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超新星爆発時における ダストの形成・放出過程

<u>野沢貴也(Takaya Nozawa)</u>

(National Astronomical Observatory of Japan)

Main collaborators:

T. Kozasa, A. Habe (Hokkaido University)

H. Umeda (University of Tokyo)

K. Maeda (Kyoto University), K. Nomoto (Kavli IPMU)

<u>N. Tominaga</u> (Konan University)

0-1. Introduction

• SNe are important sources of interstellar dust?

- abundant metal (metal : N > 5)
- low temperature (T < ~2000 K)
- high density (n > $\sim 10^6$ cm⁻³)

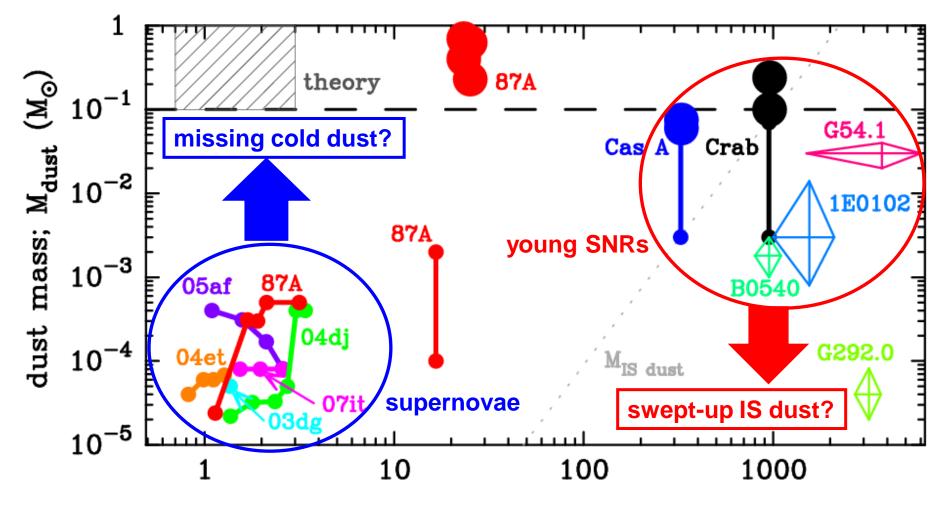
- mass-loss winds of AGB stars expanding ejecta of supernovae
- huge amounts of dust grains (>10⁸ Msun) are detected in host galaxies of quasars at redshift z > 5
 - → 0.1 Msun of dust per SN is needed to explain such massive dust at high-z (e.g. Dwek et al. 2007)
- contribution of dust mass from AGB stars and SNe

n(AGB stars) / n(SNe) ~ 10-20

Mdust = 0.01-0.05 Msun per AGB (Zhukovska & Gail 2008) Mdust = 0.1-1.0 Msun per SN (e.g., Nozawa et al. 2003, 2007)

1. Observations of Dust Formation in SNe (and SNRs)

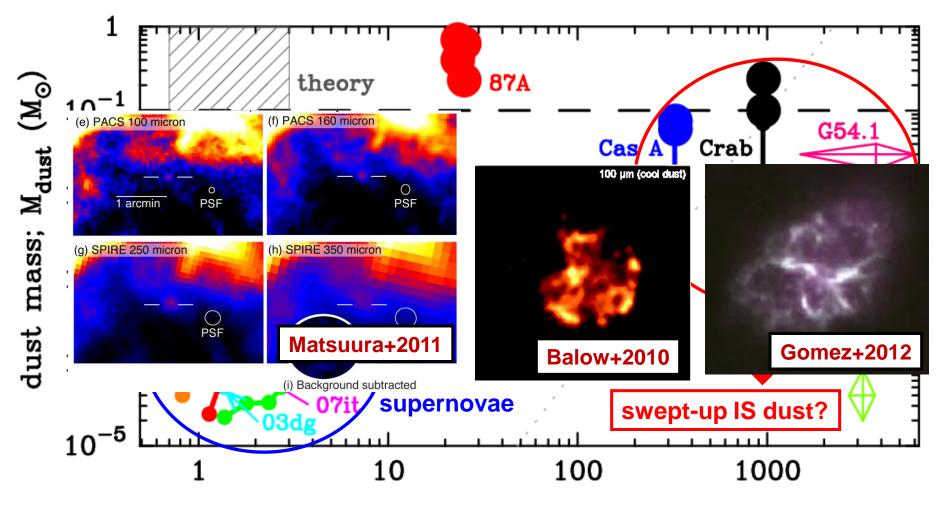
1-1. Summary of observed dust mass in CCSNe



time after the explosion (yr)

Far-IR to sub-mm observations are essential for revealing the mass of dust grains produced in the ejecta of SNe

