

種族III巨大質量星の赤色超巨星星風中 におけるダスト形成

Formation of Carbon Dust in Red-supergiant
Winds of Very Massive Population III Stars

Nozawa et al. (2014, ApJL, 787, L17)

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0-1. My research interests

- **Formation of dust in the ejecta of supernovae**
- **Dynamics, destruction, and heating of dust grains in high-velocity shock waves**
- **Evolution of grain size distribution in galaxies and interstellar extinction curves**
- **Origin of dust grains in the early universe and the roles of dust in star formation**
- **Formation of dust in stellar winds and mass-loss history at late phases of stellar evolution**

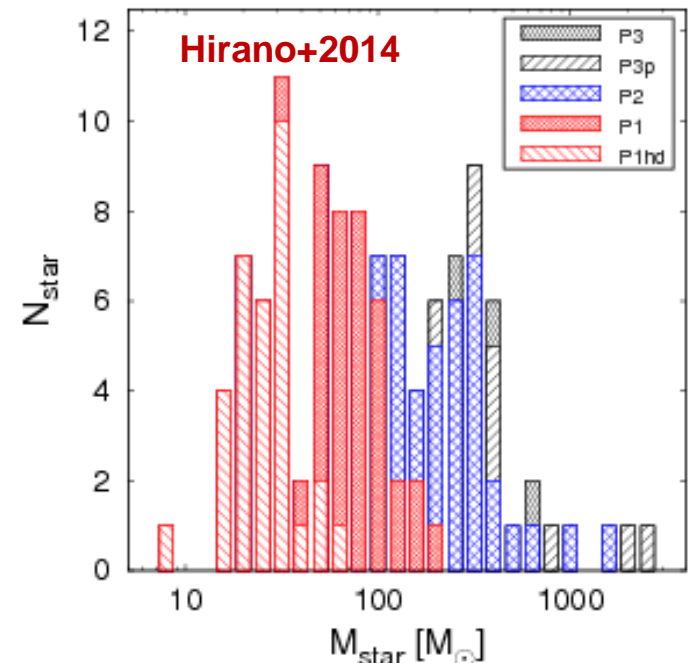
1-1. Sources of dust in the early universe

• Origin of massive dust at high redshifts ($z > 5$)

- **core-collapse supernovae (CCSNe)** may be promising sources of dust grains (e.g., Todini & Ferrara 2001; TN+2003; Dwek+2007)
- the contribution from **AGB stars** is also invoked to explain the observed dust mass (e.g., Valiante+2009; Dwek & Cherchneff 2011)
 - what stellar mass range can mainly contribute dust budget in the early universe depends on the stellar IMF

• Typical mass of Pop III stars

- Pop III stars may be much more massive than Pop I/II stars
- ~40 M_{sun} (Hosokawa+2011; Susa 2013)
- >300 M_{sun} (Omukai+2003; Ohkubo+2009)
- 10-1000 M_{sun} (Hirano+2014, Susa+2014)



1-2. Very massive Population III stars

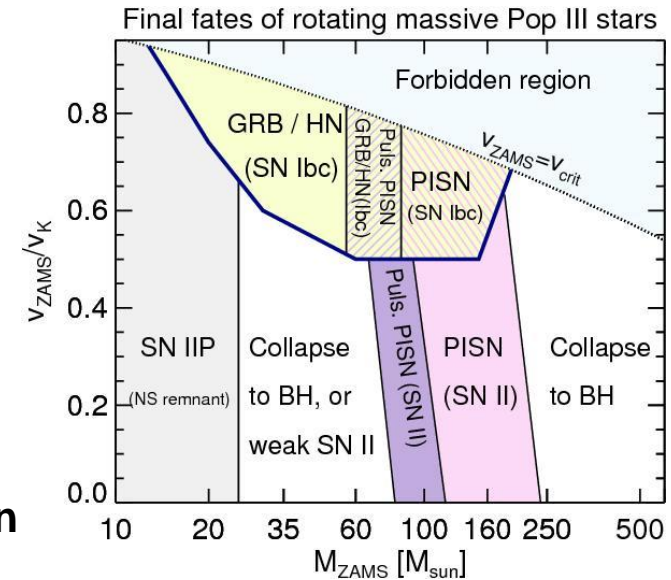
• Role of very massive stars ($M_{ZAMS} > \sim 250 M_{\text{sun}}$)

- emitting numerous ionizing photons
→ reionization of the universe
- finally collapsing into black holes
→ serving as seeds of SMBHs

