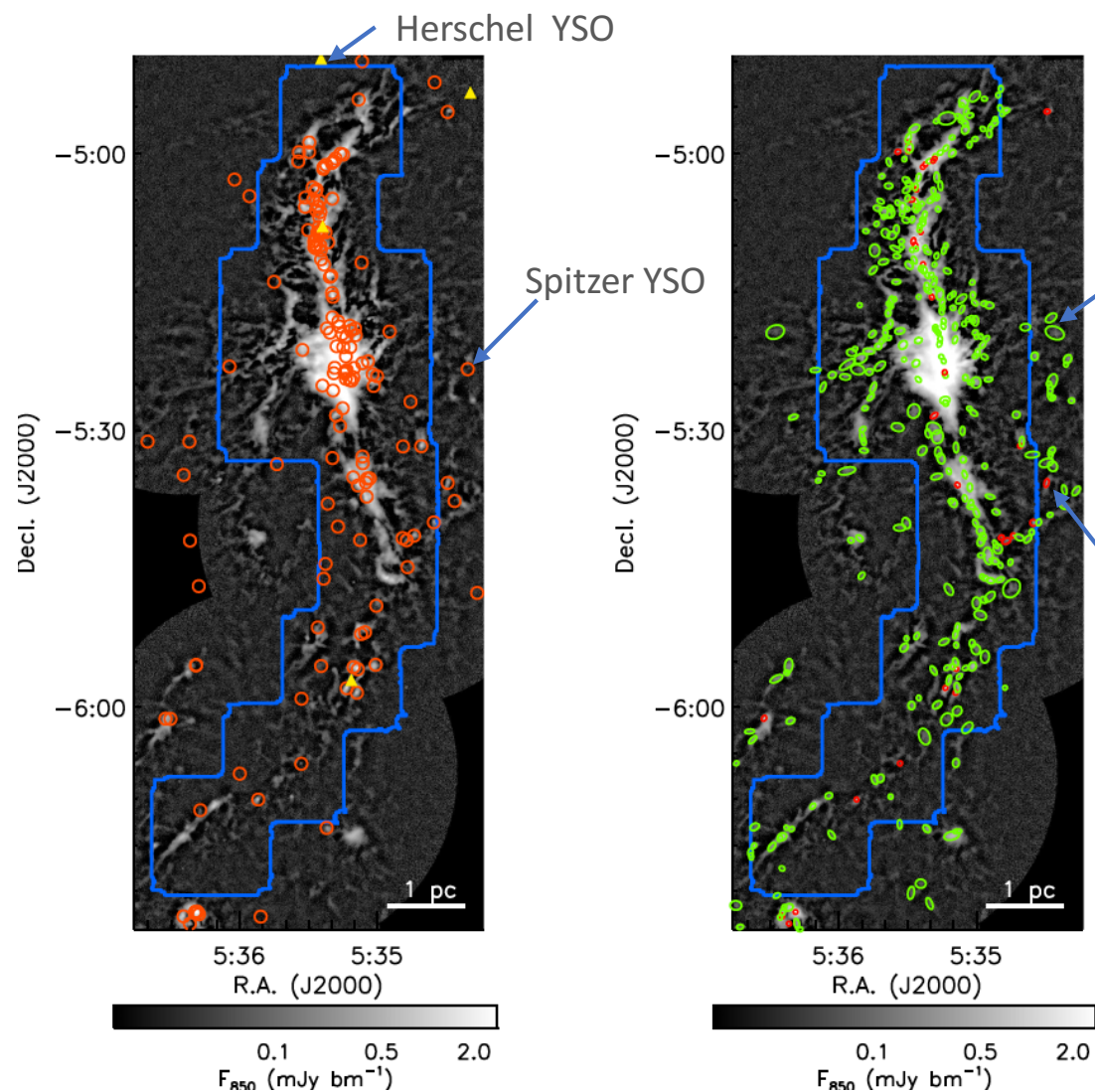
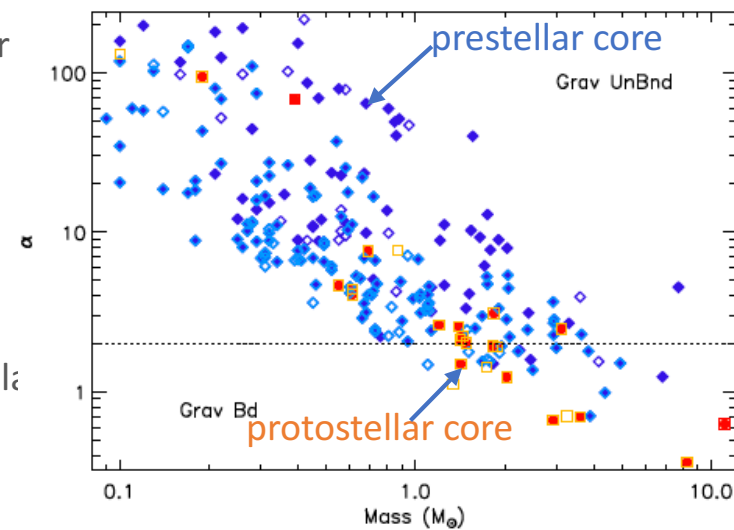