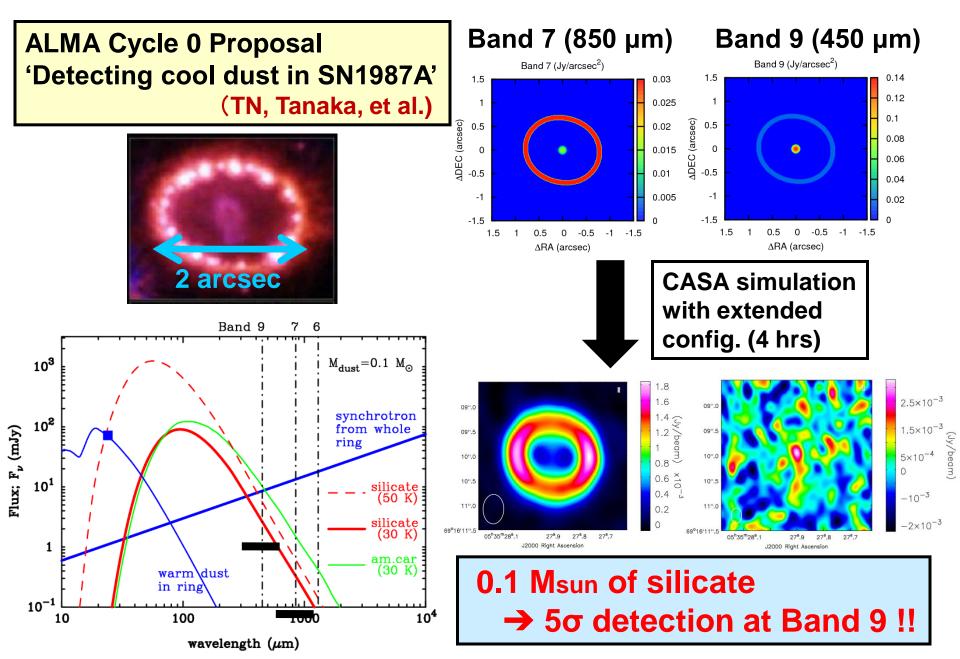
1-1. Summary of observed dust mass in CCSNe



time after the explosion (yr)

Far-IR to sub-mm observations are essential for revealing the mass of dust grains produced in the ejecta of SNe

1-2. Resolving cool dust in SN 87A with ALMA



1-3. Successful ALMA proposals for SN 1987A

2011.0.00221.5

PI	Ехес	Country	Institute
Nozawa, Takaya	EA	Japan	The University of Tokyo
COI			
Tanaka, Masaomi	EA.	Japan	The University of Tokyo
Moriya, Takashi	EA.	Japan	University of Tokyo
Minamidani. Tetsuhiro	EA	Japan	Hokkaido University
Kozasa, Takashi	EA	Japan	Hokkaido University

This proposal was ranked in the highest priority !!

PI	Exec	Country	Institute
ndebetouw. Remv	NA	United States	Virginia, University of
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AcCray, Richard	NA	United States	Colorado at Boulder. Univ of
Matsuura, Mikako	EU	United Kingdom	London. University of
Andjelic, Nilica	OTHER	Serbia	Belgrade, University of
Arbutina, Bojan	OTHER	Serbia	Belgrade, University of
Baes, Maarten	EU	Belgium	Ghent University
Bolatto, Alberto	NA	United States	Maryland, University of
Burrows, David	NA	United States	Pennsylvania State University
Chevelier, Roger	NA.	United States	Virginia, University of
Gaensler, Bryan	OTHER	Australia	Sydney, University of
Long, Knoz	NA	United States	Space Telescope Science Institute
Lundqvist. Peter	EU	Sweden	Stockholm University
Metxner, Margaret	NA	United States	Space Telescope Science Institute
Marcaide, Jon	EU	Spain	Valencia, University of
Marti-Vidal, Ivan	EU	Germany	Max-Planck-Institute for Radio Astronomy
DTSUKA, Masaaki	EA/NA	Taiwan	Academia Sinica
Sandstrom, Karin	EU	Germany	Max-Planck-Institute for Astronomy
Sonneborn, George	NA	United States	National Aeronautics and Space Administration
Staveley-Smith, Lister	OTHER	Australia	International Centre for Radio Astronomy Research
van Leon, jacco	EU	United Kingdom	Keele University
Urosevic, Dejan	OTHER	Serbia	Belgrade, University of
Vlahakis, Catherine	a	Chile	Chile, University of
Zekovic, Vladimir	OTHER	Serbia	Belgrade, University of
Zanardo, Giovanna	OTHER	Australia	International Centre for Radio Astronomy Research
Ng. Chi-Yung	NA	Canada	McGill University
Park. Sangwook	NA	United States	Texas at Arlington. University of
Barlow, Michael	EU	United Kingdom	Landon, University of
Clayton, Geoffrey	NA	United States	Louisiana State University
Wesson, Roger	EU	United Kingdom	London, University of
Dwek, ELi	NA	United States	National Aeronautics and Space Administration
Bouchet, Patrice	EU	France	CEA Gaclay
Lakicevic, Masa	EU	Germany	European Southern Observatory
Potter, Toby	OTHER	Australia	International Centre for Radio Astronomy Research

