

Constraining the ellipticity of magnetars powering superluminous supernovae

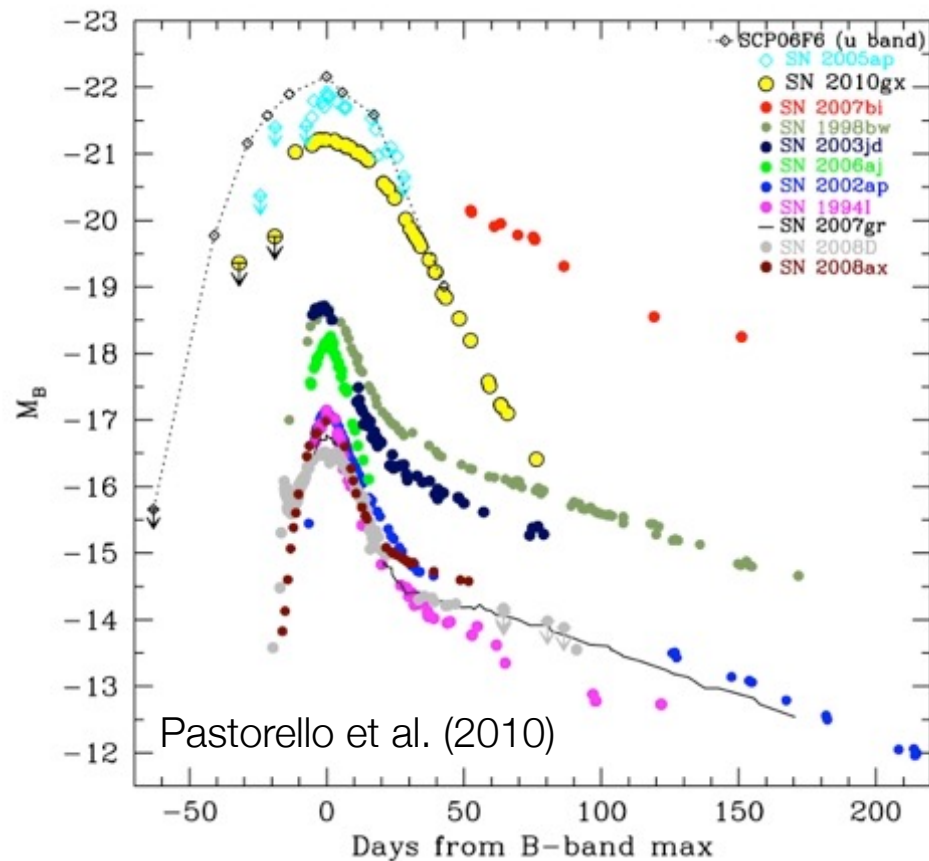
Takashi Moriya (NAOJ)

Thomas Tauris (Bonn)

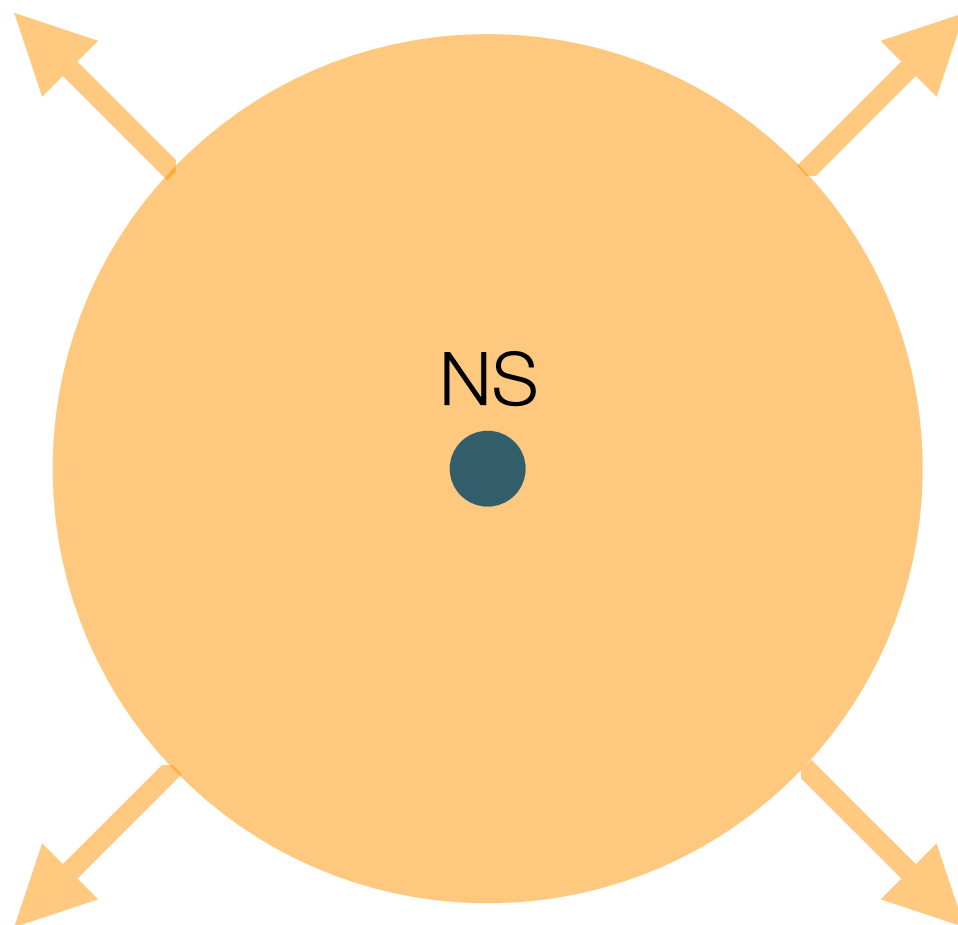
MNRAS, 460, L55 (2016)