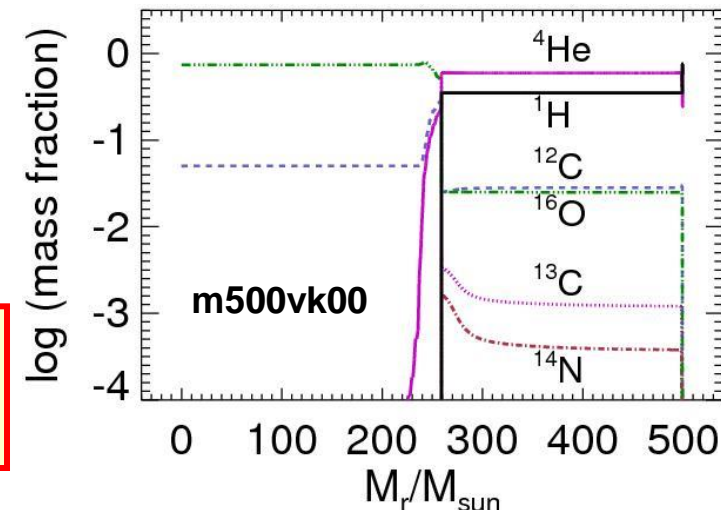
• Evolution of massive Pop III stars

- non-rotating stars with $M_{ZAMS} > 250 M_{\text{sun}}$ undergo convective dredge-up of C and O during the RSG phase (Yoon+2012)
- enriching the surrounding medium with CNO through the RSG winds
→ serving as formation sites of dust

Dust grains formed in the winds are not likely to be destroyed by the SN shocks



Yoon+2012



2-1. Model of Pop III red-supergiant winds

▪ RSG model: m500vk00 (Yoon+2012)

- MZAMS = 500 M_{sun} (no rotation)
- L = 10^{7.2} L_{sun}, T_{star} = 4440 K, R_{star} = 6750 R_{sun}
- AC = 3.11x10⁻³, A_O = 1.75x10⁻³ → **C/O = 1.78, Z = 0.034**

▪ Model of circumstellar envelope

- spherically symmetry, constant wind velocity

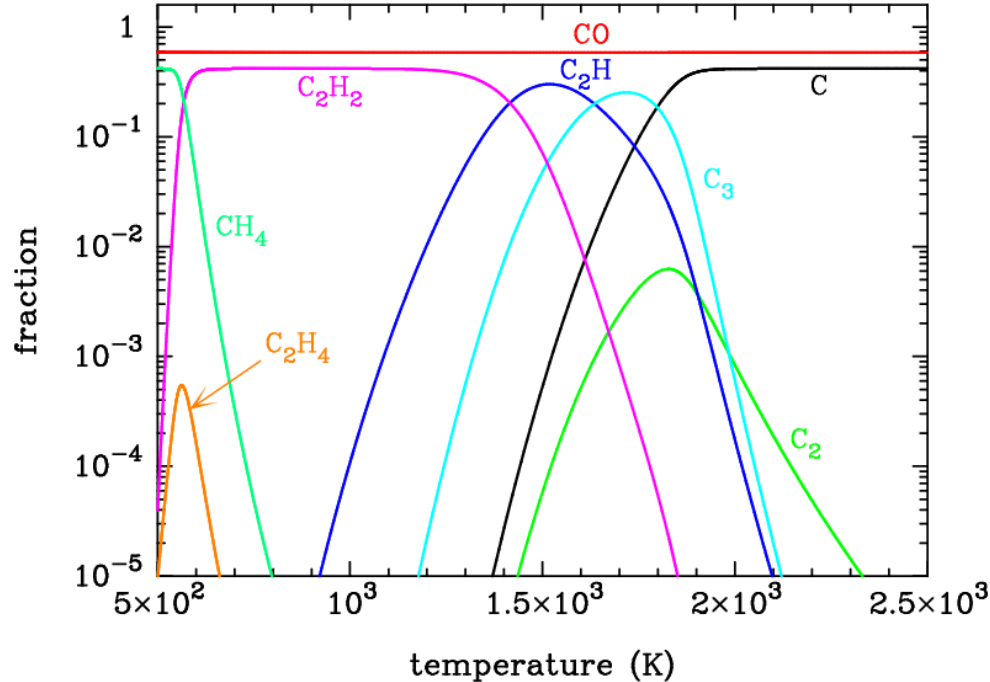
- density profile:
$$\rho(r) = \frac{\dot{M}}{4\pi r^2 v_w} = \rho_* \left(\frac{r}{R_*} \right)^{-2}$$

- temperature profile:
$$T(r) = T_* \left(\frac{r}{R_*} \right)^{-\frac{1}{2}}$$

▪ Fiducial values of M_{dot} and V_w

- **wind velocity: v_w = 20 km/s**
- **mass-loss rate: M_{dot} = 0.003 M_{sun}/yr**
 - losing 90% (208 M_{sun}) of envelope during 7x10⁴ yr