1-3. Successful ALMA proposals for SN 1987A

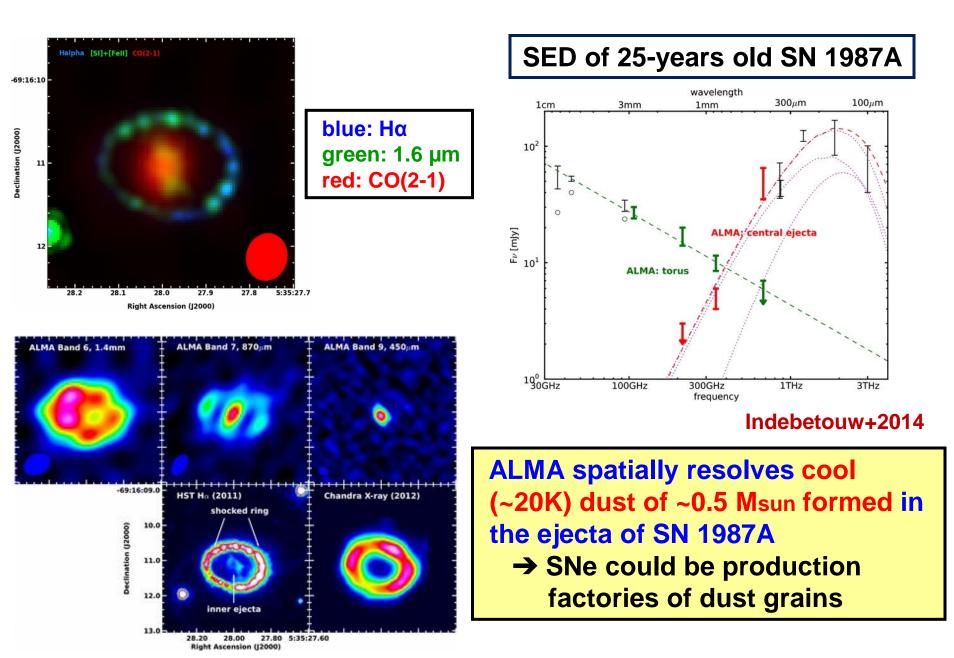
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inaka, Masa, ji	EA.	Japan	The University of Tokyo	McCray, Richard	NA	United States	Colorado at Boulder. Univ of
oriya, Takashi	EA	Japan	University of Tekyo	Matsuura, Mikako	EU	United Kingdom	London. University of
inamidani. Tetsuhin	EA	Japan	Hokksido University	Andjelic, Nilica	OTHER	Serbia	Belgrade, University of
izasa, Takashi	EA	Japan	Hokkaido University	Arbutina, Bojan	OTHER	Serbia	Belgrade, University of
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Potter, Toby

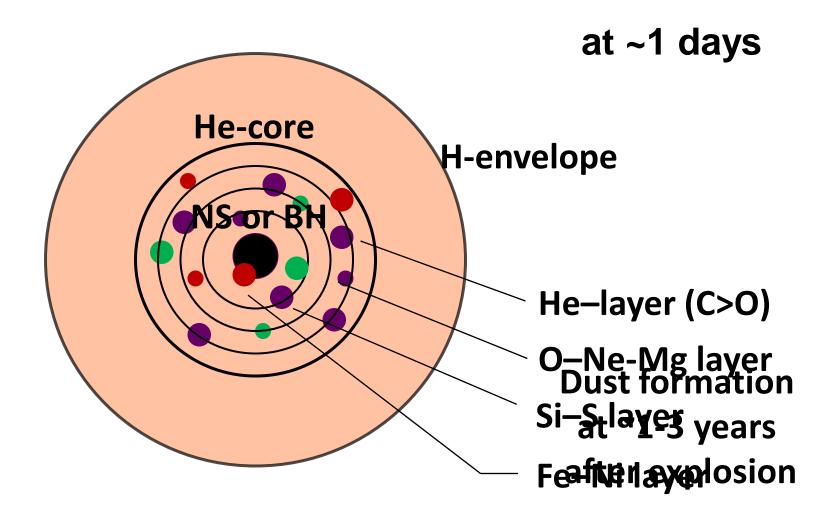
OTHER Australia

International Centre for Radio Astronomy Research

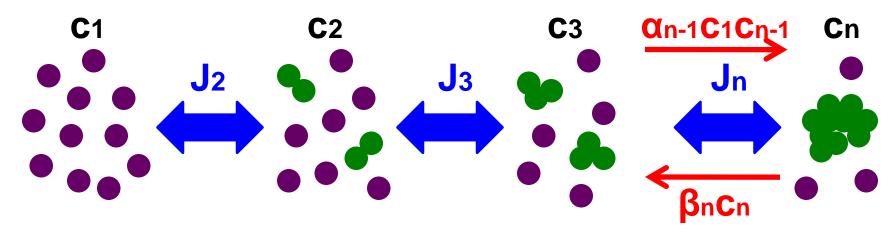
1-4. ALMA reveals dust formed in SN 1987A