Superluminous supernovae (SLSNe)

- SNe brighter than ~ -21 mag (or $\sim 1e44$ erg/s)
- total radiation energy exceeds $1e51$ erg
 - usual SNe emit $\sim 1e49$ erg
 - comparable to kinetic energy of usual SNe

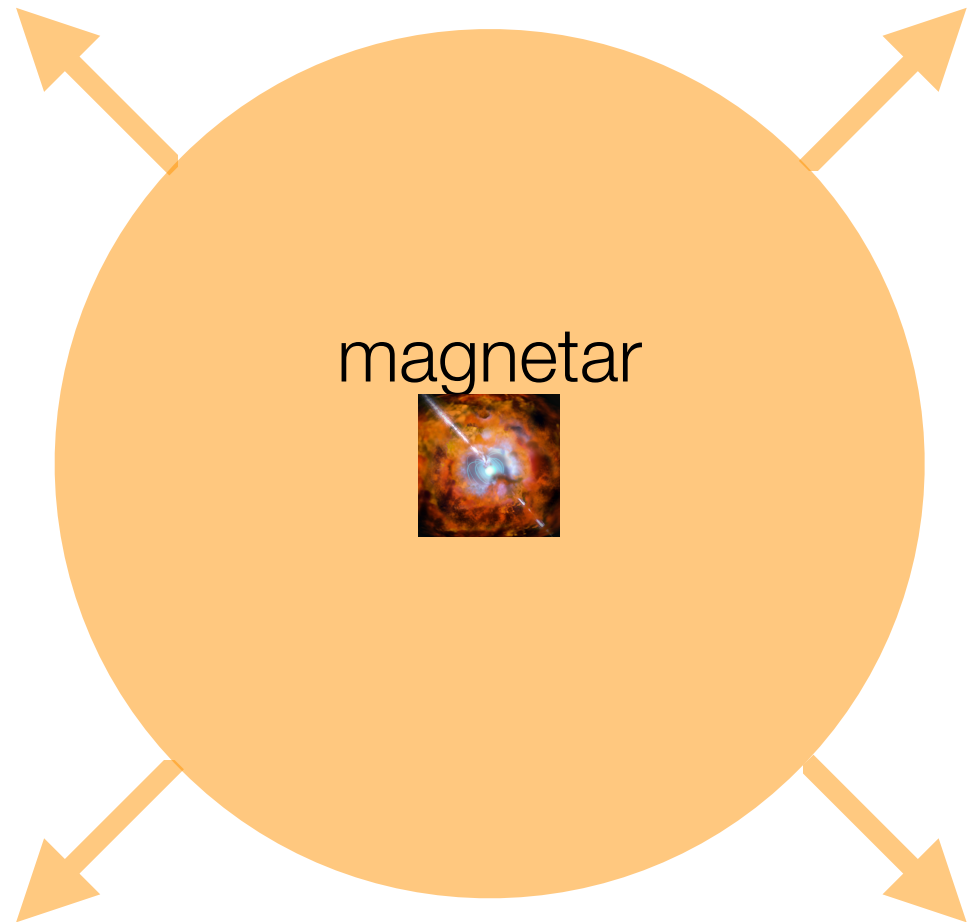


Magnetar model for SLSNe



Magnetar model for SLSNe

- magnetar



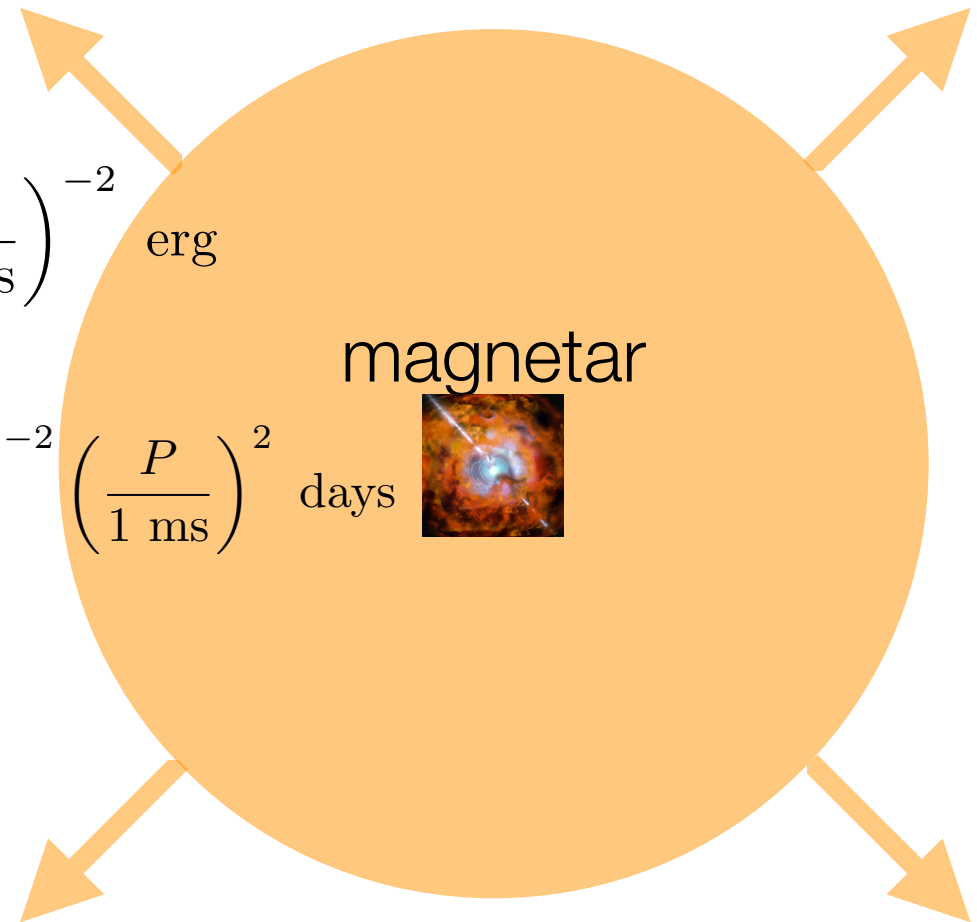
Magnetar model for SLSNe

- magnetar
 - neutron star with
 - rapid rotation

$$E_{\text{rot}} = \frac{1}{2} I_{\text{NS}} \Omega^2 \simeq 2 \times 10^{52} \left(\frac{P}{1 \text{ ms}} \right)^{-2} \text{ erg}$$