2-2. Chemical equilibrium calculations



major carbon-bearing gas species other than CO:

- atomic carbon
at $T > \sim 1800\text{K}$
- C_2H molecules
at $T = 1400\text{-}1700\text{ K}$

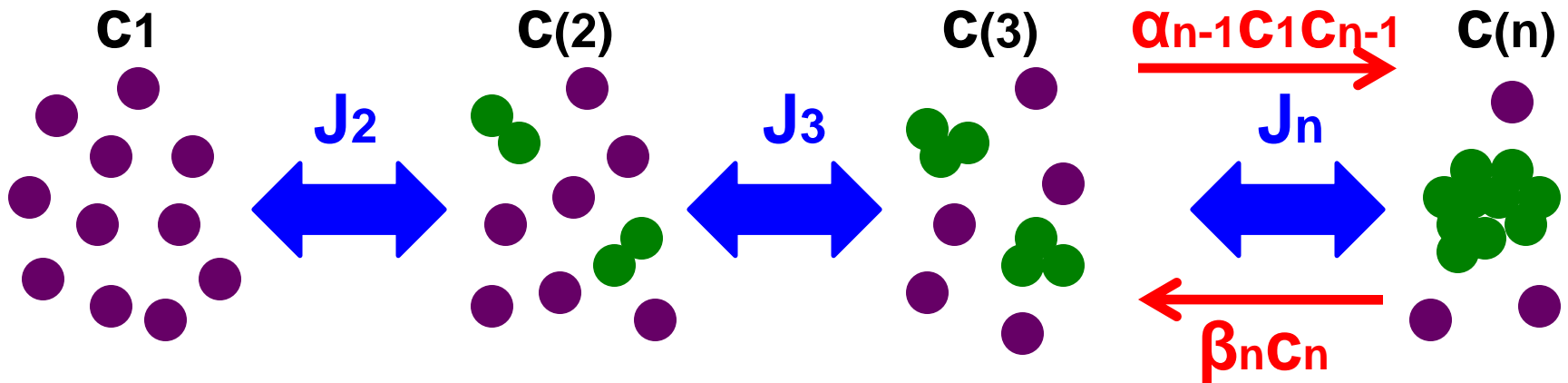
Formation of PAHs would not
be expected

chemical reactions considered in this study

(1) Model A	C	$\text{C}_{n-1} + \text{C} \rightleftharpoons \text{C}_n$	$(n \geq 2)$
(2) Model B	C_2H	$2(\text{C}_2\text{H} + \text{H}) \rightleftharpoons \text{C}_{2n} + 2\text{H}_2$	$(n = 2)$
		$\text{C}_{2(n-1)} + \text{C}_2\text{H} + \text{H} \rightleftharpoons \text{C}_{2n} + \text{H}_2$	$(n \geq 3)$

- Dust formation calculations (TN & Kozasa 2013)
formulation of non-steady-state dust formation

3-1. Concept of nucleation theory



▪ master equations

$$\frac{dc_n}{dt} = J_n(t) - J_{n+1}(t) \quad \text{for } 2 \leq n \leq n_*,$$

$$J_n(t) = \alpha_{n-1} c_{n-1} c_1 - \beta_n c_n \quad \text{for } 2 \leq n \leq n_*,$$

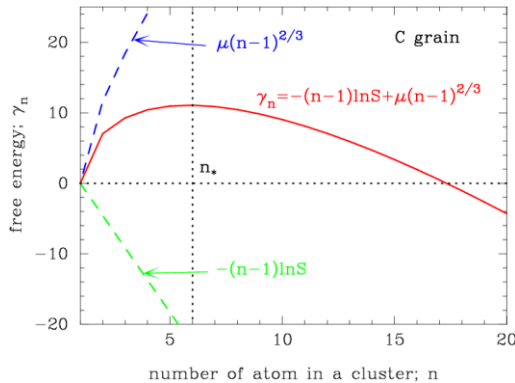
$$\alpha_n = \frac{s_n}{1 + \delta_{1n}} 4\pi a_0^2 n^{\frac{2}{3}} \left(\frac{kT}{2\pi m_n} \right)^{\frac{1}{2}},$$

$$\beta_n = \alpha_{n-1} \frac{\overset{\circ}{c}_{n-1}}{\overset{\circ}{c}_n} \overset{\circ}{c}_1,$$