21. The Green Bank Ammonia Survey: Dense Cores Under Pressure in Orion A Kirk+ ApJ in press

GAS



$$\alpha = \frac{5\sigma_{\text{tot}}^2 R}{GM},$$

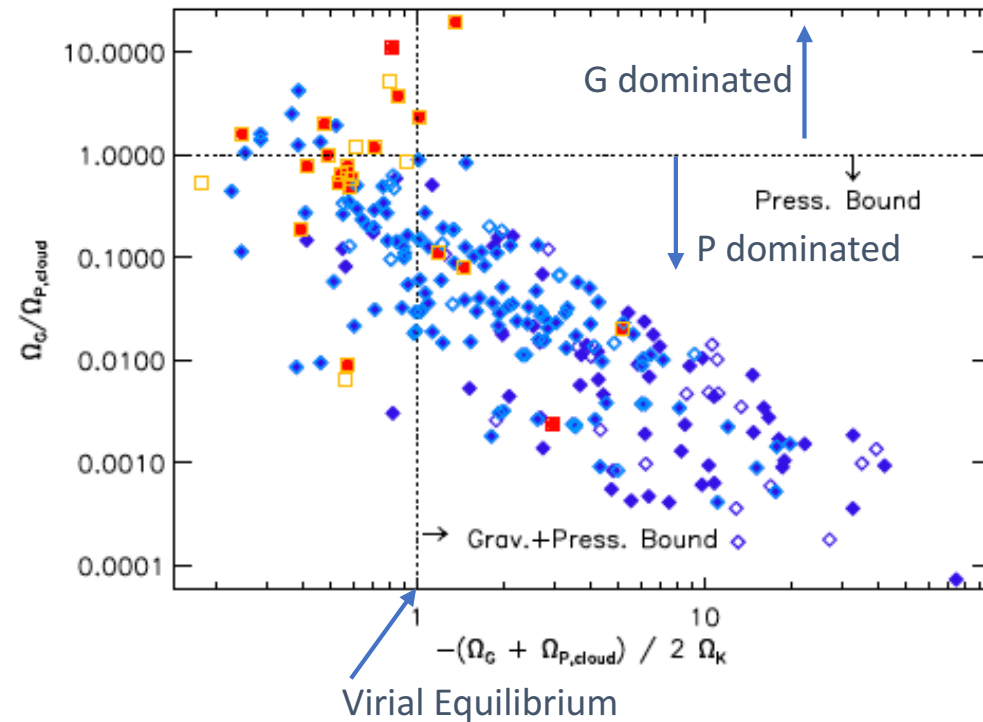


surface pressure $\Omega_P = -4\pi PR^3,$

potential $\Omega_G = \frac{-1}{2\sqrt{\pi}} \frac{GM^2}{R},$

kinetic energy $\Omega_K = \frac{3}{2}M\sigma_{\text{tot}}^2,$

$-(\Omega_G + \Omega_P) = 2\Omega_K$



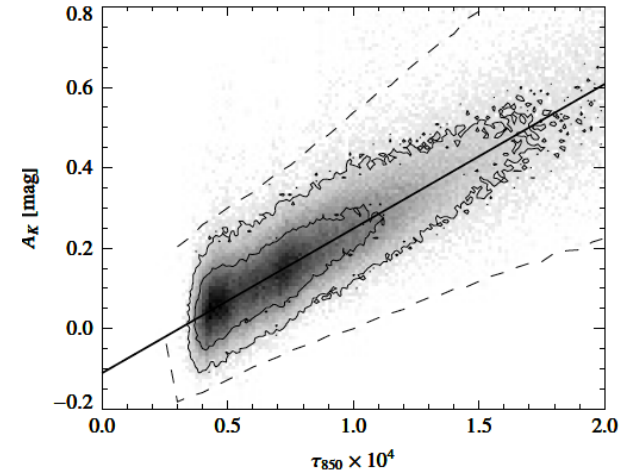
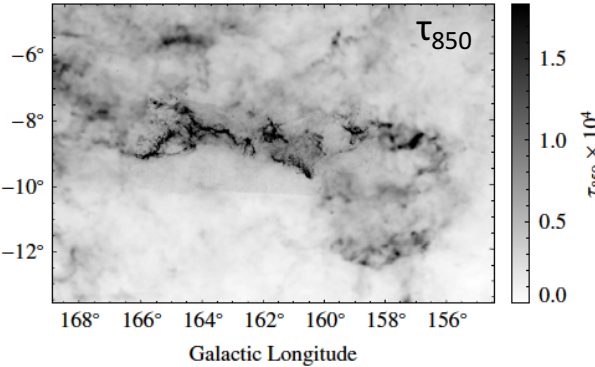
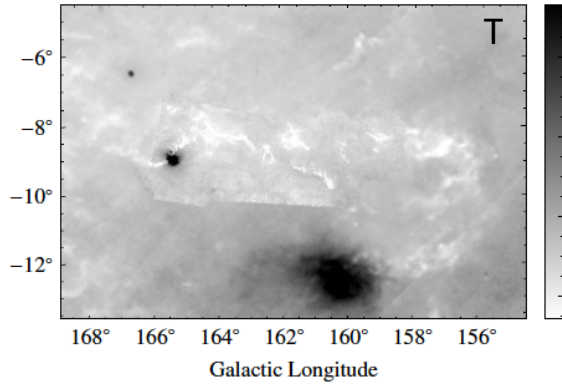
Surprisingly, we find that almost none of the dense cores are sufficiently massive to be bound when considering only the balance between self-gravity and the thermal and non-thermal motions present in the dense gas. Including the additional pressure binding imposed by the weight of the ambient molecular cloud material and additional smaller pressure terms, however, suggests that most of the dense cores are pressure-confined.