- strong magnetic field

$$t_m = \frac{6 I_{\text{NS}} c^3}{B_{\text{dipole}}^2 R_{\text{NS}}^6 \Omega^2} \simeq 5 \left(\frac{B_{\text{dipole}}}{10^{14} \text{ G}} \right)^{-2} \left(\frac{P}{1 \text{ ms}} \right)^2 \text{ days}$$



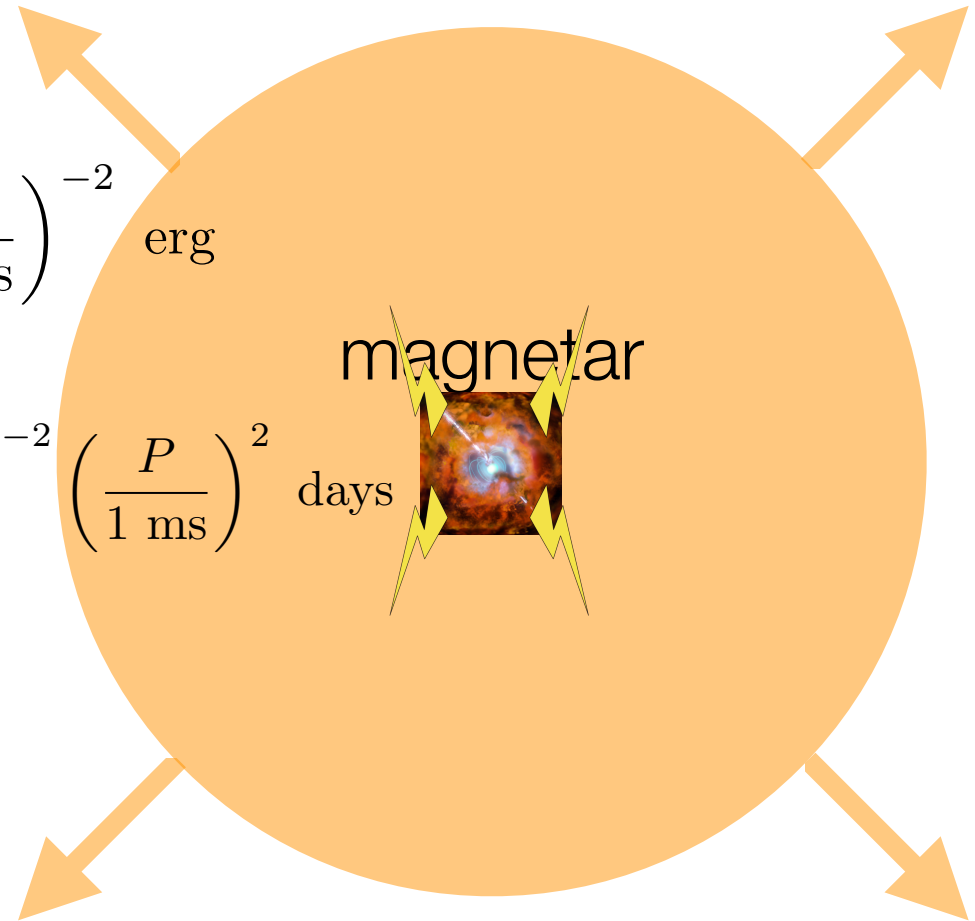
Magnetar model for SLSNe

- magnetar
 - neutron star with
 - rapid rotation

$$E_{\text{rot}} = \frac{1}{2} I_{\text{NS}} \Omega^2 \simeq 2 \times 10^{52} \left(\frac{P}{1 \text{ ms}} \right)^{-2} \text{ erg}$$