3-2. Non-steady-state nucleation

▪ steady-state nucleation rate: J_s

→ assuming $J_s = J_2 = J_3 = \dots = J_\infty$



$$(n_c - 1)^{\frac{1}{3}} = \frac{2}{3} \frac{\mu}{\ln S}$$

where $\mu = 4\pi a_0^2 \sigma / kT$

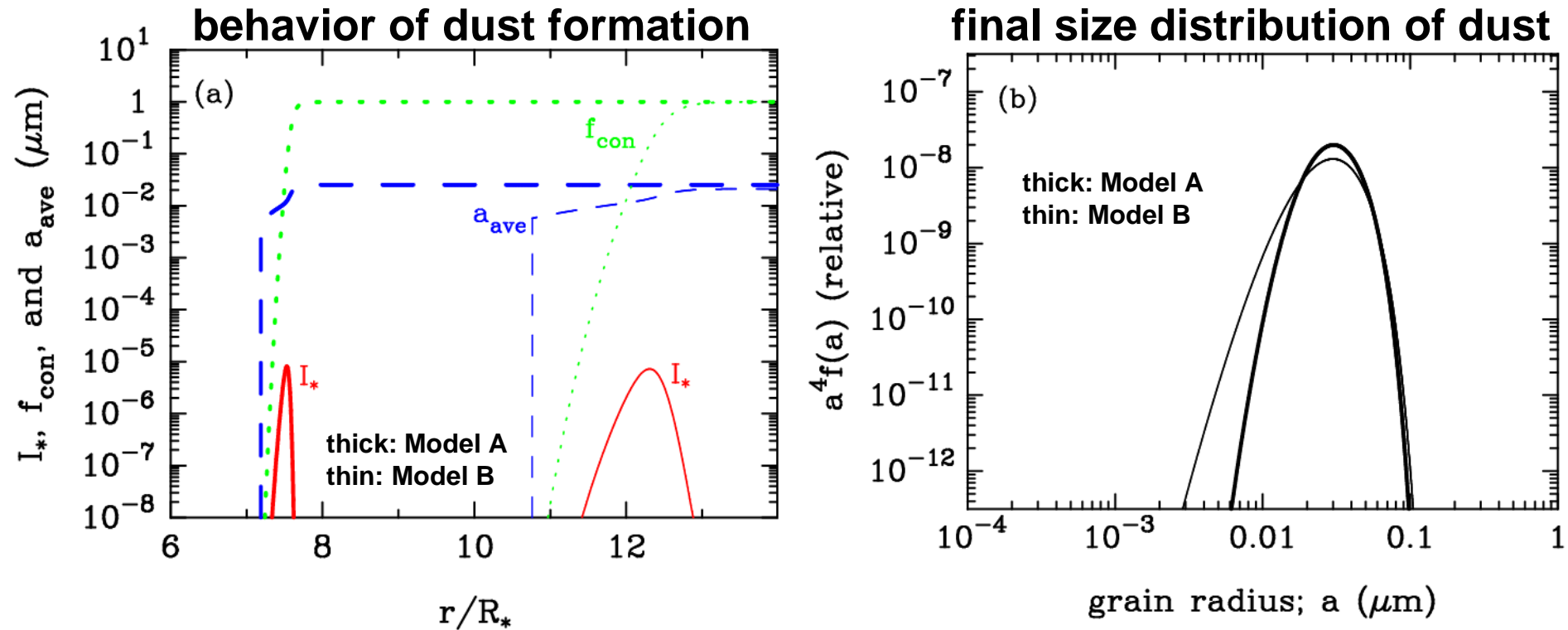
$$J_s = s \Omega_0 \left(\frac{2\sigma}{\pi m_1} \right)^{\frac{1}{2}} c_1^2 \exp \left[-\frac{4}{27} \frac{\mu^3}{(\ln S)^2} \right]$$