2. Dust Formation in the ejecta of SNe



2-1. Formulation of dust formation



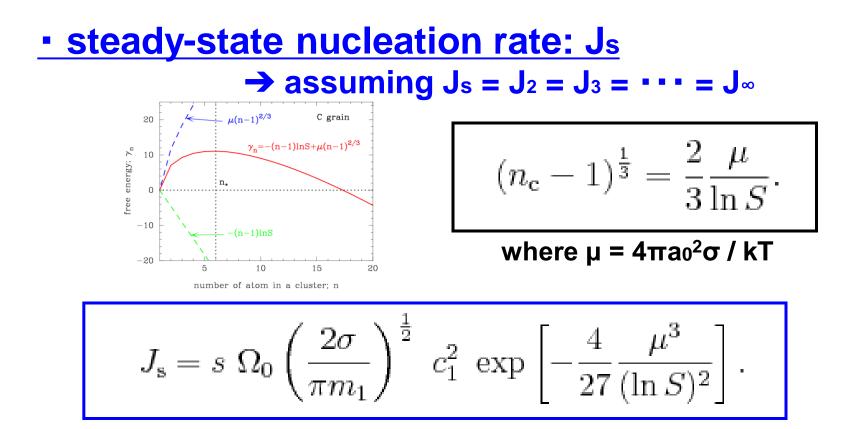
master equations

$$\frac{dc_n}{dt} = J_n(t) - J_{n+1}(t) \quad \text{for } 2 \le n \le n_*,$$

$$J_n(t) = \alpha_{n-1}c_{n-1}c_1 - \beta_n c_n \text{ for } 2 \le n \le n_*,$$

$$\alpha_n = \frac{s_n}{1+\delta_{1n}} \ 4\pi a_0^2 \ n^{\frac{2}{3}} \left(\frac{kT}{2\pi m_n}\right)^{\frac{1}{2}}, \qquad \beta_n = \alpha_{n-1} \frac{\mathring{c}_{n-1}}{\mathring{c}_n} \mathring{c}_1,$$

2-2. Non-steady-state nucleation



• non-steady-state dust formation

$$\frac{dc_n}{dt} = J_n(t) - J_{n+1}(t) \quad \text{for } 2 \le n \le n_*,$$

2-3. Basic equations for dust formation

Equation of mass conservation

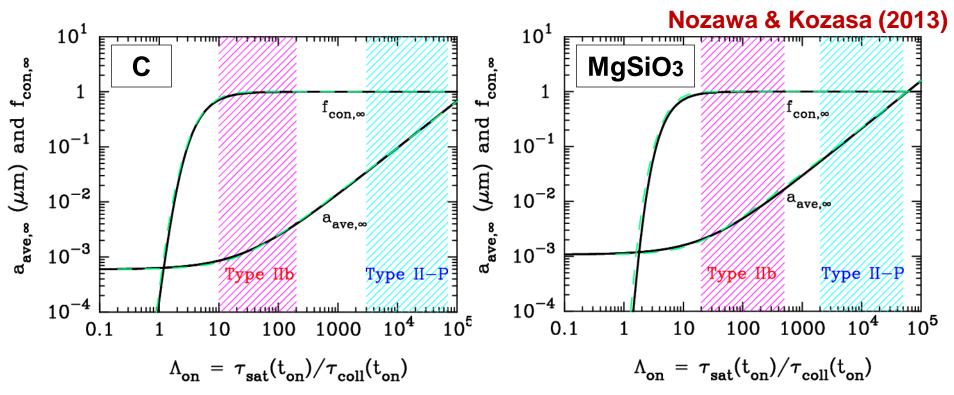
$$c_{10} - c_1 = \int_{t_0}^t J_{n_*}(t') \frac{a^3(t,t')}{a_0^3} dt',$$

Equation of grain growth

$$\frac{da}{dt} = s\Omega_0 \left(\frac{kT}{2\pi m_1}\right)^{\frac{1}{2}} c_1 \left(1 - \frac{1}{S}\right),$$
$$\frac{dV}{dt} = s\Omega_0 \ 4\pi a^2 \left(\frac{kT}{2\pi m_1}\right)^{\frac{1}{2}} c_1 \left(1 - \frac{1}{S}\right),$$