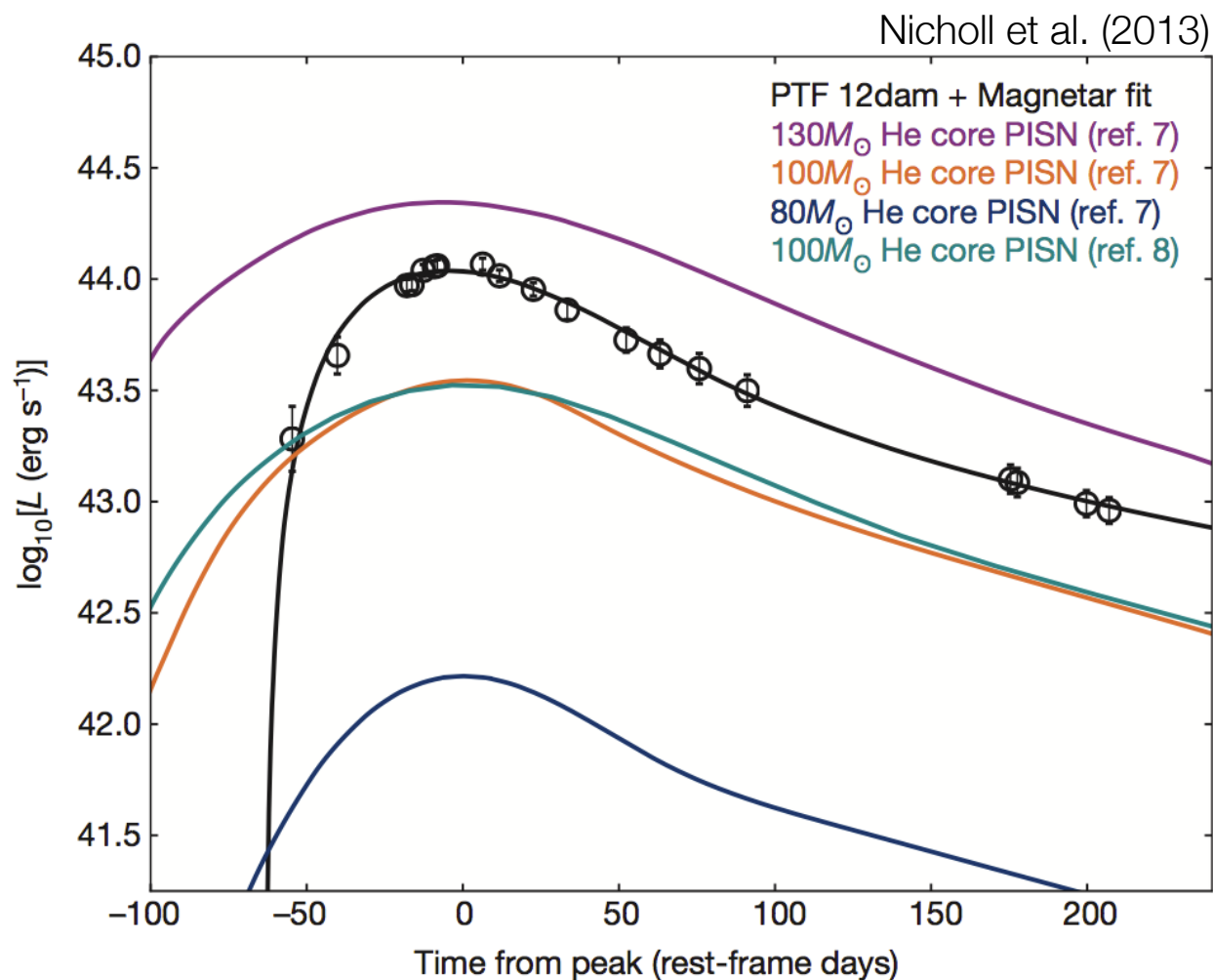
- strong magnetic field

$$t_m = \frac{6 I_{\text{NS}} c^3}{B_{\text{dipole}}^2 R_{\text{NS}}^6 \Omega^2} \simeq 5 \left(\frac{B_{\text{dipole}}}{10^{14} \text{ G}} \right)^{-2} \left(\frac{P}{1 \text{ ms}} \right)^2 \text{ days}$$