▪ non-steady-state dust formation

$n^* = 100$

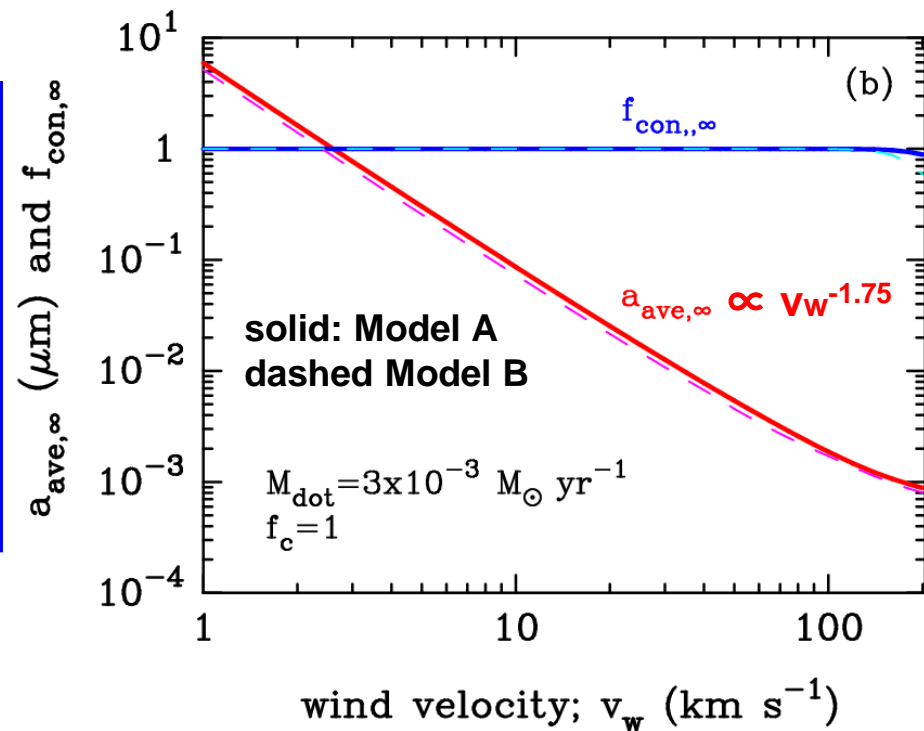
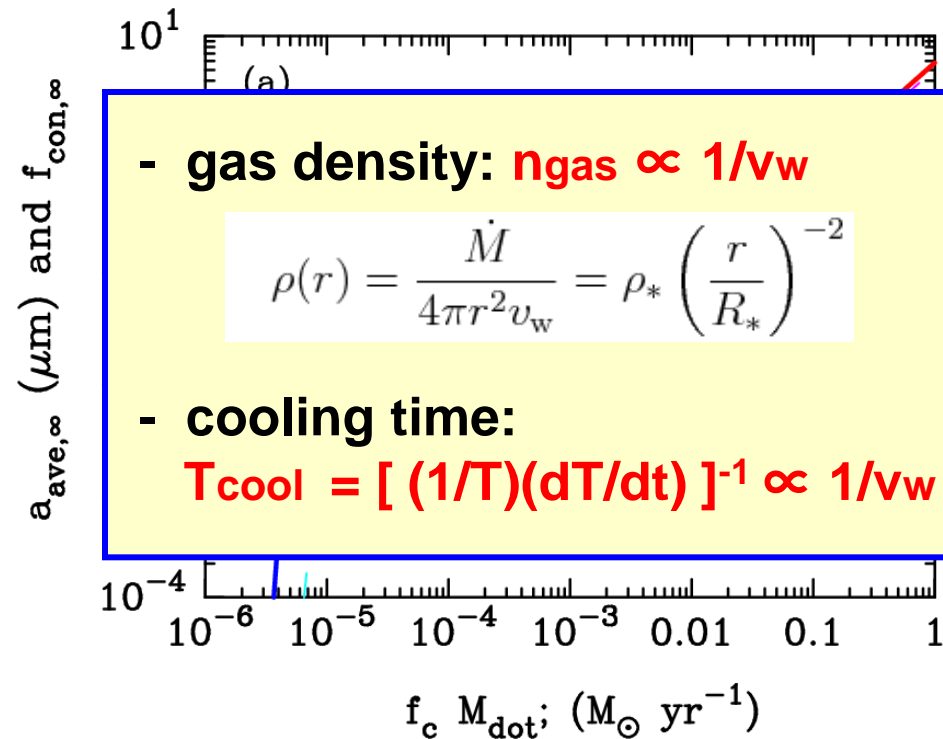
$$\frac{dc_n}{dt} = J_n(t) - J_{n+1}(t) \quad \text{for } 2 \leq n \leq n^*,$$

4-1. Results of dust formation calculations



- carbon grains form around $r = 7.5 R_{\text{star}}$ ($r = 12 R_{\text{star}}$) for Model A (Model B)
 - final condensation efficiency is unity for both of the models
 - final average radius is similar in both Model A and Model B
- the results are almost independent of chemical reactions**

4-2. Dependence on \dot{M} and v_w



- The condensation efficiency of dust is unity for the condition;

$$\left(\frac{f_c \dot{M}}{3 \times 10^{-3} M_{\odot} \text{ yr}^{-1}} \right) \left(\frac{v_w}{20 \text{ km s}^{-1}} \right)^{-2} \gtrsim 0.04.$$