22. HP2 survey: III The California Molecular Cloud– A Sleeping Giant Revisited

by Lada, Lewis, Lombardi, Alves Aap in press

SFR = 1/10 of Ori

Herschel data SED fit $\rightarrow (T, \tau_{850})$



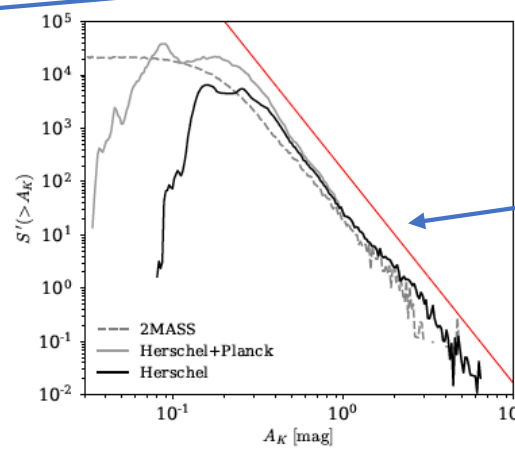
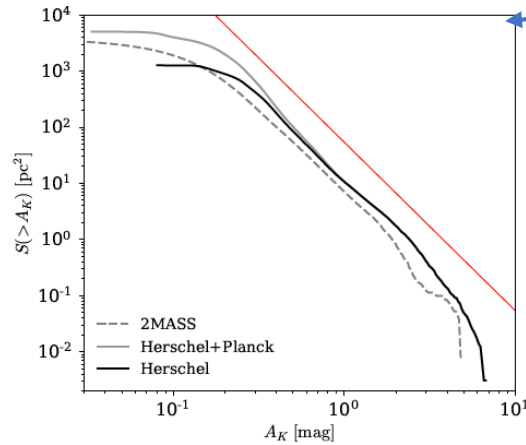
$$I_\nu = B_\nu(T)[1 - e^{-\tau_\nu}] \simeq B_\nu(T)\tau_\nu, \quad \tau_\nu = \tau_{\nu_0} \left(\frac{\nu}{\nu_0} \right)^{\beta_d},$$

$$A_K = \gamma \tau_{850} + \delta$$

$$\gamma = 3593.75 \text{ mag}$$

$$\delta = -0.110 \pm 0.006 \text{ mag}$$

Structure Function



$$S(> A_K) \equiv \int_{A_K}^{\infty} dS(A_K) \propto A_K^{-q}$$

where $dS(A_K)$ is an element of cloud surface area at an extinction A_K .

derivative of this function

$$dS(> A_K)/dA_K = S'(> A_K) \propto A_K^{-n}$$

proportional to the column density PDF

$$\Rightarrow n = 4.0 \pm 0.1$$

$$\Rightarrow q = 3.0 \quad (q=2 \text{ Ori A, B and Per})$$

\Rightarrow steeper slope than the GMCs

Fig. 8. The surface area distribution functions, $S(> A_K)$, for the California cloud. The solid red line represents a power law relation with a slope of -3 for comparison.

Fig. 9. The differential area function $-S'(> A_K)$ which is proportional to the probability density distribution for column densities in the cloud. In this plot a lognormal distribution would appear as a gaussian function whilst a power-law function would be a straight line. The red line shows the power-law $-S'(> A_K) \propto A_K^{-4.0}$. See text.

Schmidt Law

protostellar surface density: thresholded Schmidt relation

$$\Sigma_*(A_K) = \kappa A_K^{\beta} H(A_K - A_{K,0}),$$

→ $\beta = 3.31 \pm 0.23$, $\kappa = 0.36 \pm 0.09$ stars $\text{pc}^{-2} \text{mag}^{-3.31}$, $A_{K,0} = 0.51$ mag.

Table 1. Schmidt relation and PDF power-law indices for GMCs studied with *Herschel*

GMC	β	Reference	n	Reference
California	3.31 ± 0.23	1	4.0	1
Orion A	1.99 ± 0.05	2	2.9	4
Orion B	2.16 ± 0.10	2	3.0	4
Perseus	2.4 ± 0.6	3	2.7	4

References. (1) This paper; (2) Lombardi et al. (2014); (3) Zari et al. (2016) (4) Lombardi et al. (2015).

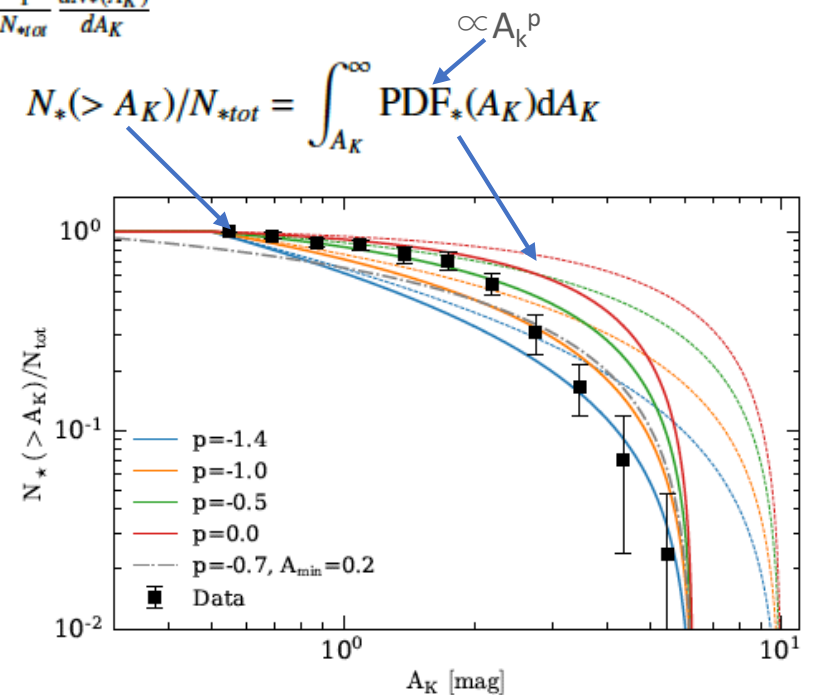
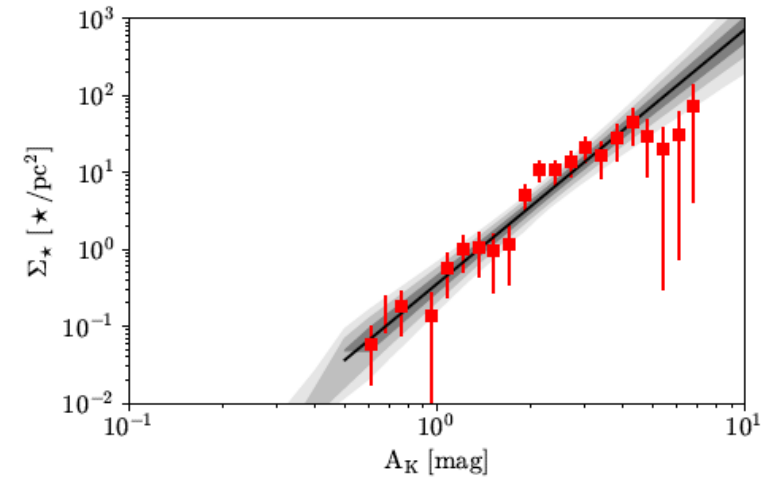
$$\Sigma_*(A_K) = \frac{dN_*(A_K)}{dS(A_K)} = \Sigma_{*0} \times \frac{\text{PDF}_*(A_K)}{\text{PDF}_N(A_K)}$$