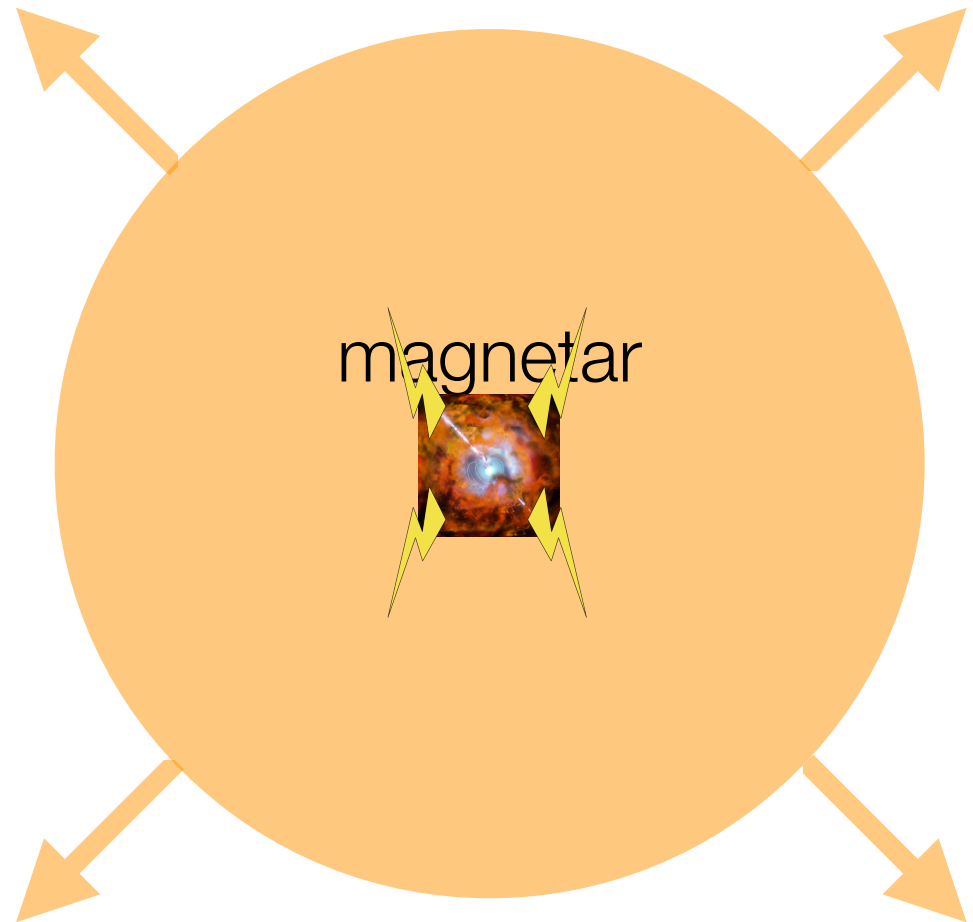
Light curve models

- $B = 1e14$ G, $P = 2.6$ ms



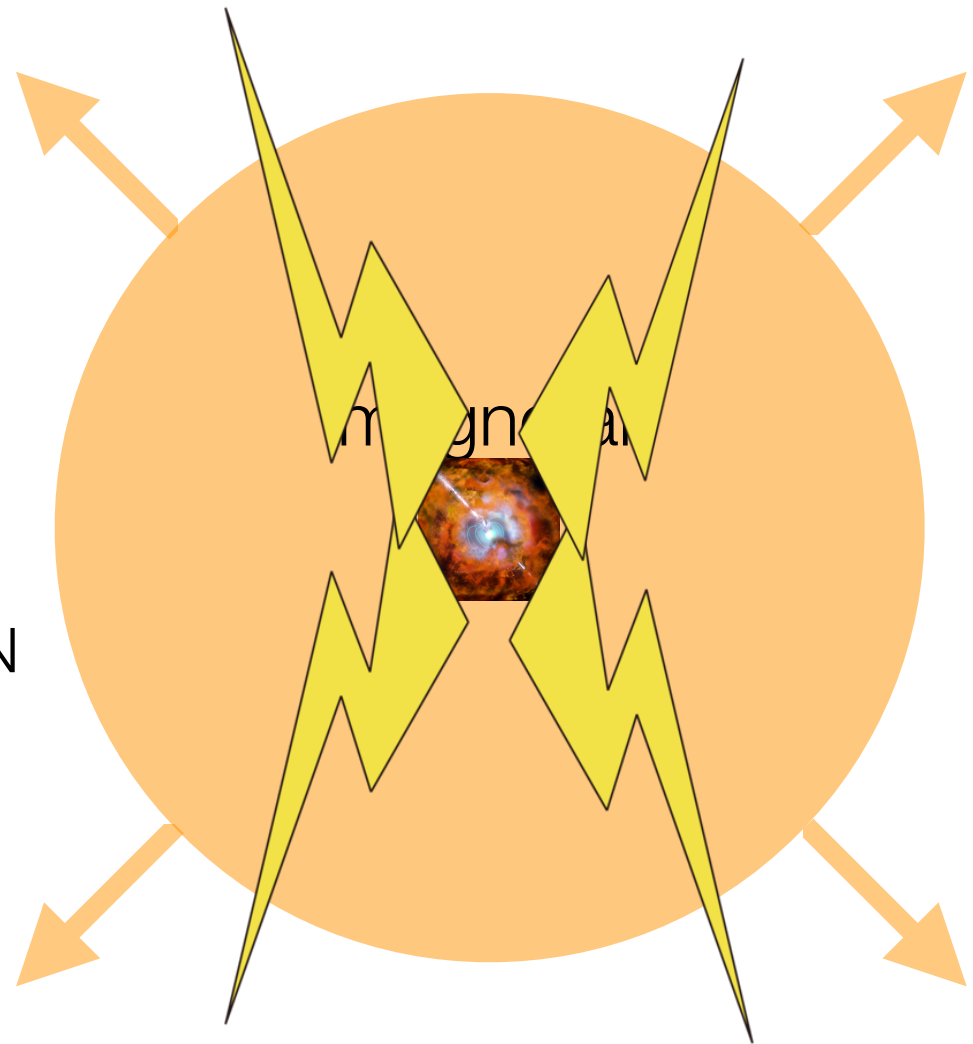
NSs are GW sources

- magnetars need to convert rotational energy to thermal energy almost in situ for SLSNe
- if magnetars are highly distorted, GW emission will be dominant way to lose their rotational energy