Growth rate is independent of grain radius

2-4. Scaling relation of average grain radius



<u> Λ on = Tsat/Tcoll</u>: ratio of supersaturation timescale to gas collision timescale at the onset time (ton) of dust formation <u> Λ on = Tsat/Tcoll \propto Tcool Ngas</u>

fcon,∞ and aave,∞ are uniquely determined by Λon

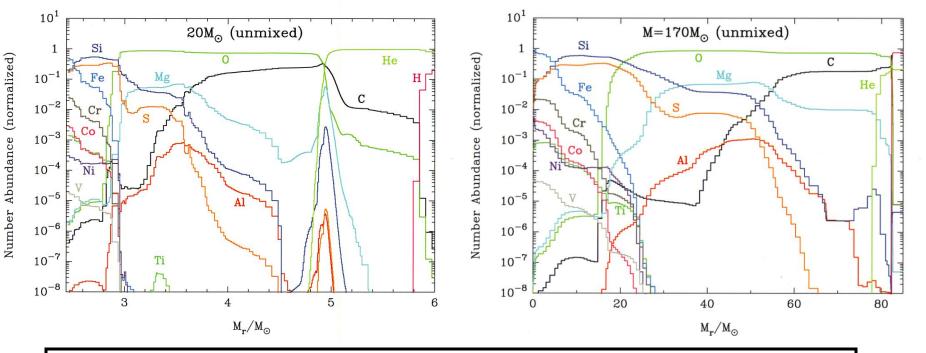
steady-state nucleation rate is applicable for Λon > 30

2-5. Dust formation in primordial SNe

Nozawa+2003, ApJ, 598, 785

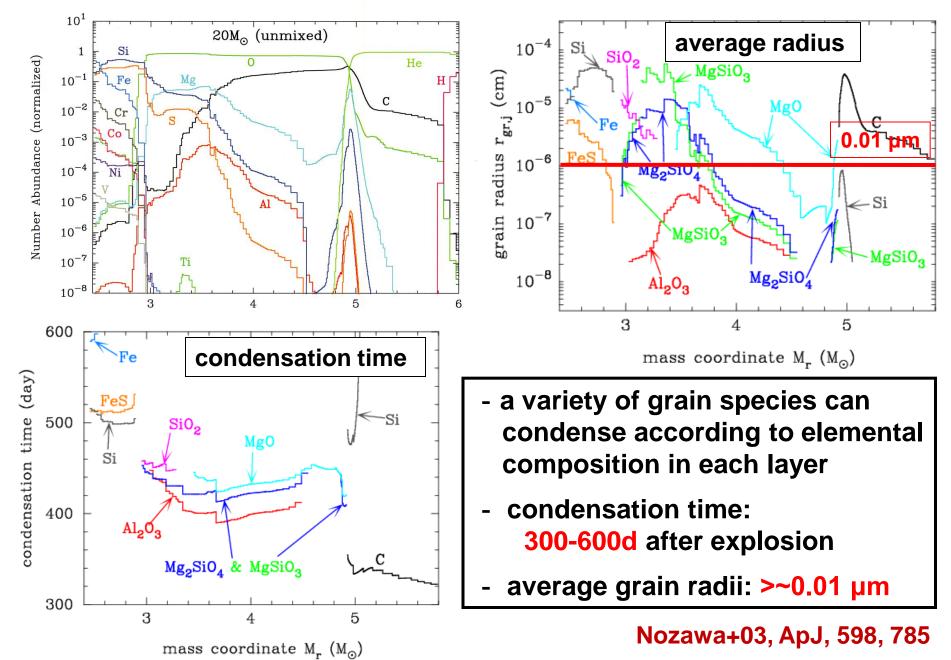
O Population III SNe model (Umeda & Nomoto 2002)

- SNe II : MZAMS = 13, 20, 25, 30 Msun (E_{51} =1)
- **PISNe** : MZAMS = 170 Msun (E_{51} =20), 200 Msun (E_{51} =28)

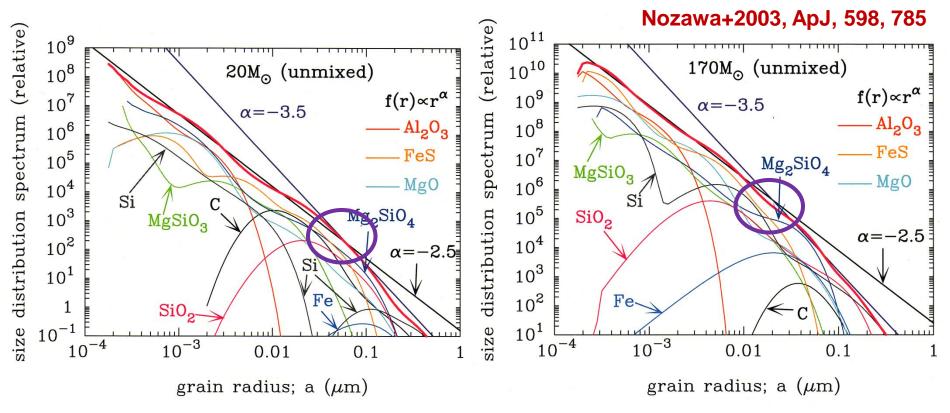


- nucleation and grain growth theory (Kozasa & Hasegawa 1987)
- no mixing of elements within the He-core
- complete formation of CO and SiO