measured (for Σ_*) unknown $\text{PDF}_*(A_K) \equiv \frac{1}{N_{*tot}} \frac{dN_*(A_K)}{dA_K}$ measured (for PDF_N)

$$\text{PDF}_*(A_K) = \frac{\Sigma_*(A_K)}{\Sigma_{*0}} \times \text{PDF}_N(A_K)$$

if power law → $p = \beta - n = -0.69 \pm 0.27$

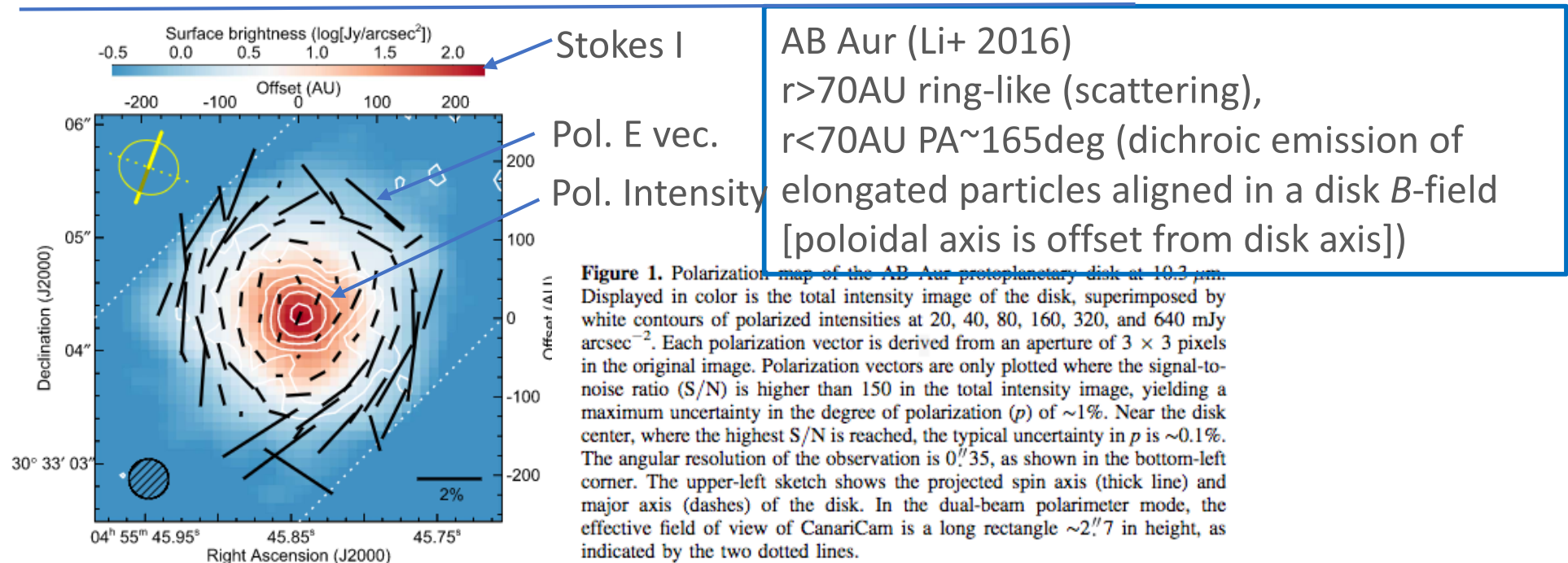
Our observations suggest that variations both in the slope of the Schmidt relation and in the sizes of the protostellar populations between GMCs are largely driven by variations in the slope, n , of $\text{PDF}_N(A_K)$.

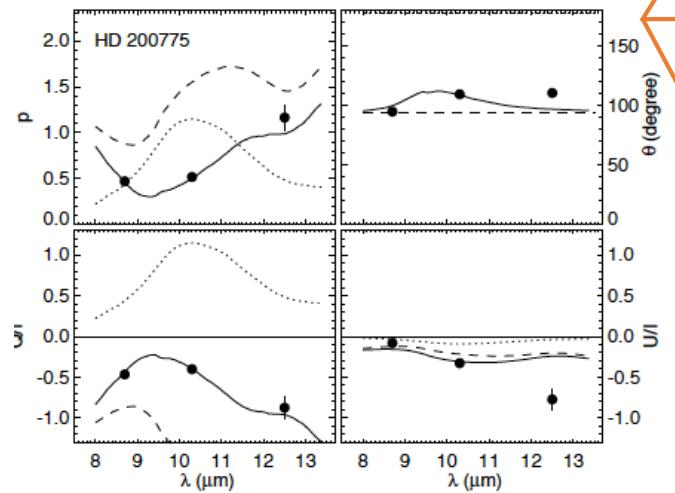
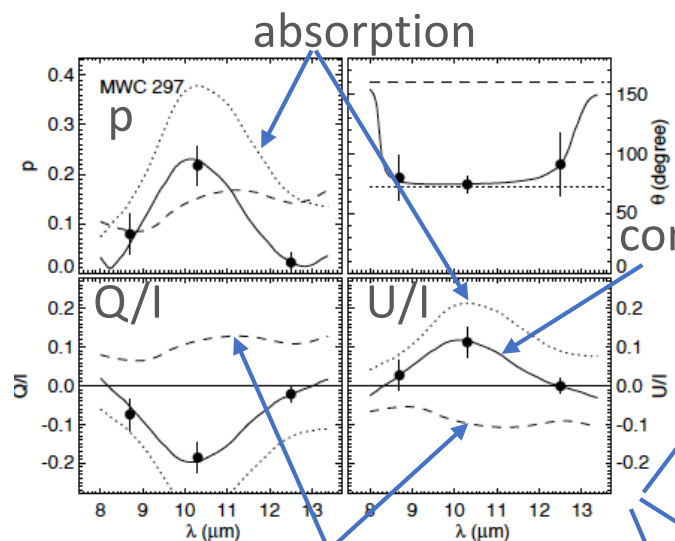


