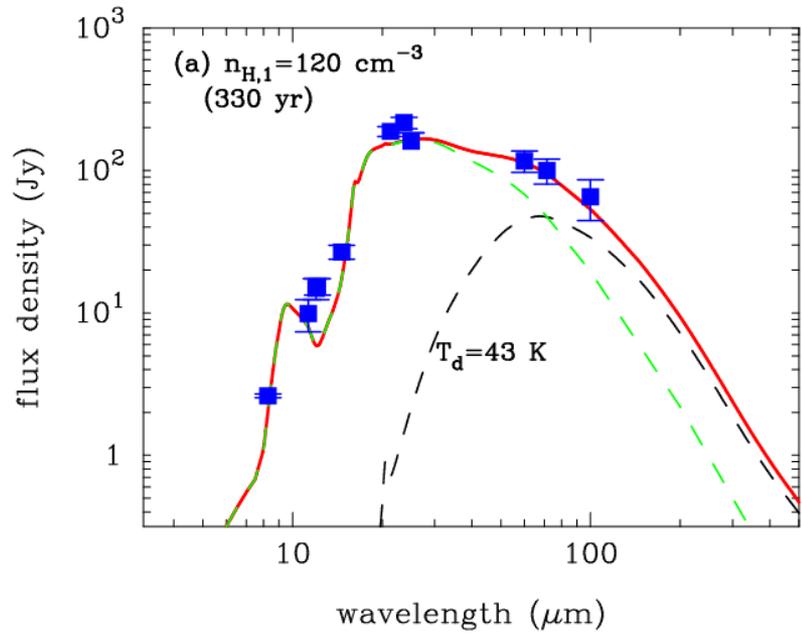


$n_{H,1} = 30, 120, 200 /cc \rightarrow dM/dt = 2.0, 8.0, 13 \times 10^{-5} M_{sun}/yr$  for  $v_w = 10$  km/s

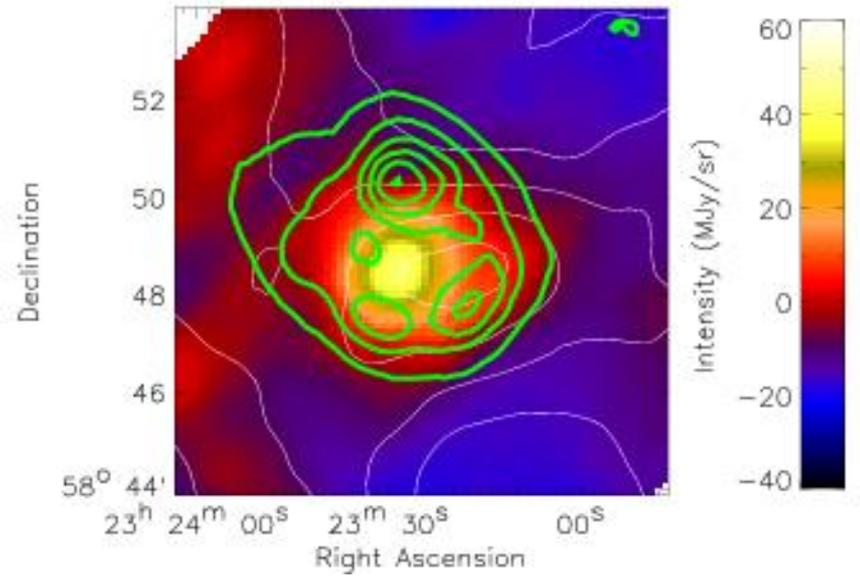
Almost all newly formed grains are destroyed in shocked gas within the SNR for CSM gas density of  $n_H > 0.1 /cc$

- small radius of newly formed dust
- early arrival of reverse shock at dust-forming region

# 5-4. IR emission from dust in Cas A SNR



AKARI corrected 90  $\mu\text{m}$  image



- total mass of dust formed  
 $M_{\text{dust}} = 0.167 M_{\text{sun}}$
- shocked dust : **0.095  $M_{\text{sun}}$**   
 $M_{\text{d,warm}} = 0.008 M_{\text{sun}}$
- unshocked dust :  
 $M_{\text{d,cool}} = 0.072 M_{\text{sun}}$   
with  $T_{\text{dust}} \sim 40 \text{ K}$

- AKARI observation**  
 $M_{\text{d,cool}} = 0.03\text{-}0.06 M_{\text{sun}}$   
 $T_{\text{dust}} = 33\text{-}41 \text{ K}$   
 (Sibthorpe+'10)

- Herschel observation**  
 $M_{\text{d,cool}} = 0.075 M_{\text{sun}}$   
 $T_{\text{dust}} \sim 35 \text{ K}$  (Barlow+'10)

Nozawa et al. 2010, ApJ, 713, 356

# 6. Summary of this talk

- SNe are important sources of dust?
  - maybe, Yes in the early universe  
(at least, to serve the seeds for grain growth in the ISM)
- Size of newly formed dust depends on types of SNe
  - H-retaining SNe (Type II-P) :  $a_{ave} > 0.01 \mu\text{m}$
  - H-stripped SNe (Type IIb/IIc and Ia) :  $a_{ave} < 0.01 \mu\text{m}$ 
    - dust is almost completely destroyed in the SNRs
    - H-stripped SNe may be poor producers of dust
- Our model treating dust formation and evolution self-consistently can reproduce IR emission from Cas A
- Mass of dust in SNe must be dominated by cool dust
  - FIR and submm observations of SNe are essential