2-6. Dust formed in Type II-P SNe



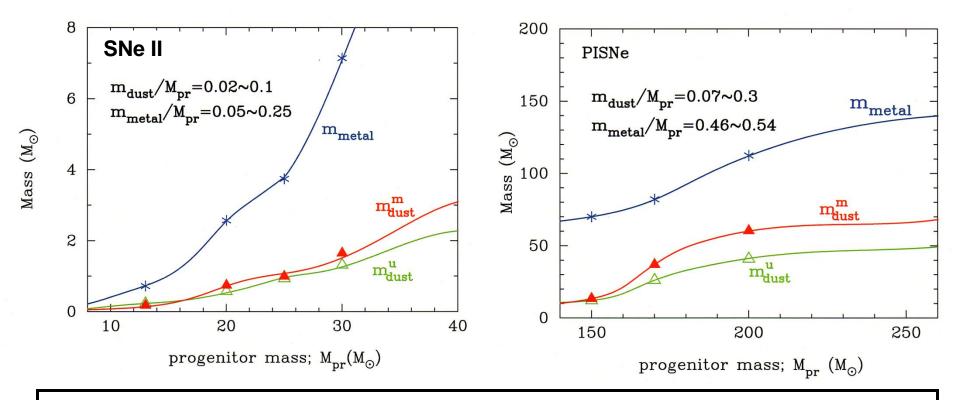
2-7. Size distribution of newly formed dust



- C, SiO2, and Fe grains have lognormal-like size distribution, while the other grains have power-law size distribution
- The composition and size distribution of dust formed are almost independent of types of supernova

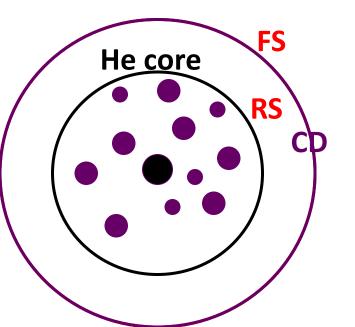
average grain radius is smaller for PISNe than SNe II-P

2-8. Total mass of dust formed in the ejecta



- Total mass of dust is higher for a higher progenitor mass (MZAMS) SNe II : mdust = 0.1-1.5 Msun, mdust / mmetal = 0.2-0.3 PISNe : mdust = 10-30 Msun, mdust / mmetal = 0.3-0.4
- almost all Fe, Mg, and Si are locked up in dust grains, while most of C and O remain in the gas-phase (such as CO)
 → dust-to-metal mass ratio is not high for SNe II

3. Evolution of dust in SN remnants



3-1. Time evolution of SNRs

• Basic equations (spherical symmetry) $\frac{\partial \rho}{\partial t} + \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \rho v) = 0$ $\frac{\partial}{\partial t}(\rho v) + \frac{1}{r^2}\frac{\partial}{\partial r}(r^2\rho v^2) = -\frac{\partial P}{\partial r}$ $\frac{\partial}{\partial t} \left(\frac{\rho v^2}{2} + \frac{P}{\gamma - 1} \right) + \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \left[\frac{\rho v^2}{2} + \frac{\gamma P}{\gamma - 1} \right] v \right)$ $= -(n_{\rm e}n_{\rm H}\Lambda_{\rm gas}(T) + \Lambda_{\rm ic}(T) + \Lambda_{\rm d}(n_{\rm H},T))$ $\Lambda_{gas}(T)$: cooling function of gas by the atomic process (Sutherland & Dopita 1993; Smith et al. 2001) $\Lambda_{\rm ic}(T)$: inverse Compton cooling (Ikeuchi & Ostriker 1986) $\Lambda_{\rm ic}(T) = 5.41 \times 10^{-32} (1+z)^4 n_e(T/10^4 \,{\rm K}) \text{ (we adopt } z = 20)$ $\Lambda_{\rm d}(n_{\rm H},T)$: cooling of gas through thermal emission of dust \cdot numerical code : flux splitting method (van Albada et al. 1982)

3-2. Dynamics of dust

• deceleration of dust due to drag force (Baines et al. 1965)