23. Mid-Infrared Polarization of Herbig Ae/Be Discs by Dan Li+, MNRAS in press

The translation of mm/submm polarization maps into the B-field morphology (projected on the plane-of-sky) was once thought to be straightforward, but recent studies have emphasized that scattered emission can also contribute to the observed mm/submm polarization if dust grains much larger than their ISM (interstellar medium) counterparts are present in discs (Kataoka et al. 2015; Yang et al. 2016; Tazaki et al. 2017).

- (1) Observation: a mid-infrared (mid-IR) polarimetry using Canari-Cam
- (2) Object: Herbig Ae/Be (HAeBe) stars (i.e., premain sequence stars of 2–8 M_{\odot}).
- (3) detected: the brightest and most compact inner regions of the discs.
- (4) Band: Si-2 ($\lambda=8.7\mu\text{m}$), Si-4 ($\lambda=10.3\mu\text{m}$), Si-6 ($\lambda=12.5\mu\text{m}$)





not inconsistent with
2 component model
→ scattering

Object	λ (μm)	F_ν (Jy)	p (per cent)	θ (degree)
MWC 297	8.7	115.9 ± 11.6	0.08 ± 0.05	80 ± 19
	10.3	103.2 ± 10.3	0.22 ± 0.05	74 ± 7
	12.5	111.6 ± 11.2	0.02 ± 0.03	91 ± 26
HD 200775	8.7	7.8 ± 0.8	0.47 ± 0.09	105 ± 2
	10.3	7.3 ± 0.7	0.52 ± 0.15	110 ± 1
	12.5	5.3 ± 0.5	1.2 ± 0.1	111 ± 1

Properties of discs and interstellar B-fields.

Object	10- μm Silicate Feature Type	Disc P.A. (degree)	Inclination (degree)	Ref.	IS B-field P.A. ^a (degree)	Ref.
MWC 1080A	Abs	135 ± 5	55 ± 5	1,2	70	19
MWC 297	Abs	165 ± 15	13 ± 5	1,3	40	19,20
HD 200775	Em	6.9 ± 1.5	54.5 ± 1.2	1,4	142	19
VV Ser	Em	15 ± 5	72 ± 5	5	60	20,21
HD 179218	Em	22 ± 3	56 ± 2	1,9	4	20
AB Aur	Em	70 ± 10	27 ± 5	6-8	70	19,22
HD 163296	Em	136 ± 2	48 ± 2	10	175	23
MWC 480	Em	148 ± 2	38 ± 5	11-13	—	
MWC 758	Em	65 ± 7	21 ± 2	14	—	
CQ Tau	Em	54 ± 1	29 ± 2	15,16	42	24
HL Tau	Abs	138.02 ± 0.07	46.72 ± 0.05	17,18	77	25

Object	Scattering	Emission	Absorption	B-field
MWC 1080A			✓	Complex or tilted
MWC 297		✓	✓	Poloidal ^a
HD 200775	✓			
VV Ser	✓	✓	✓	Toroidal
HD 179218	✓	✓		Toroidal
AB Aur	✓	✓		Tilted poloidal (L16)
HD 163296				
CQ Tau		✓		Complex or tilted
HL Tau			✓	Complex or tilted

Possible origin(s) of polarization. See Section 4.3 for details.

^a Poorly constrained. See Section 4.3.

24. The Formation of Stellar Clusters in Magnetized, Filamentary Infrared Dark Clouds

Pak Shing Li et al. MNRAS in press

InfraRed Dark Cloud $\Sigma \sim 10^{22} \text{cm}^{-2}$ $\rho \sim 10^4 \text{cm}^{-3}$ $M \sim 10^2 \sim 10^3 M_{\text{sol}}$
 $\lambda \gg 2c_s^2/G \sim 16.6 (T/10\text{K}) M_{\text{sol}}/\text{pc}$

- (1) Would stars form throughout the filament, or would the formation propagate along the filament?
- (2) How efficient is protostellar feedback in destroying or disrupting a filament?
- (3) Would the geometry of long filamentary clouds affect the properties of a protostellar cluster, such as the protostellar mass function (PMF) and the companion multiplicity?

Model and Method

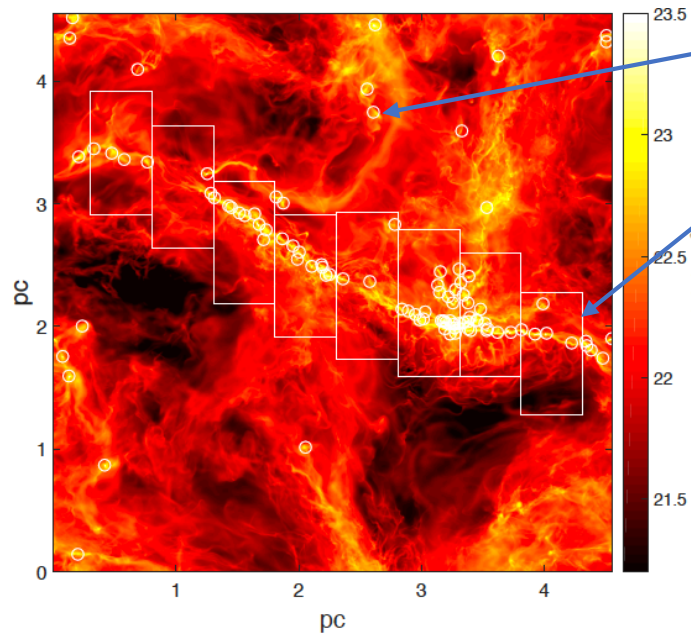
Numerical Methods: orion2 adaptive mesh refinement (AMR) code