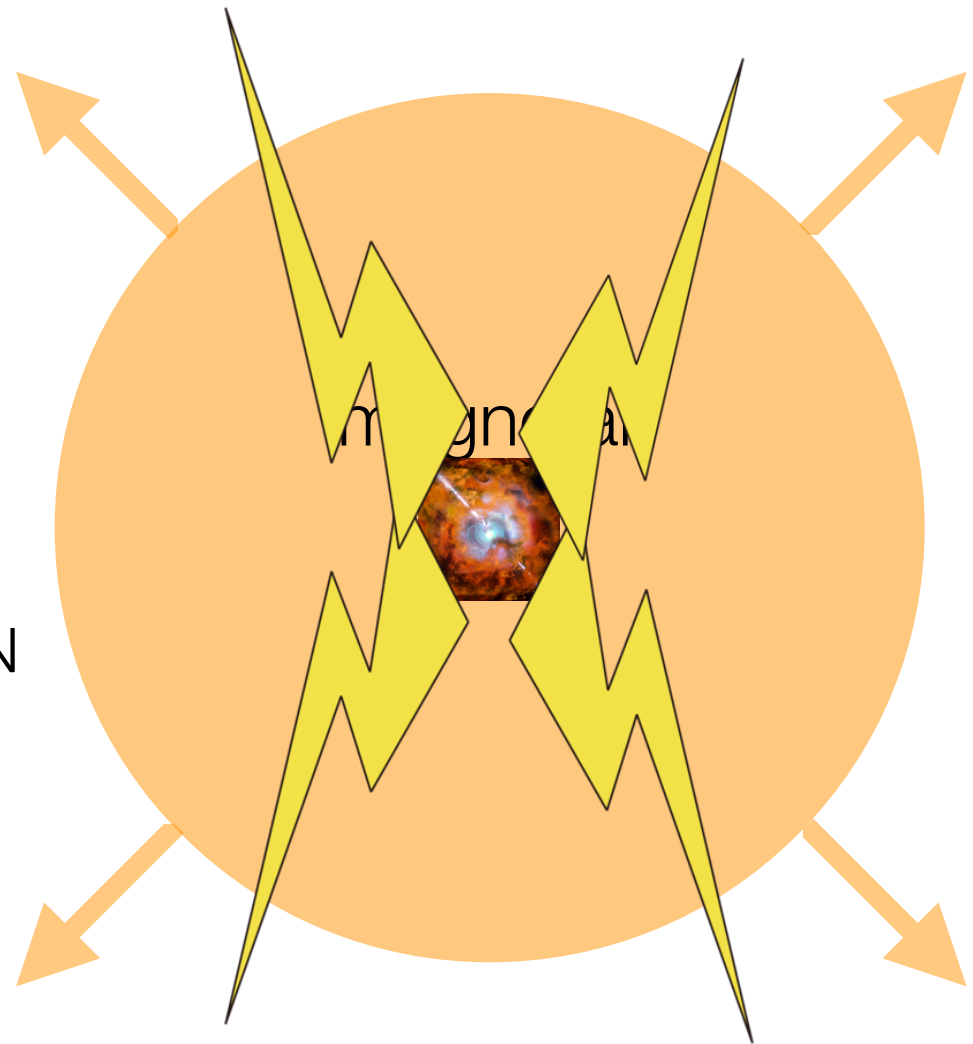
NSs are GW sources

- magnetars need to convert rotational energy to thermal energy almost in situ for SLSNe
- if magnetars are highly distorted, GW emission will be dominant way to lose their rotational energy
- GW emission will just leave the SN ejecta without heating the ejecta.



NSs are GW sources

- magnetars need to convert rotational energy to thermal energy almost in situ for SLSNe
- if magnetars are highly distorted, GW emission will be dominant way to lose their rotational energy
- GW emission will just leave the SN ejecta without heating the ejecta.
- magnetars in SLSNe should not be highly distorted!



EM emission vs GW emission

- EM emission timescale (dipole)

$$\tau_{\text{EM}} = \frac{3}{4\pi^2} \frac{I c^3 P_0^2}{B_{\text{dipole}}^2 R^6 \sin^2 \alpha}$$

- GW emission timescale

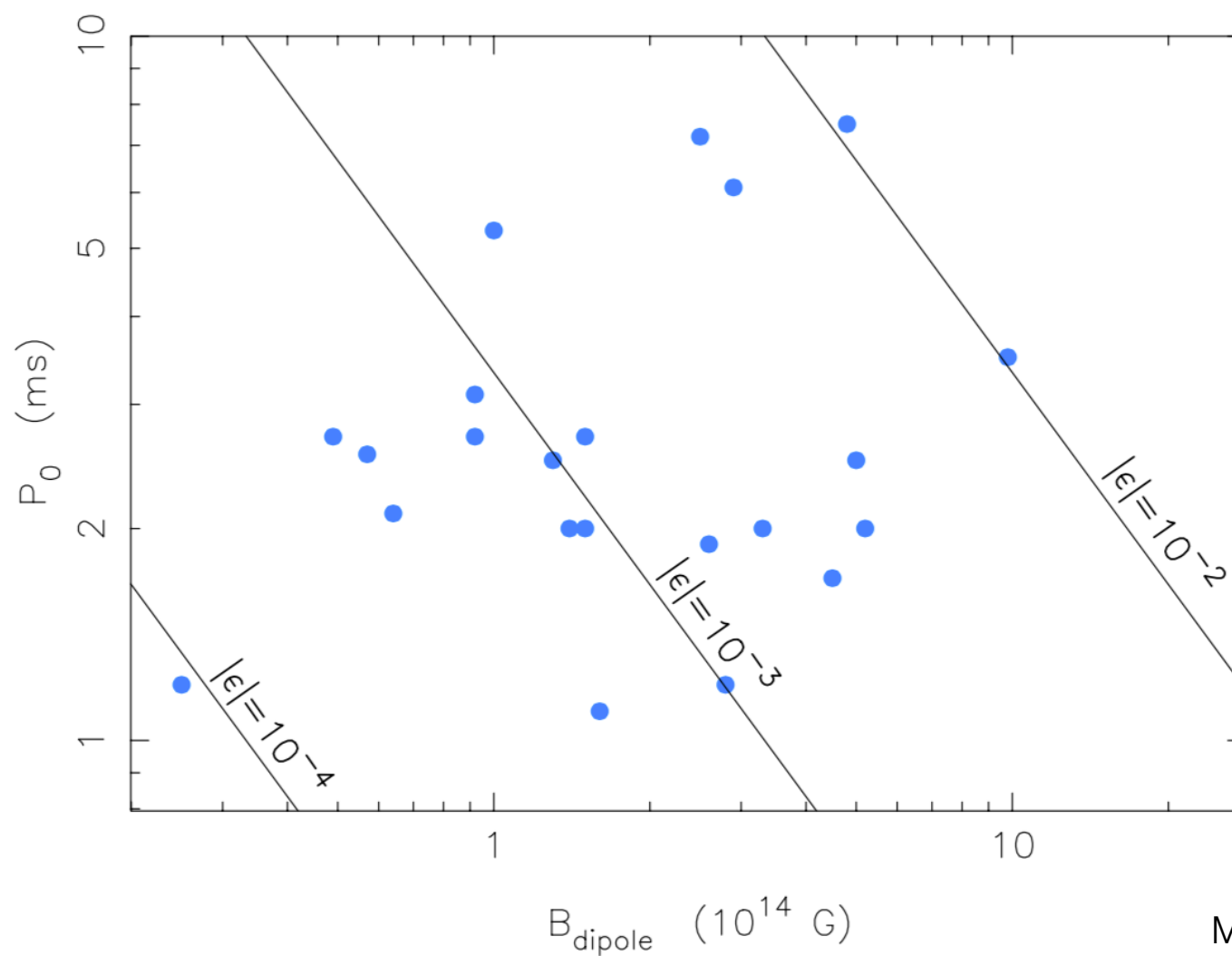
$$\tau_{\text{GW}} = \frac{5}{2^{10} \pi^4} \frac{c^5 P_0^4}{G I \varepsilon^2}$$