 $\frac{dw_{\rm d}}{dt} = \frac{F_{\rm drag}}{m_{\rm d}} = -\frac{3n_{\rm H}kT}{2a\rho_{\rm d}}\sum_{i}A_{i}G_{i}(s_{i}) \quad (w_{\rm d}: \text{relative velocity})$ $F_{\text{drag}} = m_{\text{d}} \frac{dw_{\text{d}}}{dt} = -\pi a^{2} \sum n_{i} \langle v_{i} m_{i} v_{i} \cos \theta \rangle$ $\frac{dw_{\text{d}}}{dt} = -\frac{\pi a^{2}}{\frac{4}{3}\pi a^{3}\rho_{d}} n_{\text{H}} \sum A_{i} \langle v_{i} m_{i} v_{i} \cos \theta \rangle$ $\frac{3n_{\text{H}}}{3n_{\text{H}}} - \frac{\pi a^{2}}{2} \sum n_{i} \langle v_{i} m_{i} v_{i} \cos \theta \rangle$ Dy $n_{\rm H}$ $= -\frac{3n_{\rm H}}{4a\rho_d}kT\sum A_iG_i$ $G_i(s_i) \approx \frac{8s_i}{3\sqrt{\pi}} \left(1 + \frac{9\pi}{64}s_i^2\right)^{\frac{1}{2}}$ (Draine & Salpeter 1979) where $s_i^2 = m_i w_d^2 / 2kT$

3-3. Erosion rate of dust by sputtering

• dust destruction by sputtering (e.g., Dwek, Foster & Vancura 1996)

$$\frac{da}{dt} = -\frac{m_{\rm sp}}{4\pi a^2 \rho_{\rm d}} \sum_i \mathcal{R}(Y_i(E))$$

 $Y_i(E) = 2Y_i^0(E)$; the angle-averaged sputtering yield

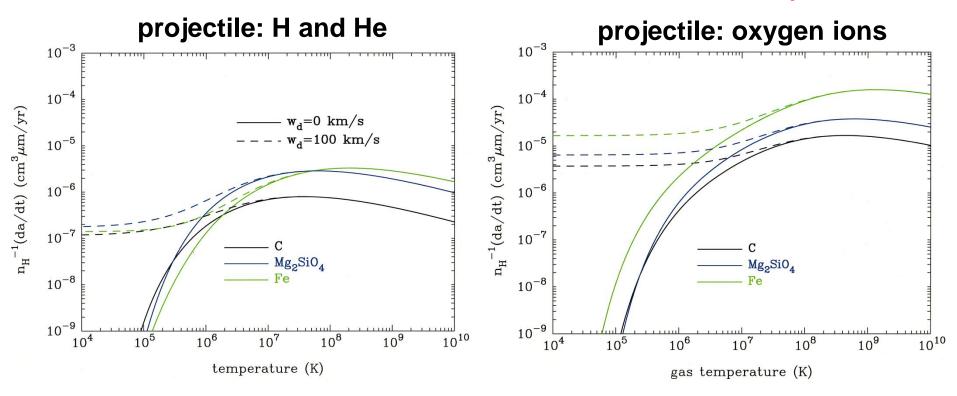
$$\frac{n}{dt} = 4\pi a^2 \frac{da}{dt} = -\pi a^2 \frac{4\pi a_0^3}{3} \sum n_i \langle v_i Y_i \rangle$$
rate
$$\frac{da}{dt} = -\frac{1}{4} \Omega n_{\rm H} \sum A_i \langle v_i Y_i \rangle$$
is taken
)

$$\frac{\mathcal{R}(\cdot)}{W} \qquad \frac{da}{dt} = -\frac{m_{\rm sp}}{2\rho_{\rm d}} n_{\rm H} \sum A_i \left(\frac{8kT}{\pi m_i}\right)^{1/2} \frac{e^{-s_i^2}}{2s_i} \\
\times \int \epsilon^{\frac{1}{2}} e^{-\epsilon_i} \sinh(2s_i\epsilon_i^{\frac{1}{2}}) Y_i^0(\epsilon_i) d\epsilon_i$$

 $X_i(\epsilon)d\epsilon$

3-4. Erosion rate of dust by sputtering

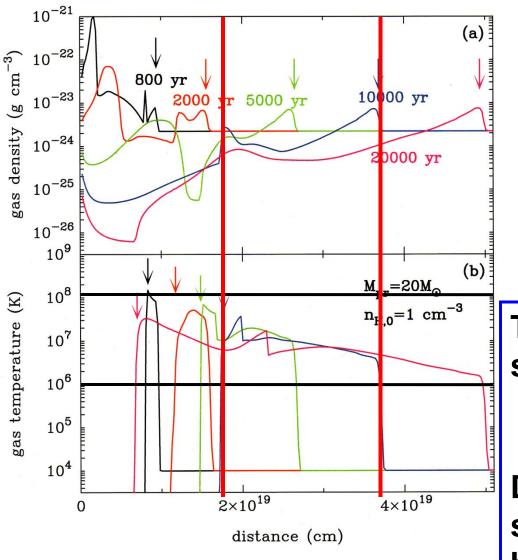
Nozawa+2006, ApJ, 648, 435



 erosion rate by sputtering quickly increases above 10⁵ K and peaks at 10⁷-10⁸ K

erosion rate : da / dt ~ 10⁻⁶ n_H µm yr⁻¹ cm³
 for the primordial gas (H and He) at T > 10⁶ K