Turbulence: $M=10$, $M_A=1$ driving turbulence, $T=10\text{K}$, $M=3110 M_{\text{sol}}$, $V=(4.55\text{pc})^3$

sink particle:

stellar feedback: Luminosity + outflow (momentum+energy)

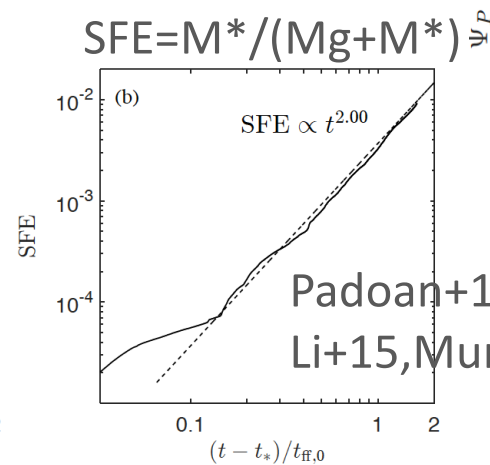
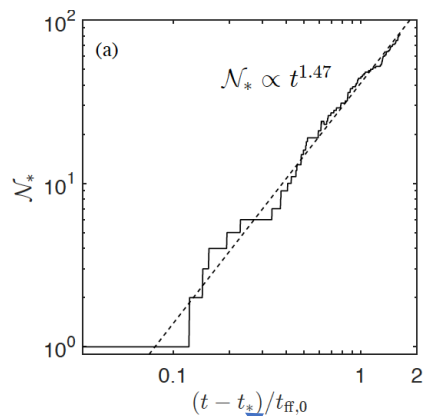
$\dot{m} = \dot{M} \epsilon_j / \dot{M} \text{ acc} = 0.3$; outflow speed $v_w = \min[v_{\text{Kep}}(R_*)/3, 100\text{km/s}]$



100 most massive clumps
zoom-in region=refine upto 6 levels

$$\epsilon_{\text{ff}} \equiv \frac{\dot{M}_{*,\text{obs}}}{\dot{M}_g/t_{\text{ff}}}$$

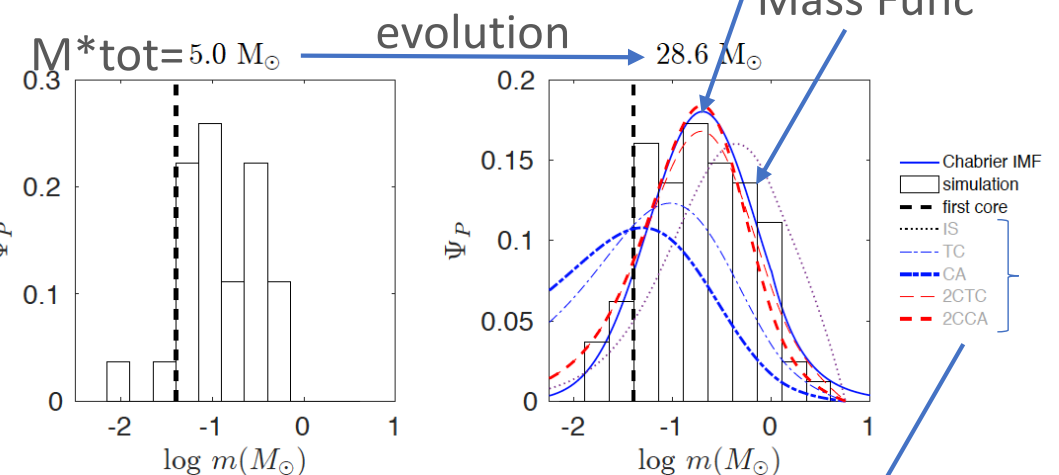
Krumholz+12 \rightarrow high z gal 0.015
Heyer+16 \rightarrow Class I 0.02
 $\epsilon_{\text{ff}}=0.067$



$$\text{SFE} = \dot{M}^*/(\dot{M}_g + \dot{M}^*) \Psi_P$$

Padoan+14 \rightarrow 1.75
Li+15, Murray+17 \rightarrow 2

1st protostar is made



Chabrier IMF05
Mass Func

theoretical IMF

Luminosity Prob \rightarrow resolved

25. The properties of the inner disk around HL Tau: Multi-wavelength modeling of the dust emission by Yao Liu+ AAP in press

Carrasco-Gonzalez+2016

VLA 7mm cont.