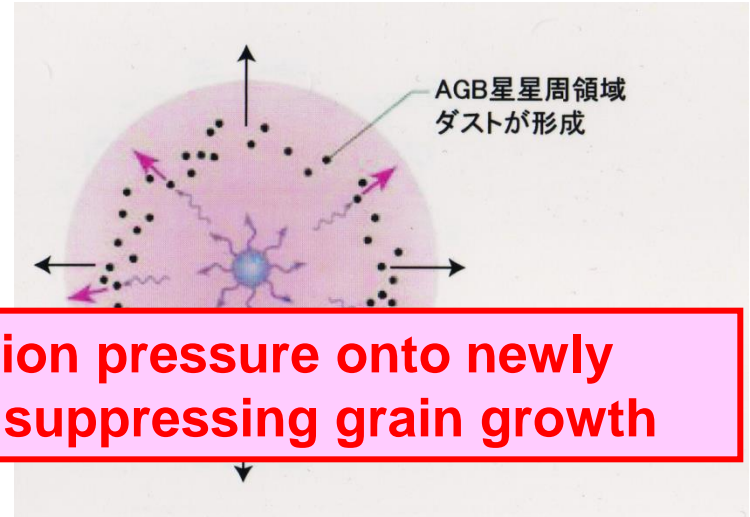
- for the fiducial case ($\dot{M} = 3 \times 10^{-3} M_{\text{sun}}/\text{yr}$, $v_w = 20 \text{ km/s}$, $f_c = 1$)
 $\rightarrow 1.7 M_{\text{sun}}$ of C grains is produced over the lifetime of the RSG

4-3. Models of dust-driven winds

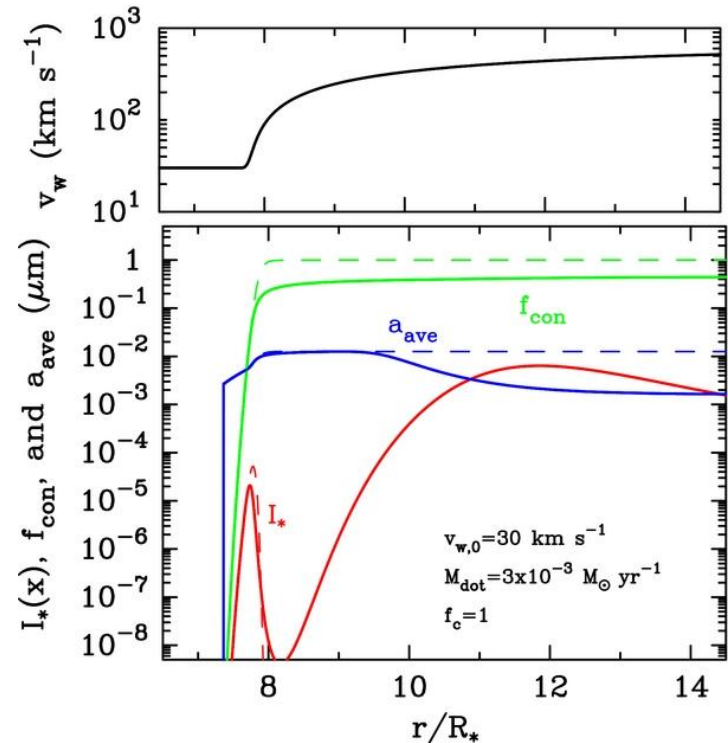
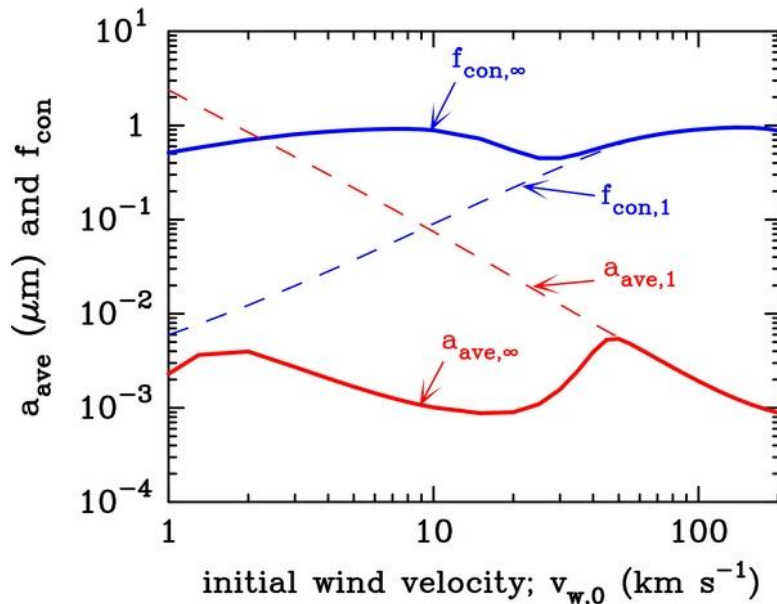
- dust-driven winds

- dust formation in the winds
- radiation pressure onto dust
- acceleration of dust

The acceleration of the wind by radiation pressure onto newly formed dust reduces the gas density, suppressing grain growth



$$v_w \frac{dv_w}{dr} = -\frac{GM_*}{r^2} \left[1 - \frac{L_* \langle \kappa_{\text{ext}}(T) \rangle D}{4\pi c GM_*} \right],$$



5-1. How efficient is dust formation?