- for SLSNe

$$\tau_{\text{GW}} > \tau_{\text{EM}}$$

Constraint on magnetar ellipticities for SLSNe

$$|\varepsilon| < \sqrt{\frac{5}{3G} \frac{cR^3 P_0 B_{\text{dipole}}}{2^4 \pi I}} \simeq 3.0 \times 10^{-4} \left(\frac{B_{\text{dipole}}}{10^{14} \text{ G}} \right) \left(\frac{P_0}{1 \text{ ms}} \right)$$

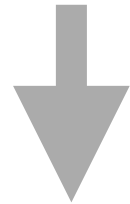


Probable origin of ellipticities: toroidal fields

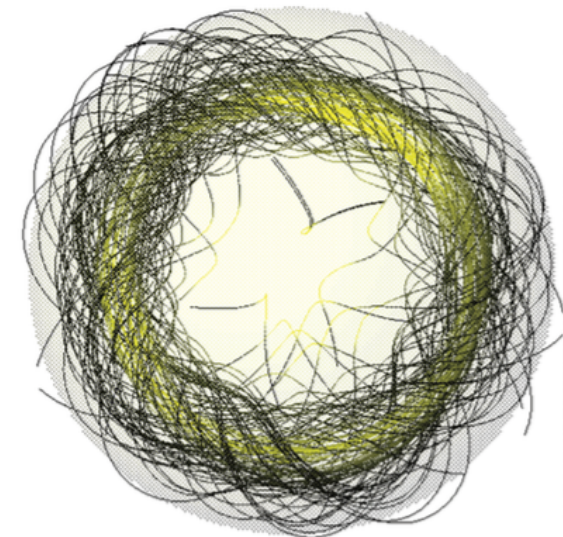
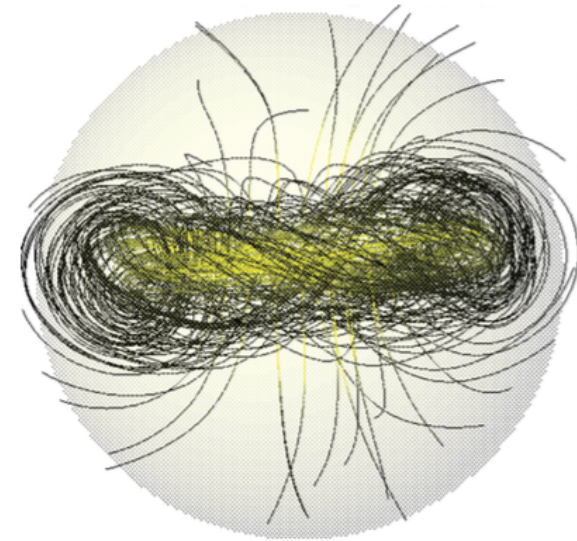
$$|\varepsilon| \simeq 1.6 \times 10^{-4} (B_{\text{toroidal}} / 10^{16} \text{ G})^2$$

Cutler (2002)

$$|\varepsilon| < \sqrt{\frac{5}{3G}} \frac{cR^3 P_0 B_{\text{dipole}}}{2^4 \pi I} \simeq 3.0 \times 10^{-4} \left(\frac{B_{\text{dipole}}}{10^{14} \text{ G}} \right) \left(\frac{P_0}{1 \text{ ms}} \right)$$



$$B_{\text{toroidal}} \lesssim 1.4 \times 10^{16} \text{ G} \left(\frac{B_{\text{dipole}}}{10^{14} \text{ G}} \right)^{1/2} \left(\frac{P_0}{1 \text{ ms}} \right)^{1/2}$$



Braithwaite (2009)