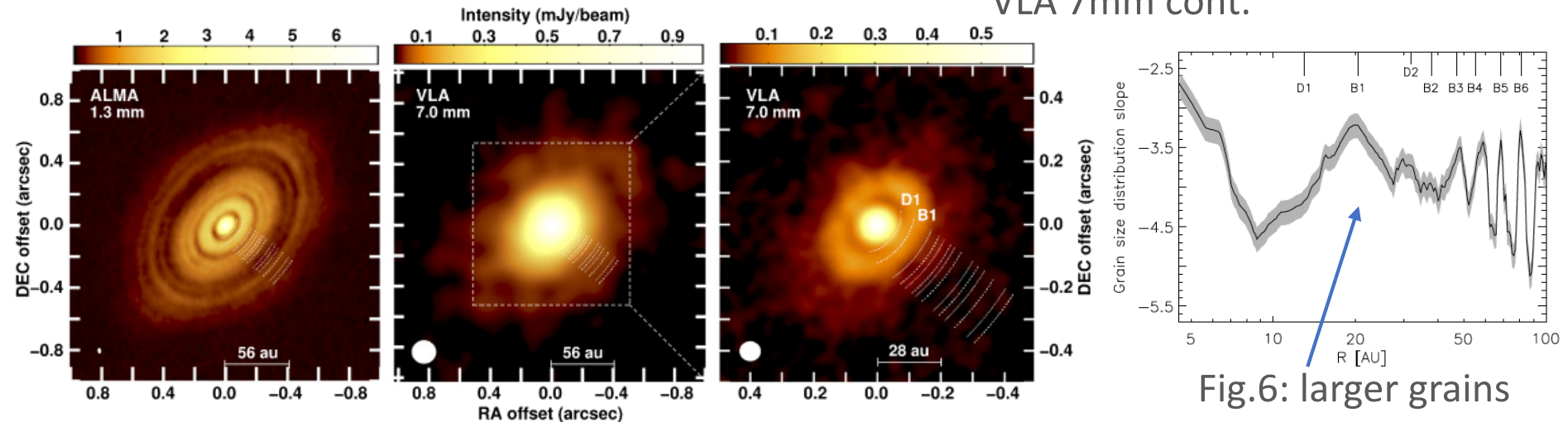


Figure 1. Comparison between the ALMA and VLA observations of the HL Tau disk. Left: ALMA image at 1.3 mm. Center: VLA image at 7.0 mm with an angular resolution of ~ 20 au ($0''.15$; tapered image). Right: close-up of the center of the disk. VLA image at 7.0 mm with an angular resolution of ~ 10 au ($\sim 0''.07$; natural weighting). In all panels, the positions of the reported dark (D1–D7; dotted lines) and bright (B1–B7; dashed lines) rings from the ALMA images (ALMA Partnership et al. 2015a) are shown. The inner disk and the first pair of dark (D1) and bright (B1) rings are clearly seen in the 7.0 mm images.

Fig.6: larger grains are trapped in the first bright ring.

ALMA wavelengths of 2.9, 1.3, and 0.87mm revealed a pattern of bright and dark rings.

- (1) embedded planets in the gaps (Dong+2015; Dipierro+ 2015; Picogna & Kley 2015)→ Large Binocular Telescope, however, excluded the presence of massive planets ($\sim 10 - 15 M_{Jup}$) in two gaps $\sim 70 AU$ (Testi+2016)
- (2) non ideal MHD (Flock+2015)
- (3) sintering (焼結) -induced dust rings (Okuzumi+ 2016),
- (4) dust coagulation triggered by condensation zones of volatiles (Zhang et al. 2015),
- (5) secular gravitational instability (Takahashi & Inutsuka 2016).

Since radiation $\lambda=2.9mm$ optically thick, longer wavelength $\lambda=7mm$ (JVLA)

Flared disk

$\xi > 0$ smaller dust has a larger scale height

$$\rho_{\text{dust}} \propto \Sigma(R) \times \exp \left[-\frac{1}{2} \left(\frac{z}{h(R)} \right)^2 \right], \quad H_{100}(a) = H_{100}(a_{\text{min}}) \left(\frac{a}{a_{\text{min}}} \right)^{-\xi}.$$

$\rightarrow \beta > 1$

dust size distribution $a = (0.01 \mu\text{m}, 7 \text{mm})$

$$h(R) = H_{100} \left(\frac{R}{100 \text{ AU}} \right)^{\beta},$$

$$\frac{dn(a)}{da} \propto a^p da, \quad p = -3.5 \text{ or spatially varies}$$

Central Star $M_* = 1.7 M_{\odot}$, $T = 4000 \text{K}$, $L = 11 L_{\odot}$

Disk parameters			
β	1.15	$1.14^{+0.018}_{-0.007}$	$1.08^{+0.03}_{-0.005}$
$H_{100}(a_{\text{min}}) [\text{AU}]$	15	$13.7^{+1.6}_{-1.0}$	$10.6^{+4.3}_{-1.1}$
$M_{\text{dust}} [10^{-3} M_{\odot}]$	1.0	$1.0^{+0.12}_{-0.1}$	$1.0^{+0.07}_{-0.16}$
p	-3.5	—	—
p_{inner}	—	$-3.45^{+0.08}_{-0.17}$	—
p_{outer}	—	$-3.92^{+0.18}_{-0.08}$	—
f_{inner}	—	$0.5^{+0.03}_{-0.11}$	—
$f_{2.82}$	—	—	$0.27^{+0.05}_{-0.07}$
$f_{0.95}$	—	—	$0.15^{+0.07}_{-0.03}$
$f_{0.45}$	—	—	$0.20^{+0.03}_{-0.06}$
$f_{0.29}$	—	—	$0.15^{+0.05}_{-0.05}$
Observational parameters			
$i [^{\circ}]$	46.7		
Position angle $[^{\circ}]$	138		
$D [\text{pc}]$	140		