Dust ejection efficiency by very massive Pop III RSGs

- $X_{VMS} = M_{dust} / M_{ZAMS} < 3.4 \times 10^{-3} = \sim 0.3 \%$
- $M_{dust} / M_{metal} < 0.24$

Dust ejection efficiency by CCSNe

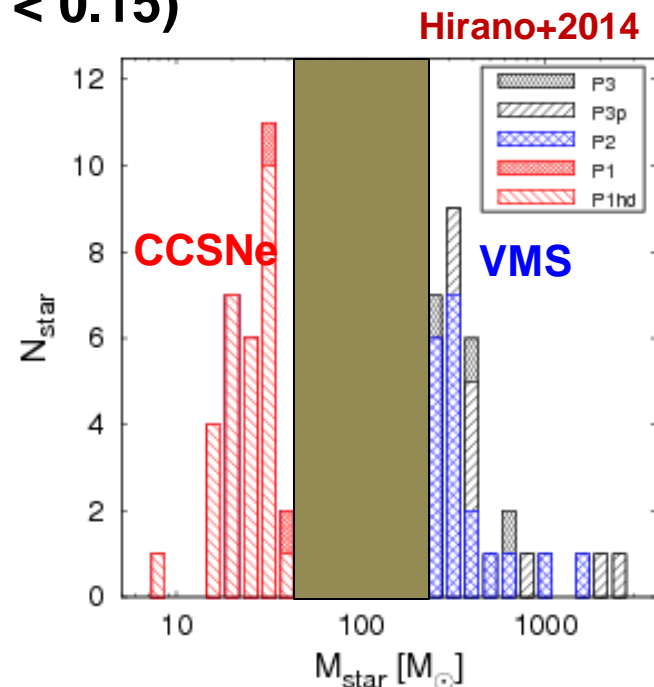
- $X_{CCSN} = (0.1-30) \times 10^{-3} = 0.1-3.0 \%$
- $M_{dust} / M_{metal} = 0.01-0.25$ ($M_{dust} / M_{metal} < 0.15$)

The ranges above reflects the destruction
efficiency of dust by the reverse shock

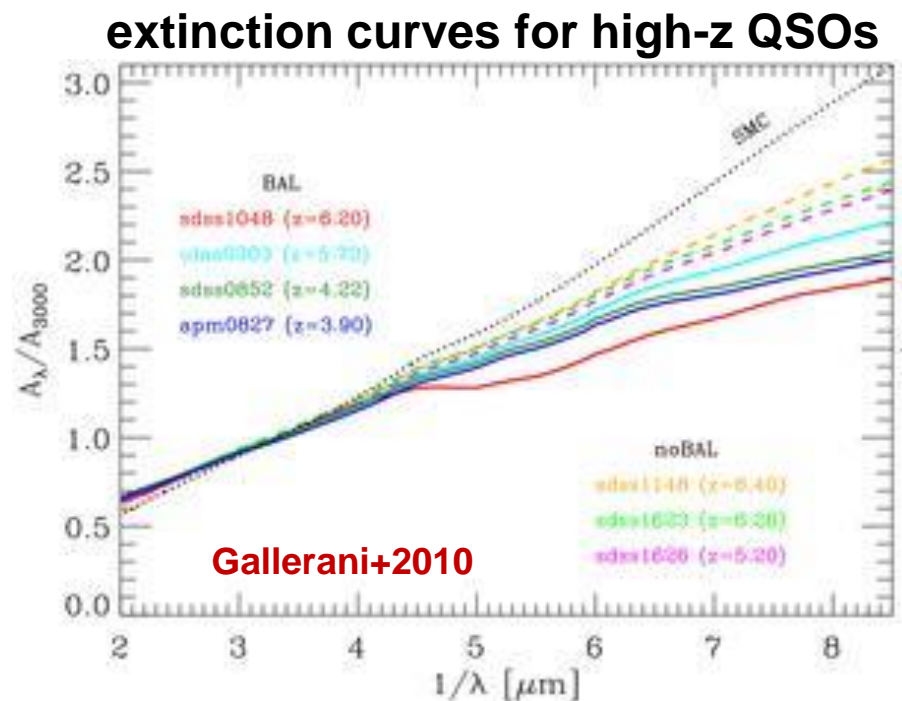
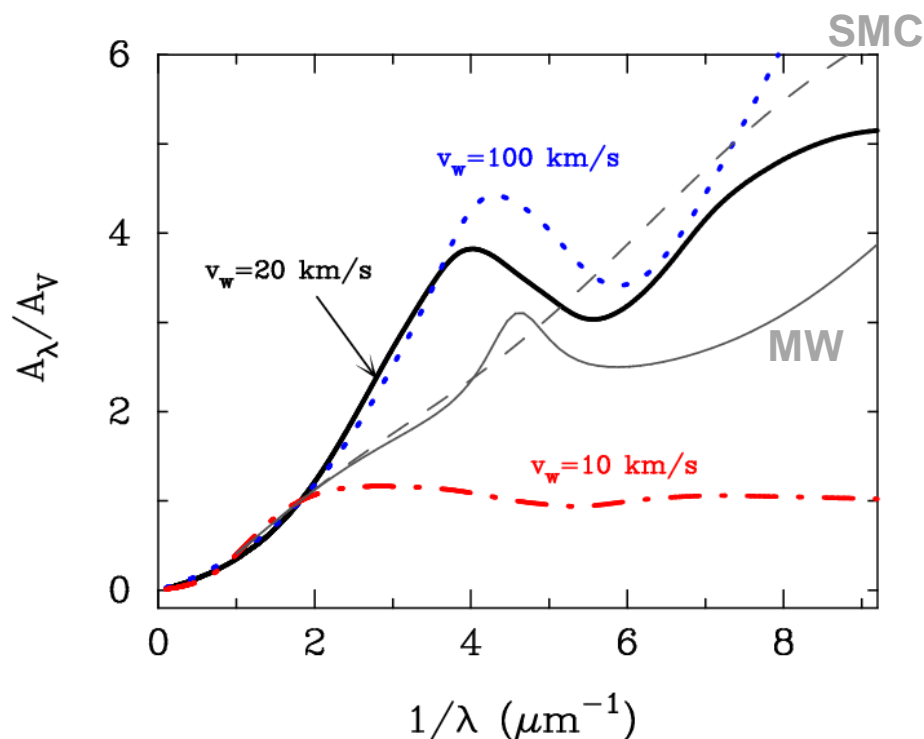
If $N_{VMS} \sim N_{CCSN}$ in the Pop III IMF ...