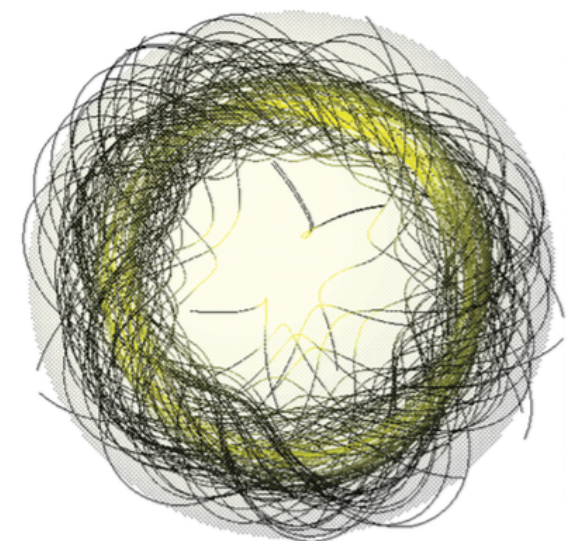
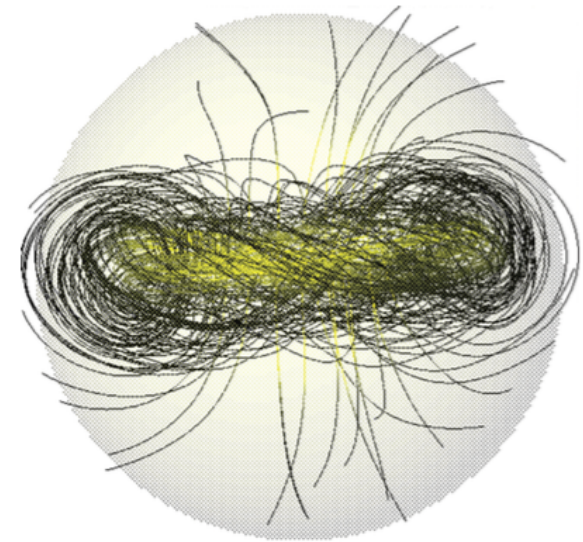
Constraint on toroidal field strengths

SN name	B_{dipole} 10^{14}G	P_0 ms	$ \varepsilon $ 10^{-3}	B_{toroidal} 10^{16}G	Reference
SN 2005ap	0.92	3.1	<0.85	<2.4	Chatzopoulos et al. (2013)
SCP06F6	1.3	2.5	<0.98	<2.5	Chatzopoulos et al. (2013)
SNLS 06D4eu	1.4	2.0	<0.85	<2.3	Howell et al. (2013)
SN 2007bi	0.92	2.7	<0.75	<2.2	Chatzopoulos et al. (2013)
SN 2010gx	5.2	2.0	<3.1	<4.5	Inserra et al. (2013)
SN 2010kd	1.5	2.7	<1.2	<2.8	Chatzopoulos et al. (2013)
SN 2010kl	9.8	3.5	<10	<8.2	Bersten et al. (2016)
PTF10hgi	2.5	7.2	<5.4	<5.9	Inserra et al. (2013)
SN 2011ke	4.5	1.7	<2.3	<3.9	Inserra et al. (2013)
SN 2011kf	3.3	2.0	<2.0	<3.6	Inserra et al. (2013)
PTF11rks	4.8	7.5	<11	<8.4	Inserra et al. (2013)
SN 2012il	2.9	6.1	<5.3	<5.9	Inserra et al. (2013)
PTF12dam	0.49	2.7	<0.39	<1.6	Chen et al. (2015)
CSS121015	1.5	2.0	<0.90	<2.4	Nicholl et al. (2014)
LSQ12dlf	2.6	1.9	<1.5	<3.1	Nicholl et al. (2014)
SSS120810	2.8	1.2	<1.0	<2.6	Nicholl et al. (2014)
SN 2013dg	5.0	2.5	<3.7	<5.0	Nicholl et al. (2014)
iPTF13ajg	1.6	1.1	<0.54	<1.9	Vreeswijk et al. (2014)
iPTF13ehe	0.57	2.55	<0.43	<1.7	Wang et al. (2015)
DES13S2cmm	1.0	5.3	<1.6	<3.2	Papadopoulos et al. (2015)
SN 2015bn	0.64	2.1	<0.40	<1.6	Nicholl et al. (2016)
ASASSN-15lh	0.25	1.2	<0.090	<0.77	Bersten et al. (2016)

$B_{\text{toroidal}} < \sim 1e16\text{ G}$

Toroidal fields are required for stable poloidal fields

- $B_{\text{dipole}} \sim 1e14 \text{ G}$ (from LC modeling)
- $B_{\text{toroidal}} < \sim 1e16 \text{ G}$ (from ellipticities)
- $B_{\text{dipole}}/B_{\text{toroidal}} > \sim 0.01$
 - still OK (dipole field can be stable)
 - all SLSNe must have stable dipole if they are powered by magnetars
 - a test for the SLSN magnetar model



Summary

- SLSNe may be powered by magnetars
 - $B_{\text{dipole}} \sim 1e14$ G, $P \sim 1$ ms
- magnetars must lose their rotational energy as EM emission
 - magnetar ellipticities must be small enough to avoid GW emission
- magnetar ellipticities must be less than $\sim 1e-3$
 - $B_{\text{toroidal}} < \sim 1e16$ G
- $B_{\text{dipole}}/B_{\text{toroidal}} > \sim 0.01$
 - dipole field can be stable
- a test for SLSN magnetar model