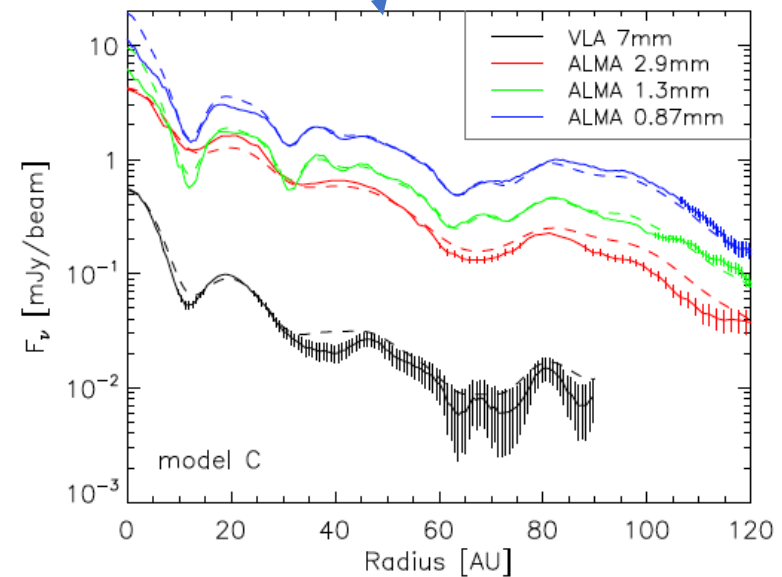
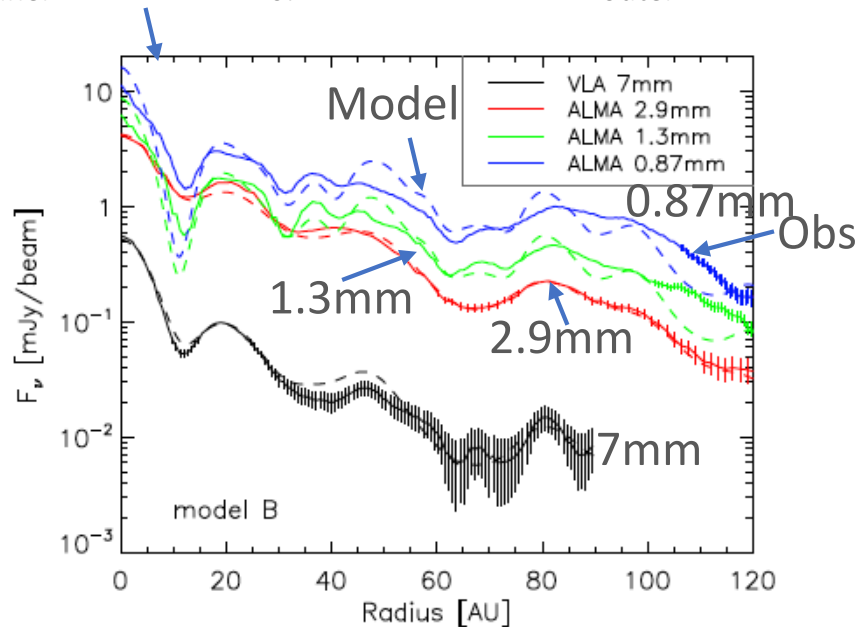
(A) Modeling with a homogeneous grain size distribution \rightarrow NO

(B) Modeling with two power-law grain size distributions

cf. Pinte+2016 model $p = -3.5$ @ $r < 75 \text{AU}$; -4.5 @ $r > 75 \text{AU}$

$p_{\text{inner}} = -3.5^{+0.1}_{-0.2}$ @ $r < 50 \text{AU}$; $p_{\text{outer}} = -3.9^{+0.2}_{-0.1}$ @ $r > 50 \text{AU}$

(C) mass fraction f_{size} @ size = 2.82, 0.95, 0.45, and 0.29mm



Notes. (1) Model A is used in the derivation of the surface density profile (see Sect. 3). (2) Model B and C correspond to the best-fit models under an assumption of two power-law grain size distributions (see Sect. 5.2) and a complex grain size distribution (see Sect. 5.3), respectively.

$$\Sigma(r) \propto r^{-q}$$

$$h(r) = h_0(r/r_0)^\gamma,$$

$$n(a) \propto a^{-3.5}$$

Model Parameters

Parameter	Value	Unit
T_*	10,000	K
R_*	2.5	R_\odot
Inclination	27	Degree
r_{in}	0.5	au
r_{out}	400	au
q	1.2	...
r_0	100	au
h_0	8.5	au
γ	1.125	...
M_{dust}	$1.2\text{e-}4$	M_\odot
a_{min}	0.01	μm
a_{max}	1.0	μm

Note. Emission of the star is assumed to be blackbody.

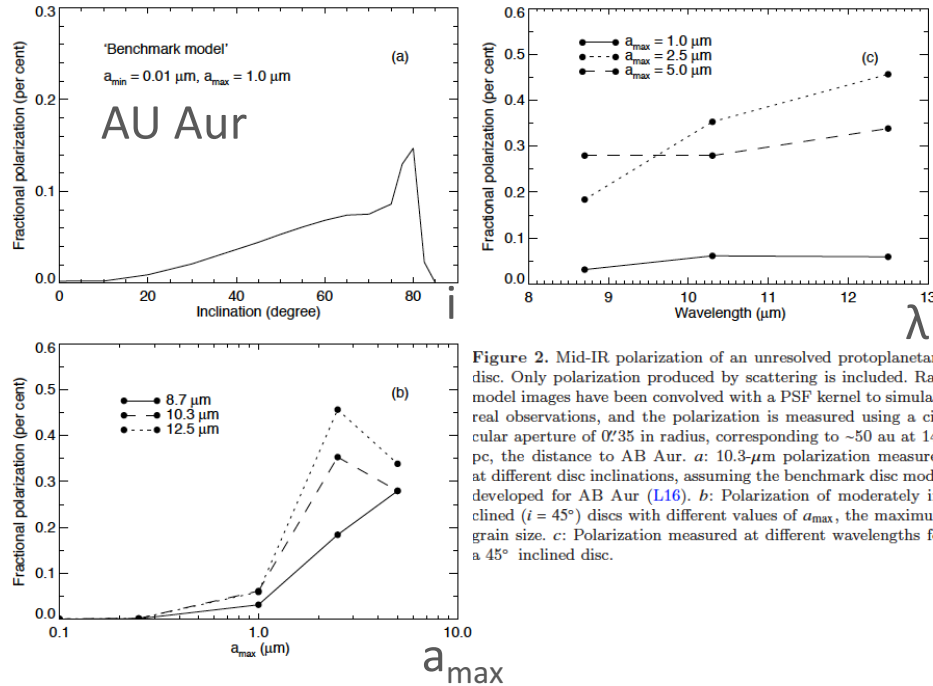


Figure 2. Mid-IR polarization of an unresolved protoplanetary disc. Only polarization produced by scattering is included. Raw model images have been convolved with a PSF kernel to simulate real observations, and the polarization is measured using a circular aperture of $0''.35$ in radius, corresponding to ~ 50 au at 144 pc, the distance to AB Aur. *a*: $10.3\text{-}\mu\text{m}$ polarization measured at different disc inclinations, assuming the benchmark disc model developed for AB Aur (L16). *b*: Polarization of moderately inclined ($i = 45^\circ$) discs with different values of a_{max} , the maximum grain size. *c*: Polarization measured at different wavelengths for a 45° inclined disc.