3-5. Temperature and density of gas in SNRs



Nozawa+07, ApJ, 666, 955

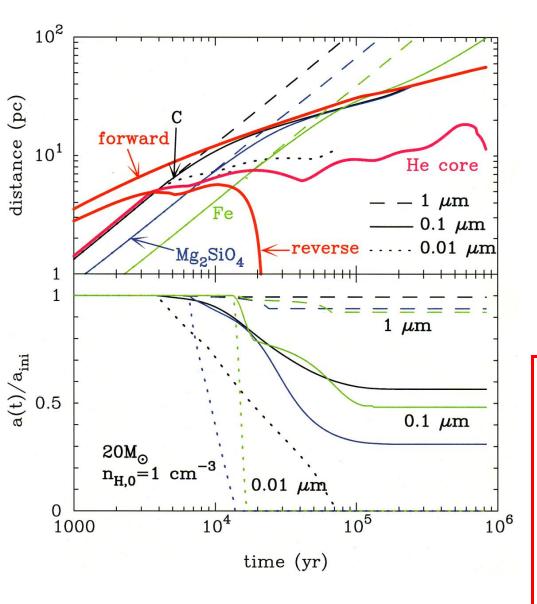
Model :
$$M_{pr}$$
= 20 Msun (E₅₁=1)
n_{H,0} = 1 cm⁻³

Downward-pointing arrows: forward shock in upper panel reverse shock in lower panel

The temperature of the gas swept up by the shocks → 10⁶-10⁸ K ↓ Dust grains residing in the

shocked hot gas are eroded by sputtering

3-6. Evolution of dust in SNRs



Nozawa+07, ApJ, 666, 955

Model : M_{pr} = 20 Msun (E₅₁=1) n_{H,0} = 1 cm⁻³

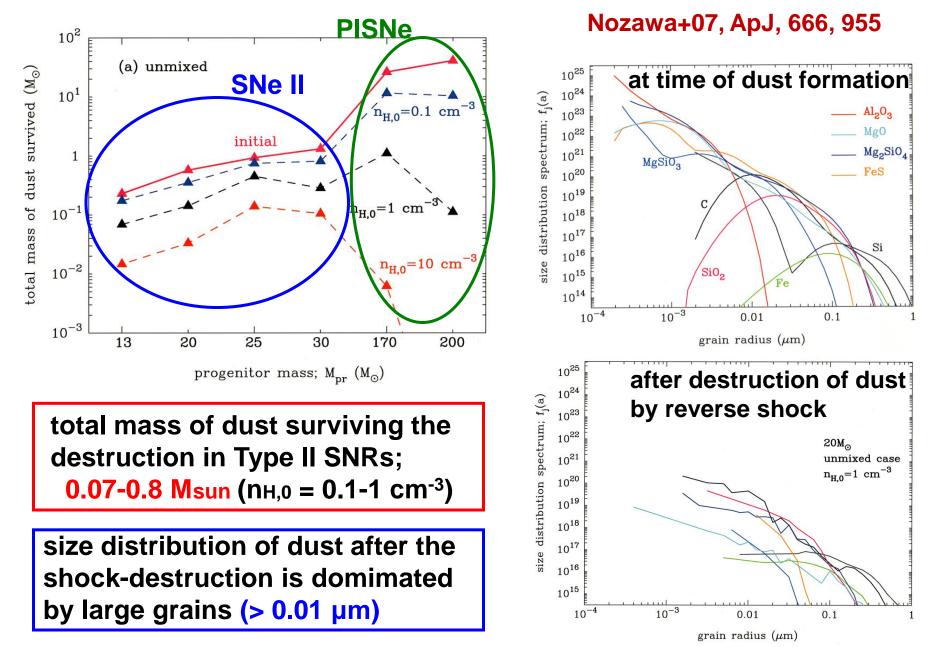
Dust grains in the He core collide with reverse shock at (3-13)x10³ yr

The evolution of dust heavily depends on the initial radius and composition

- a_{ini} = 0.01 μm (dotted lines) → completely destroyed
- a_{ini} = 0.1 μm (solid lines) → trapped in the shell

a_{ini} = 1 μm (dashed lines) → injected into the ISM

3-7. Dust mass and size ejected from SN II-P



3-8. Summary of dust production in Pop III SNe

- Various grain species can condense in the ejecta
 → almost all Fe, Mg, and Si are locked up in grains
- The fate of newly formed dust within SNRs strongly depends on the initial radii and compositions
- The size distribution of dust surviving the destruction in SNRs is weighted to relatively large size (> 0.01 μm).
- The total mass of dust injected into the ISM decreases with increasing the ambient gas density

for $n_{H.0} = 0.1-1 \text{ cm}^{-3}$

SNe II-P → Mdust = 0.1-0.8 Msun

→ significant contribution to dust budget at high z