→ The contribution of dust from very massive RSGs is comparable with, or even higher than that from CCSNe

$$(X_{VMS} N_{VMS}) / (X_{CCSN} N_{CCSN}) > \sim 1$$



5-2. Expected extinction curves



- Extinction curves derived in this study do not resemble any of the known extinction law such as those in the MW and SMC
- The extinction curves observed for high- z quasars do not show a bump structure, being inconsistent with those given here
 - The derived extinction curves can be powerful tools to probe the formation of C grains in very massive Pop III stars

5-3. Composition of low-mass UMP stars

- The ultra-metal-poor (UMP) stars with $[\text{Fe}/\text{H}] < -4$ would record chemical imprints of Population III stars
- The formation of such low-mass metal-poor stars is triggered through the cooling of gas by dust produced by Pop III SNe
(e.g., Schneider+2012a, 2012b; Chiaki+2014)

Possible channel for C-rich UMP star formation

- Very massive Pop III RSGs are sources of carbon grains as well as CNO elements
 - In the gas clouds enriched by Pop III RSGs, carbon grains enable the formation of CNO-rich low-mass stars
 - We do not predict the presence of heavier elements (Mg, Si, Fe)
 - Further observations and more quantitative theoretical studies are needed to show whether UMP stars formed through our scenario
- ## SMSS J0313-6708: $[\text{C}/\text{H}] = -2.6$, $[\text{Fe}/\text{H}] < -7$ (Keller+2014)**

6. Summary

We have examined the possibility of dust formation in a carbon-rich mass-loss wind of a Pop III RSG with $M_{ZAMS} = 500 M_{\text{sun}}$

- For a steady stellar wind, C grains can form with a lognormal-like size distribution whose average radius is sensitive to wind velocity
- The condensation efficiency is unity for

$$\left(\frac{f_c \dot{M}}{3 \times 10^{-3} M_{\odot} \text{ yr}^{-1}} \right) \left(\frac{v_w}{20 \text{ km s}^{-1}} \right)^{-2} \gtrsim 0.04.$$

→ the first dust grains in the universe ??

- The mass of C grains is $< 1.7 M_{\text{sun}}$ ($M_{\text{dust}}/M_{ZAMS} < 3.4 \times 10^{-3}$), which would be high enough to have impacts on dust enrichment history in the early universe, if the IMF of Pop III stars were top-heavy

The extinction curves expected from ejected C grains are different from any known ones

The chemical feedback by PopIII VMSs predicts a new type of UMP stars